## Self-similar solutions of semilinear wave equations

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# Motivation

## Based mainly on [Kycia].

#### Semilinear wave equation

$$\Box U(t,x) - U(t,x)^p = 0, \qquad \Box = \partial_{tt} - \Delta,$$

where  $x \in R^n, \, n \geq 3, \, p\text{-even to preserve reflection symmetry or } U^p \to |U|^{p-1}U.$ 

#### Spherical symmetry

$$U_{tt} - U_{rr} - \frac{n-1}{r}U_r - U^p = 0,$$

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There exist 'smooth' initial data that develop singularity when  $t \to T < \infty.$ 



- Example of nonglobal existence.
- Common behavior for many nonlinear PDEs, see [Eggers, Fontelos].
- Blowup dynamics is governed by self-similar solutions.

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## Self-similar solutions

$$U(t,r) = \frac{u(\rho)}{(T-t)^{\alpha}}, \qquad \rho = \frac{r}{T-t}, \qquad \alpha = \frac{2}{p-1} (>0).$$

#### Self-similar solution

**1** 
$$b_0 = \left(\frac{2(p+1)}{(p-1)^2}\right)^{\frac{1}{p-1}}$$
  
**2**  $b_\infty = \left(\frac{2(p(n-2)-n)}{(p-1)^2}\right)^{\frac{1}{p-1}}$ 

# Self-similar profile $(0, m) = h_0$ $(0, m_0, (n) = h_0 n^{-\alpha}$

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#### ODE for self-similar profiles

$$(1-\rho^2)u'' + \left(\frac{n-1}{\rho} - \frac{2(p+1)}{p-1}\rho\right)u' - \frac{2(p+1)}{(p-1)^2}u + u^p = 0,$$
 where  $' = \frac{d}{d\rho}.$ 

## • Fixed singularities at 0, $\pm 1$ , $\infty$ .

• Question: Is there a global (on [0;1]) analytic solution ?

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  - $C^1$  match at some  $\rho_0 \in (0;1)$ .

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• On substituting the formal power series  $u(\rho) = \sum_{l=0}^{\infty} a_l \rho^l$ • and using Cauchy product

$$\left(\sum_{l=0}^{\infty} a_l (x - x_0)^l\right)^p = \sum_{l=0}^{\infty} c_l (x - x_0)^l,$$
  
$$c_0 = a_0^p, \qquad c_m = \frac{1}{ma_0} \sum_{l=1}^m (lp - m + l) a_l c_{m-l},$$

that simplifies nonlinear term  $u^p$ 

- we obtain unique recurrence for  $\{a_l\}_{l=1}^{\infty}$
- and the solution

$$u(\rho) = c + \frac{1}{n} \left[ c \frac{p+1}{(p-1)^2} - \frac{1}{2} c^p \right] \rho^2 + O(\rho^4),$$

where c is initial data at  $\rho = 0$ .

#### Question: Is the formal solution a solution ?

#### Proposition [Breitenlohner, Forgács, Maison]

Consider a system of differential equations for i + j functions  $u = (u_1, \ldots, u_i)$  and  $v = (v_1, \ldots, v_j)$ ,

$$t\frac{du_l}{dt} = t^{\mu_l}f_l(t, u, v), \qquad t\frac{dv_l}{dt} = -\lambda_l v_l + t^{\nu_l}g_l(t, u, v),$$

with constants  $\lambda_l > 0$  and integers  $\mu_l, \nu_l \ge 1$  and let U be an open subset of  $\mathbb{R}^n$  such that the functions f and g are analytic in a neighborhood of t = 0, u = c, v = 0 for all  $c \in U$ . Then there exists an *i*-parameter family of solutions of that system such that

$$u_l(t) = c_l + O(t^{\mu_l}), \qquad v_l(t) = O(t^{\nu_l}),$$

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#### The Lane-Emden equation

$$U'' + \frac{n-1}{x}U' + U^p = 0, \qquad U = U(x)$$

- Widely used in astrophysics to model star structures.
- The simplest equation with power type nonlinearity and fixed singularities at 0 and ∞.
- Formal power series solution at x = 0 is of the form  $U(x) = 1 \frac{1}{2n}x^2 + O(x^4)$  with U(0) = 1, U'(0) = 0.
- Formal solutions are solutions apply the Proposition [Breitenlohner, Forgács, Maison]:

$$\begin{cases} U' = V \\ xV' = -(n-1)V - xU^p \end{cases}$$

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Movable singularities decrease the radius of convergence. They can be determined numerically - scanning complex plane [Hunter].



Figure: The Lane-Emden solution for n = 3, p = 5.

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Figure: The Lane-Emden solution for n = 3, p = 5.

p	movable singularities
1	$\infty$
2	$\pm 3.964581i$
3	$\pm 2.574840i$
4	$\pm 2.034896i$
5	$\pm 1.732051i$

Table: Numerical estimates of the location of movable singularities for n = 3 [Hunter].

Question: How large is the circle of convergence - where are the closest movable singularities ?

• Rescaling local solution at  $\rho = 0$  using  $\rho = \frac{x}{c^{\frac{p-1}{2}}}$  and U = u/c• we get

$$u(\rho) = cU\left(\frac{x}{c^{\frac{p-1}{2}}}\right) = c\left[\left(1 - \frac{1}{2n}x^2 + O(x^4)\right) + o\left(\frac{1}{c}\right)\right].$$

• The equation transforms as

$$U'' + \frac{n-1}{x}U' + U^p = \frac{1}{c^{p-1}} \left( x^2 U'' + (2\alpha + 2)xU' + \alpha(\alpha + 1)U \right).$$

- For large c asymptotic form of local solution at  $\rho = 0$  obey the Lane-Emden equation.
- Connection between location of movable singularities (the radius of convergence):

$$R_W \approx \frac{R_{LE}}{c^{\frac{p-1}{2}}}.$$

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#### ODE for self-similar profile at $\rho = 1$

$$y(2-y)u''(y) - \left[\frac{n-1}{1-y} - \frac{2(p+1)}{p-1}(1-y)\right]u'(y) - \frac{2(p+1)}{(p-1)^2}u(y) + u^p(y) = 0,$$

where  $y = 1 - \rho$  and  $' = \frac{d}{dy}$ .

There exist the parameter  $k = \frac{(n-1)p-n-3}{2(p-1)}$  which determines the form of analytic solution.



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Formal power series solutions  $u(y) = \sum_{l=0}^{\infty} a_l y^l$ :

 $\bullet$  For noninteger k we get standard formal analytic solution

$$u(y) = b + \frac{2(p+1)b - (p-1)^2 b^p}{2(1-k)(p-1)^2}y + O(y^2),$$

where b = u(y = 0) is a free parameter.

- For integer k we have:
  - First k recurrence equations for {a<sub>0</sub>,..., a<sub>k-1</sub>} give nonlinear algebraic system of equations...
  - ...which has (among others) the solution

$$a_0 = b_{\infty}, \dots, a_{k-1} = (-1)^k \frac{u_{\infty}^{(\kappa)}(\rho)|_{\rho=1}}{k!}.$$

 The free parameter b appears at y<sup>k</sup> and then formal series proceed in usual way

 $u(y) = a_0 + \ldots + a_{k-1}y^{k-1} + by^k + a_{k+1}(b, a_0, \ldots, a_{k-1})y^{k+1} + \ldots$ 

The proof of convergence uses Proposition
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Movable singularity moves from  $\rho = \infty$  to  $\rho = 1$  along the real axis when  $|b| > b_0$  increases.



Figure: Local solution at  $\rho = 1$  for n = 3, p = 7 and  $b = 1.1 > b_0 \approx 0.8736$ .

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- Use of Lyapunov functions local solutions can develop singularities only at endpoints of [0; 1].
- $C^1$  matching at  $\rho_0$  (analytic continuation):
  - C<sub>0</sub> curve of integration of local analytic solution at 0 to ρ<sub>0</sub> and parameterized by initial data c.
  - C<sub>1</sub> curve of integration of local analytic solution at 1 to ρ<sub>0</sub> and parametrized by initial data b.

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# Global existence



Figure: n = 3, p = 7,  $\rho_0 = 0.5$ 

- The curves have a countable number of intersections ⇒ countable family of global analytic solutions.
- The family can be enumerated by the pairs of initial data at both endpoints {c<sub>l</sub>, b<sub>l</sub>}<sup>∞</sup><sub>l=0</sub>.
- For integer k only solutions with asymptotics at  $\rho = 1$  for which  $u(1) = b_{\infty} (a_0 = b_{\infty}, \dots, a_{k-1} = (-1)^k \frac{u_{\infty}^{(k)}(\rho)|_{\rho=1}}{k!})$ gives countable family of global solutions; the others give only one global solution for each solution. These constraints lead to simple numerical methods to study global solutions [Kycia].

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# Thank You for Your Attention Any suggestions and comments would be extremely useful.

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# Classification

## Equation's symmetry

$$U_{\lambda}(t,r) = \lambda^{\alpha} U\left(\frac{t}{\lambda}, \frac{r}{\lambda}\right), \qquad \alpha = \frac{2}{p-1}.$$

• Scaling is critical for  $\beta = 0$ , i.e.,  $p = p_Q = \frac{n+2}{n-2}$ ,

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- subcritical for  $p < p_Q$ ,
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Energy functional

$$E[U] = \int_{\mathbf{R}^n} \left( U_t^2 + (\nabla U)^2 - \frac{1}{p+1} U^{p+1} \right) d^n x,$$

scales as

$$E[U_{\lambda}] = \lambda^{\beta} E[U],$$

where  $\beta = \frac{(n-2)p-(n+2)}{p-1}$ .

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# Classification

#### Equation's symmetry

$$U_{\lambda}(t,r) = \lambda^{\alpha} U\left(\frac{t}{\lambda}, \frac{r}{\lambda}\right), \qquad \alpha = \frac{2}{p-1}.$$

Energy functional

$$E[U] = \int_{\mathbf{R}^n} \left( U_t^2 + (\nabla U)^2 - \frac{1}{p+1} U^{p+1} \right) d^n x,$$

scales as

$$E[U_{\lambda}] = \lambda^{\beta} E[U],$$

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where  $\beta = \frac{(n-2)p-(n+2)}{p-1}$ .

- Scaling is critical for  $\beta = 0$ , i.e.,  $p = p_Q = \frac{n+2}{n-2}$ ,
- subcritical for  $p < p_Q$ ,
- supercritical for  $p > p_Q$ .