Asymptotics for the Second Painlevé equation

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This talk discusses joint work with Alexander Its on the large x asymptotical behavior of solutions u of the second Painlevé equation (PII)

$$u_{xx} = xu + 2u^3, \tag{1}$$

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Recall the following Riemann-Hilbert representation of PII: Given

$$\{s_k\}_{k=1}^6 \subset \mathbb{C}: \ s_1 - s_2 + s_3 + s_1 s_2 s_3 = 0, \ s_{k+3} = -s_k, \ s_1 = \bar{s}_3, s_2 = \bar{s}_2,$$

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put $S_k = \begin{pmatrix} 1 & 0 \\ s_k & 1 \end{pmatrix}$ if k odd, $S_k = \begin{pmatrix} 1 & s_k \\ 0 & 1 \end{pmatrix}$ if k even and consider the following Riemann-Hilbert problem (RHP) which consists in finding the piecewise analytic 2 × 2 function $\Psi(\lambda)$:



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 $\Psi(\lambda)$ is analytic on $\mathbb{C} ig \cup_1^6 \Gamma_k$



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$$\Psi_+(\lambda) = \Psi_-(\lambda)S_k, \ \lambda \in \Gamma_k$$



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$$\begin{split} \Psi_{+}(\lambda) &= \Psi_{-}(\lambda)S_{k}, \ \lambda \in \Gamma_{k} \\ \Psi(\lambda) &= (I+O(\lambda^{-1}))e^{-i(\frac{4}{3}\lambda^{3}+x\lambda)\sigma_{3}}, \ \lambda \to \infty \end{split}$$



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This RHP is for any $s \equiv (s_1, s_2, s_3)$ meromorphically w.r.t. x solvable [Bolibruch, Its, Kapaev, 2004]



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This RHP is for any $s \equiv (s_1, s_2, s_3)$ meromorphically w.r.t. x solvable [Bolibruch, Its, Kapaev, 2004] and its solution determines the solution of (1) via

$$u(x) \equiv u(x;s) = 2 \lim_{\lambda \to \infty} \left(\lambda (\Psi(\lambda) e^{i(\frac{4}{3}\lambda^3 + x\lambda)\sigma_3})_{12} \right), \quad \bar{u}(x) = u(\bar{x}).$$

Theorem (Kapaev, 1992; Bolts, 2010)

As $x \to -\infty$ and $|s_1| > 1$, the following asymptotical behavior holds for real-valued solutions of PII equation (1)

$$u(x) = \frac{\sqrt{-x}}{\sin\left(\frac{2}{3}(-x)^{3/2} + \beta \ln\left(8(-x)^{3/2}\right) + \varphi\right) + O((-x)^{-3/2})} + O((-x)^{-1})$$

where

$$eta = rac{1}{2\pi} \ln(|s_1|^2 - 1), \ \ arphi = -rg \Gammaig(rac{1}{2} + ietaig) - rg s_1.$$

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However as $x \to +\infty$ and $s_2 \neq 0$, we have

$$u(x) = \sigma \sqrt{\frac{x}{2}} \cot\left(\frac{\sqrt{2}}{3}x^{3/2} + \frac{\gamma}{2}\ln\left(8\sqrt{2}x^{3/2}\right) + \phi\right) + O(x^{-1}),$$

where

$$\sigma = \text{sgn } s_2, \ \gamma = \frac{1}{\pi} \ln |s_2|, \ \phi = -\frac{1}{2} \text{arg } \Gamma(\frac{1}{2} + i\gamma) - \frac{1}{2} \text{arg } (1 + s_2 s_3) + \frac{\pi}{2}.$$

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Sketch of the proof as $x \to -\infty$

The solution of the given RHP

$$\Psi_{+}(\lambda) = \Psi_{-}(\lambda) \underbrace{e^{-\theta(\lambda)\sigma_{3}}S_{k}e^{\theta(\lambda)\sigma_{3}}}_{\equiv G(\lambda)}, \ \lambda \in \Gamma_{k}; \quad \theta(\lambda) = i\left(\frac{4}{3}\lambda^{3} + x\lambda\right)$$
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$$\Psi(\lambda) = I + \frac{1}{2\pi i} \int_{\cup_1^6 \Gamma_k} \Psi_-(w) \big(G(w) - I \big) \frac{dw}{w - \lambda}, \quad \lambda \notin \cup_1^6 \Gamma_k,$$

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so as $|x| \to \infty$ a "steepest descent evaluation" for the oscillatory integral is needed.

We start with a scaling

$$\lambda \to z(\lambda) = \frac{\lambda}{\sqrt{-x}}, \qquad \Psi(z) \equiv \Psi(\lambda(z)),$$

 $\theta(\lambda) \to t \vartheta(z), \qquad \vartheta(z) = i \left(\frac{4}{3}z^3 - z\right), \ t = (-x)^{3/2}.$

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Then use analytic structure of $G(\lambda)$ and deform jump contour $\bigcup_{1}^{6}\Gamma_{k}$ to steepest descent contours of the exponent $\vartheta(z)$:



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a Gauss LU-decomposition (since $1-\mathit{s}_1\mathit{s}_3=1-|\mathit{s}_1|^2\neq 0)$

$$(S_3 S_4 S_5)^{-1} = \begin{pmatrix} 1 - s_1 s_3 & s_1 \\ s_1 & 1 - s_1 s_3 \end{pmatrix} = S_L S_D S_U$$

= $\begin{pmatrix} 1 & 0 \\ \frac{s_1}{1 - s_1 s_3} & 1 \end{pmatrix} (1 - s_1 s_3)^{\sigma_3} \begin{pmatrix} 1 & \frac{s_1}{1 - s_1 s_3} \\ 0 & 1 \end{pmatrix}$

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Now observe that triangularity of G(z) and the sign of $\Re \vartheta(z)$ imply for $|z-z_\pm|>r, 0< r<\frac{1}{2}$

 $\| {\boldsymbol{G}} - {\boldsymbol{I}} \|_{L^2 \cap L^\infty(\text{infinite branches})} = O\big(e^{-ct|z-z_\pm|} \big), c > 0 \ t \to \infty,$

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 $\|G-I\|_{L^2\cap L^\infty(\text{infinite branches})}=O\big(e^{-ct|z-z_\pm|}\big), c>0 \quad t\to\infty,$

i.e. main contribution in (2) arises from $[z_-, z_+]$ and neighborhood of z_{\pm} . Construct model solutions/parametrices:



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Along $[z_-, z_+]$ we solve the diagonal RHP applying the Plemelj -Sokhotskii formulae: For $z \in \mathbb{C} \setminus [z_-, z_+]$

$$\Psi^{D}(z) = \left(\frac{z-z_{-}}{z-z_{+}}\right)^{\nu\sigma_{3}} e^{-t\vartheta(z)\sigma_{3}}, \ \nu = -\frac{1}{2\pi i} \ln(1-s_{1}s_{3})$$

where arg $(1 - s_1 s_3) \in (-\pi, \pi]$ and $(\frac{z-z_-}{z-z_+})^{\nu} \to 1$ as $z \to \infty$.

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In the neighborhood of z_+ we notice the quadratic local behavior of $\vartheta(z)$:

$$\vartheta(z) = \vartheta(z_+) + 2i(z-z_+)^2 + O((z-z_+)^3), \quad z \to z_+$$

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$$\vartheta(z) = \vartheta(z_+) + 2i(z-z_+)^2 + O((z-z_+)^3), \quad z \to z_+$$

and make use of the parabolic cylinder function $D_{\nu}(\zeta)$, a solution of

$$\frac{d^2 D_{\nu}}{d\zeta^2} + \left(\nu + \frac{1}{2} - \frac{\zeta^2}{4}\right) D_{\nu} = 0, \quad D_{\nu}(\zeta) \sim \zeta^{\nu} e^{-\frac{\zeta^2}{4}}, \ \zeta \to +\infty.$$
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Skipping details, we can introduce a parametrix $\Psi^{r}(z)$ near $z = z_{+}$, such that locally its jumps coincide with those of $\Psi(z)$

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Skipping details, we can introduce a parametrix $\Psi^{r}(z)$ near $z = z_{+}$, such that locally its jumps coincide with those of $\Psi(z)$ and as $t \to \infty$ the following match-up relation between $\Psi^{r}(z)$ and $\Psi^{D}(z)$ holds

$$\Psi'(z) = \begin{pmatrix} 1 & -\frac{\nu s_3}{h_1} e^{\frac{2it}{3}} \beta(z) \frac{1}{\zeta(z)} \\ -\frac{h_1}{s_3} e^{-\frac{2it}{3}} \beta^{-1}(z) \frac{1}{\zeta(z)} & 1 \end{pmatrix} (I + O(t^{-1})) \Psi^D(z),$$

as long as $0 < r_1 \leq |z - z_+| \leq r_2 < 1$, with $h_1 = rac{\sqrt{2\pi}}{\Gamma(-\nu)} e^{i\pi\nu}$,

$$\zeta(z) = 2\sqrt{-t\vartheta(z) + t\vartheta(z_+)} = e^{i\frac{3\pi}{4}}2\sqrt{2t}(z-z_+)(1+O(z-z_+))$$

and

$$\beta(z) = \left(\zeta(z)\frac{z-z_-}{z-z_+}\right)^{2\nu}, \text{ i.e. } \beta^{\pm 1}\frac{1}{\zeta} = O\left(t^{-\frac{1}{2}\pm\Re\nu}\right), \ t \to \infty.$$

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We make the following crucial observation (recall $u = -rac{1}{2\pi i}\ln(1-|s_1|^2))$

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$$\beta^{\pm 1} \frac{1}{\zeta} = o(1), t \to \infty \text{ if } |s_1| < 1, \text{ i.e. } \Psi^r(z) = (I + O(t^{-1})) \Psi^D(z)$$

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however, here $|s_1| > 1$, so $\nu = -\frac{1}{2\pi i} \ln(|s_1|^2 - 1) - \frac{1}{2} \equiv \nu_0 - \frac{1}{2}$, i.e.

$$\beta \zeta = \frac{\beta_0(z)}{\alpha(z)} = O(1), t \to \infty \text{ i.e. } \Psi'(z) = \left(E_r(z) + O(t^{-1})\right) \Psi^D(z)$$

with

$$E_r(z) = \begin{pmatrix} 1 & 0 \\ -\frac{h_1}{s_3} e^{-\frac{2it}{3}} \frac{\alpha(z)}{\beta_0(z)} & 1 \end{pmatrix}, \quad \beta_0(z) = \left(\zeta(z) \frac{z - z_-}{z - z_+}\right)^{2\nu_0}$$

and $\alpha(z) = \frac{z-z_-}{z-z_+}$.

Now assemble the model functions

$$\hat{\Psi}(z) = \left\{ egin{array}{ll} \Psi^D(z), & |z-z_{\pm}| > r; \ \Psi^r(z), & |z-z_{+}| < r; \ \Psi^l(z) = \sigma_2 \Psi^r(-z) \sigma_2, & |z-z_{-}| < r, \end{array}
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consider the ratio $\Psi(z) = \chi(z)\hat{\Psi}(z)$ and obtain the following ratio-RHP:

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$$\begin{array}{rcl} \chi_+(z) & = & \chi_-(z)\Psi^{r,\prime}(z) \big(\Psi^D(z)\big)^{-1}, z \in C_{r,\prime} \\ \chi(z) & \to & I, \ z \to \infty \end{array}$$

but $\|\Psi^{r,l}(\Psi^D)^{-1} - I\| \not\rightarrow 0$ as $t \rightarrow \infty$.

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$$\hat{\Psi}(z) = \left\{ egin{array}{ll} \Psi^D(z), & |z-z_{\pm}| > r; \ \Psi^r(z), & |z-z_{+}| < r; \ \Psi^l(z) = \sigma_2 \Psi^r(-z) \sigma_2, & |z-z_{-}| < r, \end{array}
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$$egin{array}{rl} \chi_+(z) &=& \chi_-(z)\Psi^{r,l}(z)ig(\Psi^D(z)ig)^{-1}, z\in C_{r,l}\ \chi(z) &\to& l, \ z o\infty \end{array}$$

but $\|\Psi^{r,l}(\Psi^D)^{-1} - I\| \neq 0$ as $t \to \infty$. To this end employ an undressing transformation and pass from $\chi(z)$ to $\Phi(z)$:

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To deal with the singularities, use a final dressing transformation:

$$\Phi(z) = (zI+B)Y(z)\begin{pmatrix}\frac{1}{z-z_+} & 0\\ 0 & \frac{1}{z-z_-}\end{pmatrix}$$

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and determine B uniquely from the residue-relations (4).

Finally solve the Y-RHP asymptotically: Since

$$\|G_Y - I\|_{L^2 \cap L^\infty(ext{jump contour})} \leq rac{c}{t}, \ c > 0$$

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$$Y_{-}(z) = I + \frac{1}{2\pi i} \int_{\text{contour}} Y_{-}(w) \big(G_{Y}(w) - I \big) \frac{dw}{w - z_{-}},$$

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can be solved iteratively in $L^2(\text{jump contour})$ and its unique solution satisfies

$$\|Y_{-} - I\|_{L^2(\text{jump contour})} \leq \frac{c}{t}, \quad t \to \infty.$$

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In the end use the latter estimate together with the integral representation for Y(z)

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and extract the required asymptotics via

$$u(x) = 2\sqrt{-x} \lim_{z \to \infty} \left(z(\Psi(z)e^{t\vartheta(z)\sigma_3})_{12} \right)$$

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by tracing back: $\Psi(z) \to \chi(z) \to \Phi(z) \to Y(z).$

Reference

Thomas Bothner and Alexander Its, The Nonlinear Steepest Descent Approach to the Singular Asymptotics of the Second Painlevé Transcendent.

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