ALGEBRA, GEOMETRY AND MATHEMATICAL PHYSICS
BANACH CENTER PUBLICATIONS, VOLUME 93
INSTITUTE OF MATHEMATICS
POLISH ACADEMY OF SCIENCES
WARSZAWA 2011

## SOME SPECIAL FUNCTIONS RELATED TO UNIMODULAR PSEUDO-ORTHOGONAL GROUPS

## ALEXANDER NIZHNIKOV

Moscow Pedagogical State University
M. Piroqovskaya 1, 119991 Moscow, Russia

## ILYA SHILIN

Sholokhov Moscow State University for the Humanities
V. Radishevskaya 16–18, 109240 Moscow, Russia, and
Moscow Aviation Institute (University of Aerospace Technology)
Volokolamskoe shosse 4, 125993 Moscow, Russia
E-mail: ilyashilin@li.ru

**Abstract.** We obtain some matrix elements of basis transformations in a representation space of the unimodular pseudo-orthogonal group. Using these elements, we derive some formulas for special functions.

1. Introduction. Throughout this paper, special functions occur as matrix elements of basis transformations and representation operators. We construct some elements of the matrices which connect different bases for class 1 representations of the unimodular pseudo-orthogonal group. These matrix elements are expressed in terms of Whittaker functions for the case SO(2,2) and in terms of Vilenkin function for the general case. In this way, some integral relations are obtained for Gauss, Whittaker and Macdonald functions.

Let us assume that the linear space  $\mathbb{R}^{p+q}$  is endowed with the form

$$\vartheta(x) := x_1^2 + \ldots + x_p^2 - x_{p+1}^2 - \ldots - x_{p+q}^2.$$

We denote by X the cone  $\vartheta(x) = 0$  without the origin. By definition, the unimodular pseudo-orthogonal group SO(p,q) consists of all linear transformations of  $\mathbb{R}^{p+q}$  preserving  $\vartheta(x)$ . In the case p = 1 or q = 1, we have the special Lorentz group. If  $g \in SO(p,q)$ , then

2010 Mathematics Subject Classification: Primary 33C80; Secondary 22E30, 22E46.

Key words and phrases: unimodular pseudo-orthogonal group, representation of group, basis transformation, matrix element, special functions.

[117]

The paper is in final form and no version of it will be published elsewhere.

 $\det g = 1$  and the equation

$$ge_{p,q}g^{\mathrm{T}} = e_{p,q} \tag{1}$$

holds, where  $e_{p,q} := \text{diag}(1,\ldots,1-1,\ldots,-1)$  and tr  $e_{p,q} = p-q$ . The group SO(p,q) has 2 connected components. One of them consists of the matrices

$$g \equiv \begin{pmatrix} A & B \\ C & D \end{pmatrix},\tag{2}$$

where A is a  $p \times p$  matrix such that sign det A = sign det D = 1. This coset contains the identity element and will be under our consideration further. We denote this subgroup by G. The fixed points g under the Cartan involution  $G \to G$ ,  $g \mapsto e_{p,q} g e_{p,q}$  form a maximal compact subgroup  $K \cong SO(p) \times SO(q)$ .

Let us consider the K-orbit of the point  $(1,0,\ldots,0,1)$  on X. It is a direct product of two spheres; we denote it as  $\Gamma_K$ . The measure  $dx := \prod_{i=1}^{p+q} dx_i$  in  $\mathbb{R}^{p+q}$  is invariant with respect to  $SL(p+q,\mathbb{R})$ , the generalized function

$$\delta\left(\sqrt{x_1^2 + \ldots + x_p^2} - 1\right) \cdot \delta\left(\sqrt{x_{p+1}^2 + \ldots + x_{p+q}^2} - 1\right)$$

is invariant with respect to K, and the polynomials  $\sum_{i=1}^{p} x_i^2$  and  $\sum_{i=p+1}^{p+q} x_i^2$  are both symmetric. This leads to the following K-invariant measure on  $\Gamma_K$ :

$$(\mathrm{d}x)_K := \frac{\mathrm{d}_{\theta(1)} \dots \mathrm{d}_{\theta(p-1)} \mathrm{d}_{\vartheta(p+1)} \dots \mathrm{d}_{\vartheta(p+q-1)}}{|x_{\theta(p)}| |x_{\vartheta(p+q)}|},$$

where  $\theta$  is any permutation of the set  $\{1, \ldots, p\}$  and  $\vartheta$  is any permutation of the set  $\{p+1, \ldots, p+q\}$ . In the spherical coordinate system

$$\begin{cases} x_1 = \sin \phi_1 \dots \sin \phi_{p-1}, \\ x_2 = \sin \phi_1 \dots \sin \phi_{p-2} \cos \phi_{p-1}, \\ \dots \\ x_{p-1} = \sin \phi_1 \cos \phi_2, \\ x_p = \cos \phi_1, \\ x_{p+1} = \sin \psi_1 \dots \sin \psi_{q-1}, \\ x_{p+2} = \sin \psi_1 \dots \sin \psi_{q-2} \cos \psi_{q-1}, \\ \dots \\ x_{p+q-1} = \sin \psi_1 \cos \psi_2, \\ x_{p+q} = \cos \psi_1, \end{cases}$$

we obtain

$$(\mathrm{d}x)_K = \prod_{i=1}^{p-2} \sin^{p-i-1} \phi_i \, \mathrm{d}\phi_i \cdot \prod_{i=1}^{q-2} \sin^{q-i-1} \phi_i \, \mathrm{d}\psi_i \cdot \mathrm{d}\phi_{p-1} \, \mathrm{d}\psi_{q-1},$$

if  $p \geq 2, q \geq 2$ . Here  $\phi_1, \psi_1 \in [0; 2\pi)$  and  $\phi_2, \psi_2, \phi_3, \psi_3, \ldots \in [0; \pi)$ . So we have the following corollary for the G-invariant measure on X:

$$dx := \frac{dx_{\zeta(1)} \dots dx_{\zeta(p+q-1)}}{|x_{\zeta(p+q)}|},$$
(3)

where  $\zeta \in \mathbf{S}_{p+q}$ .

Let us introduce the binary relations  $\Omega_1$  on the set  $\{1,\ldots,p\}$  and  $\Omega_2$  on the set  $\{p+1,\ldots,p+q\}$  by the condition that  $(i,j)\in\Omega_l$  is equivalent to i< j. Let us consider now a rotation r(i,j,t) through angle t on the  $(x_i,x_j)$  plane, where the condition  $(i,j)\in\Omega_1$  or  $(i,j)\in\Omega_2$  holds. Let k(i,j) be the infinitesimal matrix of r(i,j,t). Then its matrix elements are

$$k(i,j)_{st} = \begin{cases} 1, & \text{if } (s,t) = (i,j), \\ -1, & \text{if } (s,t) = (j,i), \\ 0, & \text{if } (s,t) \neq (i,j) \text{ and } (s,t) \neq (j,i). \end{cases}$$

The system of vectors  $k(i,j), (i,j) \in \Omega_1 \cup \Omega_2$ , is linearly independent; the dimension of their linear span  $\mathfrak{k}$  is equal to  $\frac{p^2+q^2-p-q}{2}$ . Moreover,  $\mathfrak{k}$  is a Lie algebra because

$$[k(i,j),k(\tilde{i},\tilde{j})] = \begin{cases} 0, \text{ if } \delta_{i\tilde{i}}\delta_{j\tilde{j}} = 1,\\ 0, \text{ if } (i-\tilde{i})(j-\tilde{i}(j-\tilde{j}) \neq 0,\\ k(j,\tilde{j}), \text{ if } i = \tilde{i},\\ k(i,\tilde{i}), \text{ if } j = \tilde{j},\\ -k(i,\tilde{j}), \text{ if } j = \tilde{i}. \end{cases}$$

Let  $\hat{r}(i, j, t)$  be a hyperbolic rotation through angle t on the  $(x_i, x_j)$  plane, where  $(i, j) \in \{1, \dots, p\} \times \{p + 1, \dots, p + q\}$ . If the infinitesimal matrix h(i, j) corresponds to the rotation  $\hat{r}(i, j, t)$ , then its matrix elements are

$$h(i,j)_{st} = \begin{cases} 1, & \text{if } (i,j) = (s,t), \\ 1, & \text{if } (i,j) = (t,s), \\ 0, & \text{if } (i,j) \neq (s,t) \text{ and } (i,j) \neq (t,s). \end{cases}$$

These infinitesimal matrices are linearly independent and generate a linear space  $\mathfrak{h}$ , and dim  $\mathfrak{h} = pq$ .

It is not hard to prove that  $\det A \neq 0$  for A in (2). From (1) and (2), we have  $B = (A^{-1})^{\mathrm{T}}C^{\mathrm{T}}D$ . This means that any matrix  $g \in G$  depends on  $\binom{p+q}{2} = \frac{(p+q)(p+q-1)}{2}$  parameters. Therefore,  $\mathfrak{k} \oplus \mathfrak{h}$  is the tangent space of G. It is easy to verify that  $[\mathfrak{k}, \mathfrak{h}] \subset \mathfrak{h}$  and  $[\mathfrak{h}, \mathfrak{h}] \subset \mathfrak{k}$ . We denote the group  $\exp \mathfrak{h}$  by H.

It is clear that K acts transitively on  $\Gamma_K$ . Let us note that, first, G is generated by the subgroups K and H [M] and, second,  $\hat{r}^{-1}(i,j,t)(1,0,\ldots,0,1) = (e^{-t},0,\ldots,0,e^{-t})$ . This means that the action  $x \mapsto g^{-1}x$  of the group G is transitive on X. Let  $\sigma \in \mathbb{C}$  and  $D_{\sigma}$  be a linear subspace in  $C^{\infty}(X)$  consisting of  $\sigma$ -homogeneous functions. We define the representation  $T_{\sigma}$  in  $D_{\sigma}$  by left shifts:  $T_{\sigma}(g)[f(x)] := f(g^{-1}x)$ . This representation is irreducible if  $\sigma \notin \mathbb{Z}$  [M], [Ve].

For any pair  $(D_{\sigma}, D_{\tilde{\sigma}})$ , we define the bilinear functional

$$\mathsf{F}: (D_{\sigma}, D_{\tilde{\sigma}}) \to \mathbb{C}, \ (f_1, f_2) \mapsto \frac{\Gamma\left(\frac{p+q-1}{2}\right)}{2\pi^{\frac{p+q-1}{2}}} \int_{\Gamma} f_1(x) f_2(x) \, \mathrm{d}x_{\Gamma}.$$

Here  $\Gamma$  is a variety of X intersecting all or almost all generatrices. The words almost all mean here all generatrices except one of them. Every point  $x \in \Gamma$  can be represented as  $\{x_i = F_i(\xi_1, \dots, \xi_{p+q-2}), i = 1, \dots, p+q \}$ . So we can write every point  $x \in X$  as

$$\{x_i = tF_i(\xi_1, \dots, \xi_{p+q-2}), \quad i = 1, \dots, p+q,$$
 (4)

and, consequently,

$$dx = t^{p+q-3} dt d\xi, (5)$$

where  $\tilde{G}$  is a subgroup of G, which acts transitively on  $\Gamma$ , and  $d\xi$  is the  $\tilde{G}$ -invariant measure on  $\Gamma$ .

LEMMA 1.1. If  $\tilde{\sigma} = -\sigma - p - q + 2$ , then F does not depend on  $\Gamma$ .

*Proof.* It follows from homogeneity of the functions  $f_1$  and  $f_2$  and formulas (5) and (3).

2. Bases transforms and Vilenkin function. In addition to  $\Gamma_K$ , let us introduce some other contours on X intersecting almost all generatrices.

We now denote as  $\Gamma_1$  the intersection of the cone X and the the cylinder  $x_1^2 + \ldots + x_{p-1}^2 = 1$ . The subgroup  $H_1 \simeq SO(p-1) \times SO(1,q)$  acts transitively on  $\Gamma_1$ . The  $H_1$ -invariant measure

$$(dx)_{H_1} = \frac{dx_1 \dots dx_{p-2}}{|x_{p-1}|} \cdot \frac{dx_p \dots dx_{p+q-1}}{|x_{p+q}|}$$

on  $\Gamma_1$  follows from (5). In the cylinder coordinate system

$$\begin{cases} x_1 = \sin \phi_1 \dots \sin \phi_{p-3} \sin \phi_{p-2}, \\ x_2 = \sin \phi_1 \dots \sin \phi_{p-3} \cos \phi_{p-2}, \\ \dots \\ x_{p-2} = \sin \phi_1 \cos \phi_2, \\ x_{p-1} = \cos \phi_1, \\ x_p = \sinh \alpha, \\ x_{p+1} = \cosh \alpha \sin \psi_1 \dots \sin \psi_{q-1}, \\ x_{p+2} = \cosh \alpha \sin \psi_1 \dots \cos \psi_{q-1}, \\ \dots \\ x_{p+q} = \cosh \alpha \cos \psi_1, \end{cases}$$

we have

$$(\mathrm{d}x)_{H_1} = \prod_{i=1}^{p-3} \sin^{p-i-2} \phi_i \, \mathrm{d}\phi_i \cdot \prod_{i=1}^{q-2} \sin^{p-i-1} \psi_i \, \mathrm{d}\psi_i \cdot \sinh^{q-2} \alpha \, \mathrm{d}\phi_{p-2} \, \mathrm{d}\psi_{q-1} \, \mathrm{d}\alpha.$$

Here  $\alpha \in [0; +\infty)$ ,  $\phi_1, \psi_1 \in [0; 2\pi)$  and  $\phi_2, \psi_2, \phi_3, \psi_3, \ldots \in [0; \pi)$ .

Let  $\Gamma_2$  be the intersection of X and the cylinder  $x_1^2 + x_{p-1}^2 - x_{p+1}^2 = 1$ . We denote by  $H_2$  the subgroup acting transitively on  $\Gamma_2$ . Thus,  $H_2 \simeq SO(p-1,1) \times SO(1,q-1)$ . From (5), we have

$$(dx)_{H_2} = \frac{dx_1 \dots dx_{p-1}}{|x_{p+1}|} \cdot \frac{dx_{p+2} \dots dx_{p+q}}{|x_p|}.$$

In the cylinder coordinate system

$$\begin{cases} x_1 = \cosh \alpha \sin \phi_1 \dots \sin \phi_{p-3} \sin \phi_{p-2}, \\ x_2 = \cosh \alpha \sin \phi_1 \dots \sin \phi_{p-3} \cos \phi_{p-2}, \\ \dots \\ x_{p-2} = \cosh \alpha \sin \phi_1 \cos \phi_2, \\ x_{p-1} = \cosh \alpha \cos \phi_1, \\ x_p = \sinh \beta, \\ x_{p+1} = \sinh \alpha, \\ x_{p+2} = \cosh \beta \sin \psi_1 \dots \cos \psi_{q-2}, \\ \dots \\ x_{p+q} = \cosh \beta \cos \psi_1, \end{cases}$$

we have

$$(\mathrm{d}x)_{H_2} = \prod_{i=1}^{p-3} \sin^{p-i-2} \phi_i \, \mathrm{d}\phi_i \cdot \prod_{i=1}^{q-3} \sin^{p-i-2} \psi_i \, \mathrm{d}\psi_i$$
$$\cdot \sinh^{p-2} \alpha \, \sinh^{q-2} \beta \, \mathrm{d}\phi_{p-2} \, \mathrm{d}\psi_{q-2} \, \mathrm{d}\alpha \, \mathrm{d}\beta.$$

Here  $\alpha, \beta \in [0; +\infty), \ \phi_1, \psi_1 \in [0; 2\pi) \ \text{and} \ \phi_2, \psi_2, \phi_3, \psi_3, \ldots \in [0; \pi).$ 

We now define three bases in the space  $D_{\sigma}$ . These bases consist of continuations of basis functions on  $\Gamma_K$ ,  $\Gamma_1$  and  $\Gamma_2$ .

We denote as  $\{f_V^{\sigma}\}$  the basis in  $D_{\sigma}$  related to reduction  $G \supset K$ , where

$$f_V^{\sigma}(x) = (x_1^2 + \ldots + x_p^2)^{\frac{\sigma - v_1 - v_{p+1}}{2}} \, \Xi_{V^*}^p(x^*) \, \Xi_{V^{**}}^q(x^{**})$$

and the function

$$\Xi_{(y_1,\dots,\pm y_{l-1})}^l(x) = \prod_{i=1}^{l-2} (x_1^2 + \dots + x_{l-1}^2)^{y_i - y_{i+1}} \cdot C_{y_i - y_{i+1}}^{\frac{l-i}{2} + y_{i+1}} \left( \frac{x_{l-i}}{x_1^2 + \dots + x_{l-1}^2} \right) (x_2 + \mathbf{i}x_1)^{y_l - 2}$$

was defined in [Vi, 9.3.6.2].

In addition, let us introduce two bases  $\{f_L^{\sigma}\}$  and  $\{f_M^{\sigma}\}$ , where

$$\begin{split} f_L^{\sigma}(x) &= (x_1^2 + \ldots + x_{p-1}^2)^{\frac{\sigma - l_1 + q/2 - 1}{2}} \, \Xi_{L^*}^{p-1}(x^*) \, \Xi_{L^{**},\lambda}^{q+1}(x^{**}), \\ f_M^{\sigma}(x) &= (x_1^2 + \ldots + x_{p-1}^2 - x_{p+1}^2)^{\frac{\sigma + p + q}{2} - 3} \, \Xi_{M^*,\mu^*}^p(x^*) \, \Xi_{M^{**},\mu^{**}}^q(x^{**}), \end{split}$$

where

$$\Xi_{Y,\upsilon}^m(\xi) = (\xi_1^2 + \ldots + \xi_{m-1}^2)^{\frac{3-p}{4} - \frac{y_1}{2}} P_{-\frac{1}{2} + \mathbf{i}\upsilon}^{\frac{3-p}{2} - y_1}(\xi_1) \, \Xi_Y^{m-1}(\xi).$$

It is possible to unify the results about matrix elements  $c_{VL}^{\sigma}$  and  $c_{LM}^{\sigma}$  of basis transformations. In order to do it, we introduce the function  $V(z_1, z_2, z_3, z_4, z_5, z_6)$ . Let us call

it the Vilenkin function.

$$\begin{split} V(z_1,z_2,z_3,z_4,z_5,z_6) &:= 2^{-1} \sqrt{\pi^{-(z_5+z_6)}} \, \mathbf{i}^{\frac{z_5}{2}+\sigma-2z_3-z_2-2} \, \Gamma\left(\frac{z_5+z_6-1}{2}\right) \\ & \cdot \Gamma^{-1}\left(\frac{z_5}{2}+z_2-1\right) \, \left| \Gamma\left(\frac{z_6-1}{2}+z_3+\mathbf{i}z_4\right) \right| \\ & \cdot \sqrt{\frac{z_1! \, \Gamma(z_5-2)(2z_1+z_5-2)z_4 \cosh(\pi z_4)}{(z_5-2)\Gamma(z_5+z_1-2)}} \\ & \cdot \sum_{s=0}^{\left[\frac{z_1-z_2}{2}\right]} \sum_{t=0}^{\infty} \frac{(-1)^s 2^{z_1-z_2-2s}}{s! \, t! \, (z_4-z_5-2s)!} \, \Gamma\left(\frac{z_5-2}{2}+z_1-s\right) \\ & \cdot \Gamma\left(\frac{\sigma-3z_3-z_1}{2}+s+t\right) \, \Gamma^{-1}(z_1-z_2-s) \, \Gamma^{-1}\left(\frac{\sigma-3z_3-z_1}{2}+s\right) \\ & \cdot \left[\sin\left[\left(-\frac{1}{2}+\mathbf{i}z_4\right)\pi\right] \, \Gamma\left(\frac{\sigma+z_1+z_6-z_3}{2}-z_2-s-t\right) \\ & \cdot \Gamma\left(\frac{z_3-\sigma-z_1-z_6+1}{2}+z_2+s+t-\mathbf{i}z_4\right) \\ & \cdot \Gamma\left(\frac{z_3-\sigma-z_1-z_6+1}{2}-z_2+s+t+\mathbf{i}z_4\right) \\ & \cdot 2^{\frac{\sigma+z_1-3z_3}{2}-z_2-s-t+1} \, \Gamma^{-1}\left(\frac{3z_3-\sigma-z_1}{2}+s+t+z_2\right) \\ & \cdot {}_3F_2\left(z_2-z_1+2s+t,\frac{z_3-\sigma-z_1-z_6+1}{2}+z_2+s+t-\mathbf{i}z_4,\frac{z_3-\sigma-z_1-z_6+1}{2}+z_2+s+t+\mathbf{i}z_4,\frac{z_3-\sigma-z_1-z_6+1}{2}-z_2+s+t+\mathbf{i}z_4;\frac{3z_3-\sigma-z_1}{2}+s+t+z_2,\frac{z_3-\sigma-z_1-z_6+1}{2}-z_2+s+t+\mathbf{i}z_4;\frac{1}{2}\right) \\ & + \Gamma\left(\frac{\sigma+z_6-z_1-z_3}{2}+s\right) \, \Gamma^{-1}\left(\frac{z_2+z_3-z_1-\sigma-z_6}{2}+s+t\right) \\ & \cdot \Gamma^{-1}(2s+t+z_2-z_1) \, \Gamma^{-1}\left(\frac{z_6}{2}+z_3\right) \, {}_3F_2\left(\frac{\sigma+z_6-z_1-z_3}{2}+s,\frac{1}{2}\right) \right]. \end{split}$$

Here  $z_1$  and  $z_2$  are both positive integers. Quadratic brackets mean the integer part of a number.

THEOREM 2.1.  $c_{VL}^{\sigma} = V(v_1, v_2, v_{p+1}, \lambda, p, q)$ .

Proof. Since

$$c_{VL}^{\sigma} = \frac{\Gamma\left(\frac{p+q-1}{2}\right)}{2\pi^{\frac{p+q}{2}-1}} \int_{\Gamma_1} f_V^{\sigma}(x) f_L^{-\sigma-p-q+2}(x) (dx)_{H_1}$$

and

$$\int_{\xi_1^2 + \dots + \xi_{p-2}^2 = 1} \Xi_{(l_1, \dots, l_{p-2})}^{p-1}(x) \, \Xi_{(v_2, \dots, v_{p-1})}^{p-1}(x) \, (\mathrm{d}x)_K \neq 0$$

only for  $l_i = v_{i+1}$ , then (see [Vi])

$$c_{VL}^{\sigma} \sim \int_{0}^{+\infty} (\mathbf{i} \sinh \alpha)^{\frac{q}{2} + \sigma - v_2 - 2v_{p+1}} C_{v_1 - v_2}^{\frac{p}{2} + v_2 - 1} (\coth \alpha) P_{-\frac{1}{2} + \mathbf{i}\lambda}^{1 - \frac{q}{2}} (\coth \alpha) d\alpha.$$

Now we can use the formulas

$$C_a^b(x) = \Gamma^{-1}(b) \sum_{j=0}^{\left[\frac{a}{2}\right]} \frac{(-1)^j \Gamma(a+b-j)}{j! (a-2j)! \Gamma(a-j)} (2x)^{a-2j}$$

and [E, 18.2.10].

Theorem 2.2.  $c_{LM}^{\sigma} = V(l_1, \mu^*, l_{p+1}, l_{p+2}, p-1, q+1).$ 

The formulas for  $c_{VL}^{\sigma}$  and  $c_{LM}^{\sigma}$  lead to some relations for special functions. For example, we derive the representation of the Gauss hypergeometric function.

THEOREM 2.3. If  $2 - p - q < \text{re } \sigma < 0 \text{ and } \alpha \neq 0 \text{ then}$ 

$${}_{2}F_{1}\left(-\sigma - \frac{1}{2}, \sigma + \frac{3}{2}; \frac{1}{2} + l; \frac{1 - \cosh \alpha}{2}\right)$$

$$= (-1)^{l-1} 2^{-\sigma - \frac{5}{2}} \pi^{-\frac{1}{2}} e^{-\alpha} \sinh \alpha \sin(-\pi \sigma)$$

$$\cdot \left(\frac{\cosh \alpha + 1}{\cosh \alpha - 1}\right)^{\frac{l}{2} + \frac{1}{4}} \Gamma(\sigma + 1 - l) \Gamma\left(l - \frac{3}{2}\right) \int_{0}^{\infty} \rho^{-\sigma - 1} K_{\sigma + 1}(\rho e^{-\alpha})$$

$$\cdot \sum_{s=0}^{\infty} (-1)^{n} \Gamma^{-2}(s + 1) \Gamma^{-1}(s - \sigma) G_{13}^{21} \left(\frac{\rho^{2}}{4} \middle| \frac{-s}{-\sigma - 1}, 0\right) d\rho.$$

*Proof.* Let q=1 and let  $\tilde{x}$  belong to the hyperboloid  $\vartheta(\tilde{x})=1$ . Let  $\Phi(x)=(x_1\tilde{x}_1+\ldots+x_p\tilde{x}_p-x_{p+1}\tilde{x}_{p+1}-\ldots x_{p+q}\tilde{x}_{p+q})^{\sigma}$ . Then

$$f_L^{-\sigma-p-q+2}(x) = \sum_V c_{LV}^{-\sigma-p-q+2} f_V^{-\sigma-p-q+2}(x).$$

The equality

$$\int_{\Gamma_1} f_L^{-\sigma - p - q + 2}(x) \,\Phi(x) \,(\mathrm{d}x)_{H_1} = \int_{\Gamma_K} \sum_V c_{LV}^{-\sigma - p - q + 2} f_V^{-\sigma - p - q + 2}(x) \,\Phi(x) \,(\mathrm{d}x)_K$$

leads to our formula.

**3.** A formula for Meijer and Legendre functions. Since the groups SO(p,q) and SO(q,p) are isomorphic, we can assume that  $p \geq q$ . Let us consider all p-1 partitions  $\Omega$  of the direct product  $\{1,\ldots,p\} \times \{p+1,\ldots,p+q\}$  into p classes

$$\Lambda_{\Omega,i} = \{(x_{i1}, p+1)\}, \dots, (x_{iq}, p+q) : \text{if } s \neq t \text{ then } x_{is} \neq x_{it}\}.$$

The set of matrices h(m,n), where  $(i,j) \in \Lambda_{\Omega,i}$ , generate a maximal  $\mathbb{R}$ -diagonalizable subalgebra  $A_{\Omega,i}$  in  $\mathfrak{h}$ . For different  $\Lambda_{\Omega,i}$  and  $\Lambda_{\hat{\Omega},\hat{i}}$ , such subalgebras are conjugate under Cartan involution. For any  $g \in G$ , we have  $g = g_1g_2g_3$ , where  $g_1, g_3 \in K$  and  $g_2 \in H_{\Omega,i} = \exp A_{\Omega,i}$ .

LEMMA 3.1. If  $\hat{\sigma} = -\sigma - p - q + 2$ , then the functional F is invariant under the pair  $(T_{\sigma}, T_{\hat{\sigma}})$ , i. e.  $\mathsf{F}(T_{\sigma}(g)(f_1), T_{-\sigma-p-q+2}(g)(f_2)) = \mathsf{F}(f_1, f_2)$ .

*Proof.* Without loss of generality, let us prove this lemma for the simplest case p=2, q=1. In this case, we have the only partition  $\Omega$  into two classes  $\Lambda_{\Omega,1}=\{(1,3)\}$  and  $\Lambda_{\Omega,2}=\{(2,3)\}$ . We will deal with the class  $\Lambda_{\Omega,1}$ . It is sufficient to argue separately for restrictions of  $T_{\sigma}$  to K and  $H_{\Omega,1}$ . According to (4), we can write an arbitrary point  $x\in X$  as  $x=(t\cos\phi,t\sin\phi,t)$ , where the point  $(\cos\phi,\sin\phi,1)$  belongs to  $\Gamma_K$ . If  $g\in K$ , then

$$g^{-1}(\alpha)x = \begin{pmatrix} \cos \alpha & \sin \alpha & 0 \\ -\sin \alpha & \cos \alpha & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} t\cos \phi \\ t\sin \phi \\ t \end{pmatrix} = \begin{pmatrix} t\cos(\phi - \alpha) \\ t\sin(\phi - \alpha) \\ t \end{pmatrix} \equiv \begin{pmatrix} t'\cos \phi' \\ t'\sin \phi' \\ t'. \end{pmatrix}$$
(6)

If  $g \in H_{\Omega,1}$  then

$$g^{-1}(s)x = \begin{pmatrix} \cosh s & 0 & -\sinh s \\ 0 & 1 & 0 \\ -\sinh s & 0 & \cosh s \end{pmatrix} \begin{pmatrix} t\cos\phi \\ t\sin\phi \\ t \end{pmatrix}$$
$$= \begin{pmatrix} t\cosh s\cos\phi - t\sinh s \\ t\sin\phi \\ -t\sinh s\cos\phi + t\cosh s \end{pmatrix} \equiv \begin{pmatrix} t'\cos\phi' \\ t'\sin\phi' \\ t'. \end{pmatrix} . \tag{7}$$

We obtain immediately from (7) that

$$\cos \phi' = \frac{t \cosh s \cos \phi - t \sinh s}{t'},\tag{8}$$

$$\sin \phi' = \frac{t \sin \phi}{t'},\tag{9}$$

$$t' = -t\sinh s\cos\phi + t\cosh s. \tag{10}$$

Let us find the partial derivative of  $\cos \phi'$  with respect to  $\phi$  from (8):

$$-\sin\phi'\,\mathrm{d}\phi' = -\frac{t^2\sin\phi\,\mathrm{d}\phi}{t'^2}.\tag{11}$$

Formulas (9), (10) and (11) lead to

$$\mathrm{d}\phi = \frac{t'\,\mathrm{d}\phi'}{t}.\tag{12}$$

Therefore,

$$\begin{split} \mathsf{F}(T_{\sigma}(g)(f_1), T_{-\sigma-1}(g)(f_2)) \\ &= \int_{\Gamma_K} f_1(t'\cos\phi', t'\sin\phi', t') \mid_{t=1} f_2(t'\cos\phi', t'\sin\phi', t') \mid_{t=1} \mathrm{d}\phi \\ &= \int_{\Gamma_K} t'^{\sigma} f_1(\cos\phi', \sin\phi', 1) \, t'^{-\sigma-1} f_2(\cos\phi', \sin\phi', 1) \, t' \, \mathrm{d}\phi' \\ &= \int_{\Gamma_K} f_1(\cos\phi', \sin\phi', 1) \, f_2(\cos\phi', \sin\phi', 1) \, \mathrm{d}\phi' = \mathsf{F}(f_1, f_2). \end{split}$$

Formula (6) leads to t' = t, dt' = dt,  $d\phi' = d\phi$ . Thus,

$$\begin{split} \mathsf{F}(T_{\sigma}(g)(f_1), T_{-\sigma-1}(g)(f_2)) \\ &= \int_{\Gamma_K} f_1(t'\cos\phi', t'\sin\phi', t') \mid_{t=1} f_2(t'\cos\phi', t'\sin\phi', t') \mid_{t=1} \mathrm{d}\phi \\ &= \int_{\Gamma_K} f_1(\cos\phi', \sin\phi', 1) \, f_2(\cos\phi', \sin\phi', 1) \, \mathrm{d}\phi' = \mathsf{F}(f_1, f_2). \ \blacksquare \end{split}$$

Theorem 3.2. For  $-1 < \text{re } \sigma < 0$ ,

$$\int_{0}^{+\infty} J_{0}(e^{t}\rho) G_{13}^{21} \left(\frac{\rho^{2}}{4} \middle| 0 \atop \sigma, 0, 0\right) G_{13}^{21} \left(\frac{e^{2t}\rho^{2}}{4} \middle| 0 \atop -\sigma - 1, 0, 0\right) d\rho$$

$$= -2^{\frac{5}{2}} \pi \sigma \sin^{-1}(-\pi \sigma) \sinh^{\frac{1}{2}}(-\beta) \Gamma(-\sigma - 1)$$

$$\cdot \sum_{s=0}^{\infty} (-1)^{s} \Gamma^{-1}(-\sigma - 3 - s) P_{-\sigma - \frac{3}{2}}^{-\frac{1}{2} - s}(\cosh \beta), \quad (13)$$

where  $\cosh \beta = \cosh t + \frac{d^2 e^t}{2}$ .

*Proof.* Let p = 1 and q = 3. Then

$$f_V^{\sigma}(x) = x_1^{\sigma - v_2} C_{v_1 - v_2}^{v_2 + \frac{1}{2}} \left(\frac{x_4}{x_1}\right) (x_3 + \mathbf{i}x_2)^{v_2},$$
  
$$f_L^{\sigma}(x) = (x_1 + x_4)^{\sigma} \exp \frac{\mathbf{i}(l_1 x_2 + l_2 x_3)}{x_1 + x_4}.$$

Consider the restriction  $T_{\sigma}^*$  of representation  $T_{\sigma}$  to the subgroup  $N \times \tilde{H}$ . Here the subgroup N consists of the matrices

$$\begin{pmatrix} 1 + \frac{d^2}{2} & d\cos\tau & d\sin\tau & \frac{d^2}{2} \\ d\cos\tau & 1 & 0 & d\cos\tau \\ d\sin\tau & 0 & 1 & d\sin\tau \\ -\frac{d^2}{2} & -d\cos\tau & -d\sin\tau & 1 - \frac{d^2}{2} \end{pmatrix}$$

and  $\tilde{H}$  consists of the matrices  $\hat{r}(1,4,t)$ . According to lemma 3.1, we obtain the matrix elements  $t_{L\hat{L}}^{*\sigma}(g)$  of  $T_{\sigma}^{*}$  in the following way:

$$t_{L\hat{L}}^{*\sigma}(n\hat{r}) = \mathsf{F}(T_{\sigma}(n\hat{r})(f_L^{\sigma}, f_{-\hat{L}}^{-\sigma-2}) = \mathsf{F}(T_{\sigma}(\hat{r})(f_L^{\sigma}), T_{-\sigma-2}(n^{-1})(f_{-\hat{L}}^{-\sigma-2})).$$

Let us integrate over  $\Gamma_N$ . Further, let us derive the matrix elements  $t^{\sigma}_{V\hat{V}}1(n\hat{r})$  and use relation between  $t^{\sigma}_{L\hat{L}}(n\hat{r})$  and  $t^{\sigma}_{V\hat{V}}(n\hat{r})$ . In this way, we obtain the relation containing the Bessel function, the Legendre function, the Gauss hypergeometric function, and two Meijer functions. The simplest form of this formula corresponds to  $V \equiv (v_1, v_2) = (0, 0)$  and coincides with formula (13).

4. Two formulas for Gauss, Whittaker, and Macdonald functions. Let us consider the case p = q = 2. We put here

$$f_V^{\sigma}(x) = (x_1^2 + x_2^2)^{\frac{\sigma - v_1 - v_2}{2}} (x_2 + \mathbf{i}x_1)^{v_1} (x_4 + \mathbf{i}x_3)^{v_2},$$
  
$$f_L^{\sigma}(x) = (x_2 + x_4)^{\sigma} \exp \frac{\mathbf{i}(l_1 x_1 - l_2 x_3)}{x_2 + x_4}.$$

LEMMA 4.1. If  $|l_1| < l_2$ , then

$$\begin{split} c_{VL}^{\sigma} &= 2^{-3} \, \pi^{-1} \, (l_2^2 - l_1^2)^{-\frac{\sigma}{2} - 1} \, \Gamma^{-1} \left( \frac{v_1 + v_2 - \sigma}{2} \right) \, \Gamma^{-1} \left( \frac{v_2 - v_1 - \sigma}{2} \right) \\ &\quad \cdot \, W_{\frac{v_2 + v_1}{2}, \, \frac{\sigma + 1}{2}} (l_2 - l_1) \, W_{\frac{v_2 - v_1}{2}, \, \frac{\sigma + 1}{2}} (l_2 + l_1). \end{split}$$

If  $|l_2| < l_1$ , then

$$\begin{split} c_{VL}^{\sigma} &= 2^{-3} \, \pi^{-1} \, (l_1^2 - l_2^2)^{-\frac{\sigma}{2} - 1} \, \Gamma^{-1} \left( -\frac{v_1 + v_2 + \sigma}{2} \right) \, \Gamma^{-1} \left( \frac{v_2 - v_1 - \sigma}{2} \right) \\ &\quad \cdot \, W_{-\frac{v_2 + v_1}{2}, \, \frac{\sigma + 1}{2}} (l_1 - l_2) \, W_{\frac{v_2 - v_1}{2}, \, \frac{\sigma + 1}{2}} (l_2 + l_1). \end{split}$$

If  $l_1 < 0$  and  $|l_2| < |l_1|$ , then

$$\begin{split} c_{VL}^{\sigma} &= 2^{-3} \, \pi^{-1} \, (l_1^2 - l_2^2)^{-\frac{\sigma}{2} - 1} \, \Gamma^{-1} \left( \frac{v_1 + v_2 - \sigma}{2} \right) \, \Gamma^{-1} \left( \frac{v_1 - v_2 - \sigma}{2} \right) \\ &\quad \cdot \, W_{\frac{v_2 + v_1}{2}, \, \frac{\sigma + 1}{2}} (l_2 - l_1) \, W_{\frac{v_1 - v_2}{2}, \, \frac{\sigma + 1}{2}} (|l_2 + l_1|). \end{split}$$

If  $l_2 < 0$  and  $|l_1| < |l_2|$  then

$$\begin{split} c_{VL}^{\sigma} &= 2^{-3} \, \pi^{-1} \, (l_2^2 - l_1^2)^{-\frac{\sigma}{2} - 1} \, \Gamma^{-1} \left( -\frac{v_1 + v_2 + \sigma}{2} \right) \, \Gamma^{-1} \left( \frac{v_1 - v_2 - \sigma}{2} \right) \\ &\quad \cdot \, W_{-\frac{v_2 + v_1}{2}, \, \frac{\sigma + 1}{2}} (l_1 - l_2) \, W_{\frac{k_1 - k_2}{2}, \, \frac{\sigma + 1}{2}} (|l_2 + l_1|). \end{split}$$

*Proof.* Let us choose the integration contour  $\Gamma = \Gamma_N$  in formula  $c_{VL}^{\sigma} = \mathsf{F}(f_V^{\sigma}, f_L^{-\sigma-2})$ , where the tangent space of the subgroup N is generated by the matrices

$$\begin{pmatrix} 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix} \qquad \text{and} \qquad \begin{pmatrix} 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 1 \\ 0 & 0 & -1 & 0 \end{pmatrix}.$$

Let  $(t, \frac{1-t^2+s^2}{2}, s, \frac{1+t^2-s^2}{2})$  be a parametrization of  $\Gamma_N$ . We compute this integral by formula [E, 3.2.12].

THEOREM 4.2. If  $\beta > 0$  and  $l_2 > |l_1|$ , then

$$\begin{split} {}_2F_1\left(1,\sigma+2;2;1+\frac{l_1}{l_2}\tanh\beta\right) &= 2^{-1}\,\pi\,\sigma\,(\sigma+1)^{-1} \\ &\quad \cdot \, \left(\mathrm{i}l_2\right)^{\sigma+2}\left(l_2-l_1\right)^{-\frac{\sigma+2}{2}}\,\sin(\sigma+1)\,\sinh^{-1}\beta\,\tanh^{-1}\beta\,e^{-\beta} \\ &\quad \cdot \, \sum_{v_1=0}^{\infty}\,\sum_{v_2=-v_1}^{v_1}(v_1^2+v_2^2)^{-\frac{\sigma+1}{2}}\,\Gamma^{-1}\left(\frac{v_1+v_2-\sigma}{2}\right)\,\Gamma^{-1}\left(\frac{v_2-v_1-\sigma}{2}\right) \\ &\quad \cdot \, W_{\frac{v_2+v_1}{2},\,\frac{\sigma+1}{2}}(l_2-l_1)\,W_{\frac{v_2-v_1}{2},\,\frac{\sigma+1}{2}}(l_2+l_1)\,K_{\sigma+1}\left((v_1^2+v_2^2)^{\frac{1}{2}}\,e^{-\beta}\right). \end{split}$$

Proof. Let  $u(x) := (x_1 \cosh \beta - x_3 \sinh \beta)^{\sigma}$ . Then  $\mathsf{F}(u, f_L^{-\sigma-2})_{\Gamma := \Gamma_N} = \mathsf{F}(u, f_L^{-\sigma-2})_{\Gamma := \Gamma_K}$ . In order to compute the corresponding integral along the contour  $\Gamma_K$ , we use

$$f_L^{-\sigma-2} = \sum_{v_1=0}^{\infty} \sum_{v_2=-v_1}^{v_1} c_{LV}^{-\sigma-2} f_V^{-\sigma-2} = \sum_{v_1=0}^{\infty} \sum_{v_2=-v_1}^{v_1} c_{VL}^{\sigma} f_V^{-\sigma-2}. \quad \blacksquare$$

In the same way, we derive

THEOREM 4.3. If  $\zeta > 1$  and  $l_2 > |l_1|$ , then

$$\int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} (\mathbf{i}l_2)^{\sigma} (l_2^2 - l_1^2)^{-\frac{\sigma}{2} - 1} W_{\frac{v_2 + v_1}{2}, \frac{\sigma + 1}{2}} (l_2 - l_1) 
\cdot W_{\frac{v_2 - v_1}{2}, \frac{\sigma + 1}{2}} (l_2 + l_1) {}_2F_1 \left( 1, -\sigma; 2; 1 + \frac{l_1}{l_2} \cdot \frac{\zeta^2 - 1}{\zeta^2 + 1} \right) dl_1 dl_2 
= 32 \left( v_1^2 + v_2^2 \right)^{\frac{\sigma + 1}{2}} \zeta^{-1} \left( \frac{\zeta^2 + 1}{2\zeta} \right)^{\sigma + 2} \left( \frac{\zeta^2 - 1}{\zeta^2 + 1} \right)^{\sigma + 1} \Gamma \left( \frac{v_1 + v_2 - \sigma}{2} \right) 
\cdot \Gamma \left( \frac{v_2 + v_1 - \sigma}{2} \right) \Gamma^{-1} (\sigma + 2) \Gamma^{-1} (-\sigma) K_{-\sigma - 1} \left( \frac{1}{2} \sqrt{v_1^2 + v_2^2} \right).$$

**Acknowledgments.** The second author would like to thank the Ministry of Science and Education of the Russian Federation for grant NK 586-P30, the Sholokhov Moscow State University for financial support, the organizers of the conference for a pleasant atmosphere.

## References

- [Vi] N. Ja. Vilenkin, Special Functions and Theory of Group Representations, AMS, Providence, R.I., 1968.
- [Ve] Y. A. Verdiev, Invariants of the representations of the group of pseudo-orthogonal matrices, Preprint 46 of the Inst. for Theor. Phys. of the Nat. Acad. of Sci. of Ukraine, 1977 (in Russian).
- [M] V. F. Molchanov, Representations of pseudo-orthogonal groups associated with a cone, Mat. Sb. 81 (1970), 358–375 (in Russian); English transl.: Math. USSR Sb. 10 (1970), 333–347.
- [E] A. Erdelyi, Tables of Integral Transforms, McGraw-Hill, New York, 1954.

Received January 30, 2010; Revised March 5, 2011