

Deformation quantization of Hamiltonian actions in Poisson geometry

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Outline

Hamiltonian actions

- Hamiltonian actions in canonical setting
- Hamiltonian actions in Poisson geometry
- Poisson Reduction

Quantization

- Approach
- Quantum Hamiltonian actions
- Quantum reduction
- Examples

Towards?

- Symplectic groupoids and Hamiltonian actions
- Quantization

Canonical action

Definition

Let G be a Lie group acting on a Poisson manifold (M, π) . The action $\Phi : G \times M \rightarrow M$ is said canonical if

$$\Phi_g^* \{f, h\} = \{\Phi_g^* f, \Phi_g^* h\} \quad \forall f, h \in C^\infty(M)$$

What is Hamiltonian action?

Momentum map

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the map $\mu : M \rightarrow \mathfrak{g}^*$ defined by

$$H_\xi(m) = \langle \mu(m), \xi \rangle$$

is the momentum map for the canonical action.

Hamiltonian action

A momentum map $\mu : M \rightarrow \mathfrak{g}^*$ is equivariant if the correspondent $H : \mathfrak{g} \rightarrow C^\infty(M)$ is a Lie algebra homomorphism.

Theorem

A canonical action is Hamiltonian if and only if there is a Lie algebra homomorphism $\psi : \mathfrak{g} \rightarrow C^\infty(M)$ such that $X_{\psi(\xi)} = \xi_M$ for all $\xi \in \mathfrak{g}$. If ψ exists, an equivariant momentum map μ is determined by $H = \psi$. Conversely, if μ is equivariant, we can take $\psi = H$.

Symplectic Reduction

Marsden and Weinstein idea: we can reduce the size of the phase space by taking advantage of the momentum map and the invariance of the system under the given symmetry group.

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Theorem (Marsden-Weinstein Reduction)

Let $\Phi : G \times M \rightarrow M$ be a Hamiltonian action of the Lie group G on the Poisson manifold (M, π) with momentum map $\mu : M \rightarrow \mathfrak{g}^$. Let $u \in \mathfrak{g}^*$ be a regular value of μ and suppose that G_u acts freely and properly on the manifold $\mu^{-1}(u)$. Then there is a Poisson structure π_u on the reduced space $M//G := \mu^{-1}(u)/G_u$.*

Ingredients

What is a Hamiltonian action in this context?

Ingredients:

- Poisson Lie groups
- Lie bialgebras

Poisson Lie groups

Definition

A Poisson Lie group (G, π) is a Lie group equipped with a Poisson structure π which preserves multiplication and inverse.

Example

$\pi = 0$ is obviously multiplicative, hence any Lie group G with the trivial Poisson structure is a Poisson Lie group.

Lie bialgebra

Definition

A Lie bialgebra is a Lie algebra \mathfrak{g} with a linear map $\delta : \mathfrak{g} \rightarrow \mathfrak{g} \wedge \mathfrak{g}$ such that

- ${}^t\delta : \mathfrak{g}^* \otimes \mathfrak{g}^* \rightarrow \mathfrak{g}^*$ defines a Lie bracket on \mathfrak{g}^* , and
- δ is a 1-cocycle on \mathfrak{g} relative to the adjoint representation of \mathfrak{g} on $\mathfrak{g} \otimes \mathfrak{g}$

$$ad_{\xi}(\delta(\eta)) - ad_{\eta}(\delta(\xi)) - \delta([\xi, \eta]) = 0$$

Dual Poisson Lie group

Theorem

If (G, π_G) is a Poisson Lie group, then the linearization of π_G at e defines a Lie algebra structure on \mathfrak{g}^ . Conversely, if G is connected and simply connected, then every Lie bialgebra (\mathfrak{g}, δ) defines a unique multiplicative Poisson structure π_G on G .*

This implies that G^* is also a Poisson Lie group, called dual.

Poisson action

Definition

The action of (G, π_G) on (M, π) is called Poisson action if the map $\Phi : G \times M \rightarrow M$ is Poisson, where $G \times M$ is a Poisson manifold with structure $\pi_G \oplus \pi$.

Generalization of canonical action! If $\pi_G = 0$, the action is Poisson if and only if it preserves π .

Momentum map

Definition (Lu)

A momentum map for the Poisson action $\Phi : G \times M \rightarrow M$ is a map $\mu : M \rightarrow G^*$ such that

$$\xi_M = \pi^\sharp(\mu^*(\theta_\xi))$$

where θ_ξ is the left invariant 1-form on G^* defined by the element $\xi \in \mathfrak{g} = (T_e G^*)^*$ and μ^* is the cotangent lift $T^*G^* \rightarrow T^*M$.

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A Hamiltonian action is a Poisson action induced by an equivariant momentum map.

Infinitesimal momentum map

Let's focus on the map $\alpha : \mathfrak{g} \rightarrow \Omega^1(M)$

Definition

Let M be a Poisson manifold and G a Poisson Lie group. An infinitesimal momentum map is a morphism of Gerstenhaber algebras

$$\alpha : (\wedge^\bullet \mathfrak{g}, \delta, [,]) \longrightarrow (\Omega^\bullet(M), d_{DR}, [,]_\pi).$$

Poisson Reduction

Theorem

Let $\Phi : G \times M \rightarrow M$ be a Hamiltonian action with momentum map $\mu : M \rightarrow G^*$ and $u \in G^*$ a regular value of μ . The Poisson reduction of (M, G) is the quotient

$$M//G := \mu^{-1}(\mathcal{O}_u)/G$$

$M//G$ inherits a Poisson structure from M .

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Example

Suppose that $\pi_G = 0$. Then

$$C^\infty(M//G) \simeq (C^\infty(M)/\mathcal{I})^G$$

Deformation quantization approach

Goal: quantize Hamiltonian actions and Poisson reduction.

Steps:

- 1 Quantize a Poisson action
- 2 Quantize Momentum map
- 3 Quantize Poisson reduction

Quantization of Poisson manifold

Given (M, π) we define

$$f \star g = f \cdot g + \sum_{n=1}^{\infty} \hbar^n P_n(f, g)$$

where

$$P_1(f, g) - P_1(g, f) = \{f, g\}$$

Quantization of Lie bialgebra

Given (\mathfrak{g}, δ) we associate the Hopf algebra $(\mathcal{U}(\mathfrak{g}), \Delta)$, where

$$\Delta X = X \otimes 1 + 1 \otimes X$$

Quantum group $(\mathcal{U}_{\hbar}(\mathfrak{g}), \Delta_{\hbar}, [\cdot, \cdot]_{\hbar})$

$$\Delta_{\hbar} = \Delta + \sum_{n=1}^{\infty} \hbar^n \Delta_n$$

Quantum action

How can we define a quantum action of $\mathcal{U}_{\hbar}(\mathfrak{g})$ on \mathcal{A}_{\hbar} ?

- Hopf algebra action
- $\hbar \rightarrow 0$ Poisson action

Quantum action

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- $\hbar \rightarrow 0$ Poisson action

Definition

The quantum action is a linear map

$$\Phi_{\hbar} : \mathcal{U}_{\hbar}(\mathfrak{g}) \rightarrow \text{End } \mathcal{A}_{\hbar} : \xi \mapsto \Phi_{\hbar}(\xi)(f)$$

such that

- 1 Hopf algebra action
- 2 Algebra homomorphism

Quantum Hamiltonian action

- 1 Quantum momentum map which, as in the classical case, factorizes the quantum action
- 2 $\hbar \rightarrow 0$ classical momentum map

Non commutative forms

The non-commutative analogue of the de Rham complex is $(\Omega(\mathcal{A}_\hbar), d)$ with the universal derivation

$$d : \mathcal{A}_\hbar \rightarrow \Omega(\mathcal{A}_\hbar)$$

Quantum momentum map

The map

$$adb \longmapsto a[b, \cdot]_*$$

induces a non commutative product on $\Omega(\mathcal{A}_{\hbar})$ and natural morphism of differential graded algebras

$$\Omega^1(\mathcal{A}_{\hbar}) \longrightarrow C^1(\mathcal{A}_{\hbar}, \mathcal{A}_{\hbar})$$

This induces a first definition of the momentum map.

Quantum momentum map

Definition

A quantum momentum map is defined to be a linear map

$$\mu_{\hbar} : \mathcal{U}_{\hbar}(\mathfrak{g}) \rightarrow \Omega^1(\mathcal{A}_{\hbar}) : \xi \mapsto \sum_i a_{\xi}^i db_{\xi}^i.$$

Such μ_{\hbar} defines an action of $\mathcal{U}_{\hbar}(\mathfrak{g})$ on \mathcal{A}_{\hbar} via the map

$$\Omega^1(\mathcal{A}_{\hbar}) \rightarrow C^1(\mathcal{A}_{\hbar}, \mathcal{A}_{\hbar}).$$

Extension

Definition

A quantum momentum map is defined to be a linear map

$$\mu_{\hbar} : T(\mathcal{U}_{\hbar}(\mathfrak{g})[1]) \rightarrow \Omega^{\bullet}(\mathcal{A}_{\hbar}) : \xi_1 \otimes \cdots \otimes \xi_n \mapsto a_1 db_1 \otimes \cdots \otimes a_n db_n$$

such that

$$\Phi_{\hbar}(\xi_1 \otimes \cdots \otimes \xi_n)(f_1, \dots, f_n) = \frac{1}{\hbar^n} a_1[b_1, f_1] \dots a_n[b_n, f_n]$$

Quantum Reduction

Definition

Let \mathcal{I}_{\hbar} be the left ideal of \mathcal{A}_{\hbar} generated by μ_{\hbar} . The action of $\mathcal{U}_{\hbar}(\mathfrak{g})$ descends to an action on $\mathcal{A}_{\hbar}/\mathcal{I}_{\hbar}$ and we define the reduced algebra by

$$\mathcal{A}_{\hbar}^{red} = (\mathcal{A}_{\hbar}/\mathcal{I}_{\hbar})^{\mathcal{U}_{\hbar}(\mathfrak{g})}$$

Hopf algebra action condition

Assume that ξ acts by

$$\Phi_{\hbar}(\xi) = \frac{1}{\hbar} a[b, \cdot]$$

for some $a, b \in C_{\hbar}^{\infty}(M)$. Note that $a \neq 0$ as soon as ξ is not killed by the cocycle δ .

Hopf algebra action $\implies \Phi_{\hbar}(\eta) = \frac{1}{\hbar} a[a^{-1}, \cdot]$

$$\Delta_{\hbar}(\xi) = \xi \otimes 1 - \hbar \eta \otimes \xi + 1 \otimes \xi$$

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Algebra homomorphism condition

We calculate the bracket of generators to get the deformed algebra structure of \mathfrak{g} :

$$\begin{aligned} [\Phi_{\hbar}(\xi), \Phi_{\hbar}(\eta)] f &= \frac{1}{\hbar^2} (a[b, a[a^{-1}, f]] - a[a^{-1}, a[b, f]]) \\ &= a[b, a][a^{-1}, f] + a^2[[b, a^{-1}], f]. \end{aligned}$$

Imposing that Φ_{\hbar} is a Lie algebra homomorphism we obtain different algebra structures that we discuss case by case.

Two dimensions: $[a, b] = 0$

Consider the Lie bialgebra $\mathfrak{g} = \mathbb{R}^2$ with generators ξ, η and a deformation quantization $C_{\hbar}^{\infty}(M)$ of a Poisson manifold M .

Algebra homomorphism $\implies \mathcal{U}_{\hbar}(\mathbb{R}^2)$ generated by $[\xi, \eta] = 0$

deformation quantization of

Abelian Lie bialgebra $\mathfrak{g} = \mathbb{R}^2$ with cobracket

$$\delta(\xi) = -\frac{1}{2}\eta \wedge \xi$$

$$\delta(\eta) = 0$$

Two dimensions: $[a, b] = 0$

Classical action

$$\Phi(\xi) = a_0 \{b_0, \cdot\}$$

$$\Phi(\eta) = a_0 \{a_0^{-1}, \cdot\}.$$

Quantum reduction

$$(C_{\hbar}^{\infty}(M)/\mathcal{I}_{\hbar})^{\mathcal{U}_{\hbar}(\mathbb{R}^2)} = \{a = \lambda, b = \mu\}^{\mathcal{U}_{\hbar}(\mathbb{R}^2)}$$

Quantization of the Poisson reduced algebra

$$(C^{\infty}(M)/\mathcal{I})^{\mathbb{R}^2} = \{a_0 = \lambda, b_0 = \mu\}^{\mathbb{R}^2}$$

Three dimensions: $\mathfrak{su}(2)$

Consider $a, b, c \in C_{\hbar}^{\infty}(M)$ satisfying

$$aba^{-1} = e^{2\hbar} b$$

$$aca^{-1} = e^{-2\hbar} c$$

$$[b, c] = \frac{\hbar^2}{e^{-\hbar} - e^{\hbar}} a^{-2} - (1 - e^{2\hbar}) cb$$

and the generators ξ, η, ζ acting respectively by

$$\Phi_{\hbar}(\xi)f = \frac{1}{\hbar} a[b, f]$$

$$\Phi_{\hbar}(\eta)f = \frac{1}{\hbar} [c, f]a$$

$$\Phi_{\hbar}(\zeta)f = afa^{-1}.$$

Three dimensions: $\mathfrak{su}(2)$

1 Lie algebra homomorphism

$$\zeta \xi \zeta^{-1} = e^{2\hbar} \xi$$

$$\zeta \eta \zeta^{-1} = e^{-2\hbar} \eta$$

$$[\xi, \eta] = \frac{\zeta^{-1} - \zeta}{e^{-\hbar} - e^{\hbar}}$$

2 Hopf algebra action

$$\Delta_{\hbar}(\zeta) = \zeta \otimes \zeta$$

$$\Delta_{\hbar}(\xi) = \xi \otimes 1 + \zeta \otimes \xi$$

$$\Delta_{\hbar}(\eta) = 1 \otimes \eta + \eta \otimes \zeta^{-1}.$$

Three dimensions: $\mathfrak{su}(2)$

Let

$$\Lambda = a^{-2} - e^{\hbar} \frac{(1 - e^{2\hbar})^2}{\hbar^2} cb$$

The ideal \mathcal{I}_{\hbar} generated by Λ in \mathcal{A}_{\hbar} is $\mathcal{U}_{\hbar}(\mathfrak{su}(2))$ -invariant, and

$$(C_{\hbar}^{\infty}(M)/\mathcal{I}_{\hbar})^{\mathcal{U}_{\hbar}(\mathfrak{su}(2))}$$

is the deformation quantization of the Poisson reduction

$$M//SU(2)$$

corresponding to the symplectic leaf $a_0^{-2} - 4b_0c_0 = 0$ in $SU(2)^* = SB(2, \mathbb{C})$.

Idea

Given a Poisson manifold (M, π) we can associate $\Sigma(M) \rightrightarrows M$.

Question: given a Poisson action of (G, π_G) on (M, π) , can we associate an action of (G, π_G) on $\Sigma(M)$?

Theorem

Given a Poisson action $G \times M \rightarrow M$ there exists a lifted Poisson action of G on $\Sigma(M)$ which is Hamiltonian with momentum map $J : \Sigma(M) \rightarrow G^$.*

Groupoid

Groupoid Γ over M is defined by

$$\Gamma \rightrightarrows M$$

Groupoid

Groupoid Γ over M is defined by

$$\Gamma \rightrightarrows M$$

Composition map $m : \Gamma_2 \rightarrow \Gamma$ where

$$\Gamma_2 = \{(g, h) \in \Gamma \times \Gamma \mid s(g) = t(h)\}$$

Unit map $u : M \rightarrow \Gamma : x \mapsto 1_x$

Inverse map $i : \Gamma \rightarrow \Gamma : g \mapsto g^{-1}$

Lie groupoid

Definition

A Lie groupoid is a groupoid $(\Gamma \rightrightarrows M, m, i, u)$ where

- Γ and M are manifolds
- s, t, m, i and u are smooth maps
- s and t are submersions

Symplectic groupoid

Definition

A Poisson groupoid is a Lie groupoid $\Gamma \rightrightarrows M$ with a multiplicative structure π on Γ .

When π is non degenerate, $\Omega = \pi^{-1}$ is a symplectic form. Thus, $(\Gamma \rightrightarrows M, \Omega)$ symplectic groupoid.

Poisson action on symplectic groupoid

A momentum map for the action of a Poisson Lie group (G, π_G) on symplectic groupoid $\Gamma \rightrightarrows M$ is a map $J : \Gamma \rightarrow G^*$ such that

$$\xi_M = \pi^\sharp(J^*(\theta_\xi))$$

Theorem

If $G \times \Gamma \rightarrow \Gamma$ is Hamiltonian action with momentum map $J : \Gamma \rightarrow G^$ such that $J(M) = e$ then the following are equivalent:*

- 1 $J : \Gamma \rightarrow G^*$ is a groupoid morphism
- 2 twisted multiplicativity

Quantization of J

Hamiltonian action $G \times \Gamma \rightarrow \Gamma$ with momentum map $J : \Gamma \rightarrow G^*$

- 1 Quantize symplectic groupoid given by quantum groupoid

$$C_{\hbar}^{\infty}(M) \rightrightarrows C_{\hbar}^{\infty}(\Gamma)$$

- 2 Quantize action $\mathfrak{g} \rightarrow \text{End } C^{\infty}(\Gamma)$
- 3 Quantize groupoid homomorphism $\alpha : \mathfrak{g} \rightarrow \Omega^1(\Gamma)$

Lifted momentum map

Hamiltonian action $G \times M \rightarrow M$ with momentum map

$$\mu : M \rightarrow G^*$$

If (M, π) is integrable, we associate $\Sigma(M) \rightrightarrows M$ and we get a lifted Hamiltonian action $G \times \Sigma(M) \rightarrow \Sigma(M)$ with momentum map

$$J : \Sigma(M) \rightarrow M : J(x) = \mu(t(x))\mu(s(x))^{-1}$$

Thanks!

Now we have our “towards” :)