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SPHERICAL HARMONICS ON GRASSMANNIANS

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#### Abstract

We propose a generalization of the theory of spherical harmonics to the context of symmetric subgroups of reductive groups acting on flag manifolds. We give some sample results for the case of the orthogonal group acting on Grassmann manifolds, especially the case of 2-planes.


1. Introduction. The theory of spherical harmonics is a classical piece of mathematics, with manifold applications to physics, in any situation where the dependence of a quantity on direction in space needs to be described.

A representation-theoretic understanding of the theory is that it describes the behavior of polynomials under the action of the orthogonal group. If $\mathcal{P}\left(\mathbb{R}^{n}\right)=\mathcal{P}$ is the space of complex-valued polynomial functions in $n$ (real) variables, then the group $\mathrm{GL}_{n}(\mathbb{R})=\mathrm{GL}_{n}$ of linear transformations on $\mathbb{R}^{n}$ also acts on the polynomial functions by the standard recipe:

$$
\begin{equation*}
g(p)(\vec{x})=p\left(g^{-1}(\vec{x})\right) . \tag{1.1}
\end{equation*}
$$

Here $g \in \mathrm{GL}_{n}, p \in \mathcal{P}$, and $\vec{x}$ is a point in $\mathbb{R}^{n}$. (We do not distinguish notationally here between $g$ as a linear transformation on $\mathbb{R}^{n}$ and on $\mathcal{P}$; which is meant should be clear from context.)

Let $\mathcal{P}^{d}\left(\mathbb{R}^{n}\right)=\mathcal{P}^{d}$ be the space of polynomials homogeneous of degree $d$. Then it is well known that

$$
\begin{equation*}
\mathcal{P} \simeq \bigoplus_{d=0}^{\infty} \mathcal{P}^{d} \tag{1.2}
\end{equation*}
$$

that the action (1.1) of $\mathrm{GL}_{n}$ preserves each space $\mathcal{P}^{d}$, and that the $\mathcal{P}^{d}$ are irreducible representations for $\mathrm{GL}_{n}$. However, if $\mathrm{O}_{n} \subset \mathrm{GL}_{n}$ is the group of orthogonal transformations, defined to be the linear transformations which preserve the usual Euclidean (squared) length $r^{2}=\sum_{i=1}^{n} x_{i}^{2}$, then $\mathrm{O}_{n}$ does not act irreducibly on the homogeneous components $\mathcal{P}^{d}$. Indeed, it is easy

[^0]to find operators on $\mathcal{P}$ that commute with the action of elements of $\mathrm{O}_{n}$, and these operators allow one to display subspaces of $\mathcal{P}^{d}$ that are invariant under $\mathrm{O}_{n}$.

The most obvious such operator is multiplication by $r^{2}: p \mapsto r^{2} p$. Since $r^{2}$ is homogeneous of degree two, multiplying by it raises the degree of a polynomial by two, so we get mappings

$$
\begin{equation*}
r^{2}: \mathcal{P}^{d} \rightarrow \mathcal{P}^{d+2} \tag{1.3}
\end{equation*}
$$

for each $d$. Since operators $g$ defined in (1.1) are algebra automorphisms of $\mathcal{P}$, it is easy to check that $g\left(r^{2} p\right)=r^{2} g(p)$ for $g \in \mathrm{O}_{n}$. It follows that the space $r^{2}\left(\mathcal{P}^{d}\right) \subset \mathcal{P}^{d+2}$ is invariant under $\mathrm{O}_{n}$.

Another operator that commutes with $\mathrm{O}_{n}$, less obvious than $r^{2}$, but which presented itself in physics, is the Laplacian $\Delta=\sum_{i=0}^{n} \partial^{2} / \partial x_{i}^{2}$. The Laplacian reduces the degree of polynomials by two; it defines mappings

$$
\begin{equation*}
\Delta: \mathcal{P}^{d} \rightarrow \mathcal{P}^{d-2} \tag{1.4}
\end{equation*}
$$

Since these maps commute with the action of $\mathrm{O}_{n}$, it follows that the kernel

$$
\begin{equation*}
\mathcal{H}^{d}\left(\mathbb{R}^{n}\right)=\mathcal{H}^{d}=\left\{p \in \mathcal{P}^{d}: \Delta p=0\right\} \tag{1.5}
\end{equation*}
$$

of $\Delta$, commonly known as the harmonic polynomials, is an $\mathrm{O}_{n}$-invariant subspace of $\mathcal{P}^{d}$.

The main assertions of the theory of spherical harmonics are
Proposition 1.1 (Theory of Spherical Harmonics).
(a) The space $\mathcal{H}^{d}$ of harmonic polynomials of degree $d$ is an irreducible representation of $\mathrm{O}_{n}$.
(b) $\mathcal{P}^{d} \simeq \mathcal{H}^{d} \oplus r^{2} \mathcal{P}^{d-2}$.
(c) $\mathcal{P}^{d} \simeq \bigoplus_{k=0}^{[d / 2]} r^{2 k} \mathcal{H}^{d-2 k}$ is a decomposition of $\mathcal{P}^{d}$ into irreducible representations for $\mathrm{O}_{n}$.
Evidently, the main statement (c) follows from (a) and (b).
One way to prove these results is to study the interaction between the operators $r^{2}$ and $\Delta$. They do not commute with each other, but it turns out that there is an elegant formula for their commutator:

$$
\begin{equation*}
\left[\Delta, r^{2}\right]=\Delta r^{2}-r^{2} \Delta=4 E+2 n \tag{1.6}
\end{equation*}
$$

where $E=\sum_{i=1}^{n} x_{i} \partial / \partial x_{i}$ is the Euler degree operator, which acts on $\mathcal{P}^{d}$ by the scalar $d$. From this, one can show by induction that, if $h \in \mathcal{H}^{\ell}$, then

$$
\begin{equation*}
\Delta\left(r^{2 k} h\right)=2 k(2(\ell+k-1)+n) r^{2(k-1)} h \tag{1.7}
\end{equation*}
$$

From (1.7) it is evident that, if the decomposition (c) is true for $\mathcal{P}^{d-2}$, then $\operatorname{ker} \Delta \cap r^{2} \mathcal{P}^{d-2}=\{0\}$. Hence $\operatorname{dim} \operatorname{ker} \Delta \leq \operatorname{dim} \mathcal{P}^{d}-\operatorname{dim} \mathcal{P}^{d-2}$. On the other hand, since $\Delta$ maps $\mathcal{P}^{d}$ to $\mathcal{P}^{d-2}$, it is likewise clear that $\operatorname{dim} \operatorname{ker} \Delta \geq$
$\operatorname{dim} \mathcal{P}^{d}-\operatorname{dim} \mathcal{P}^{d-2}$. Hence $\operatorname{dim} \operatorname{ker} \Delta=\operatorname{dim} \mathcal{P}^{d}-\operatorname{dim} \mathcal{P}^{d-2}$, and equation (b) holds for $\mathcal{P}^{d}$, whence equation (c) does also.

In this situation, a remarkable fact is that the three operators $r^{2}, \Delta$ and

$$
\begin{equation*}
\left[\Delta, r^{2}\right]=4 E+2 n \tag{1.8}
\end{equation*}
$$

span a three-dimensional Lie algebra isomorphic to the three-dimensional simple Lie algebra $\mathfrak{s l}_{2}$. Furthermore, the associative algebra generated by this Lie algebra is the full algebra of polynomial coefficient differential operators which commute with the action of $\mathrm{O}_{n}$ on $\mathcal{P}$ (GW], Ho ). This in turn implies statement (a) of Proposition 1.1. (There are also ways of proving (a) that do not involve observing the existence of the commuting $\mathfrak{s l}_{2}$.)

We are interested in a generalization of the theory of spherical harmonics. This theory is usually thought of as describing functions on the sphere, but statement (c) of Proposition 1.1 in fact describes the relationship between the space $\mathcal{H}^{d}$, which is the $\mathrm{O}_{n}$ components of the functions on the sphere, and the space $\mathcal{P}$ of all polynomials. Therefore, it is reasonable in thinking about generalizing the theory of spherical harmonics to consider the nature of $\mathcal{P}$, in particular as a module for $\mathrm{GL}_{n}$.

The polynomial ring $\mathcal{P}$ can be thought of as the "homogeneous coordinate ring" ( $[\mathrm{F}]$ ) of the projective space $\mathbb{P}^{n-1}$ of lines through the origin in $\mathbb{R}^{n}$. Projective space is the simplest of the flag manifolds associated to $\mathrm{GL}_{n}$. Other flag manifolds for $\mathrm{GL}_{n}$ include the Grassmann varieties of $k$-dimensional subspaces of $\mathbb{R}^{n}$, or higher flag manifolds associated to nested sequences of subspace of specified dimensions. Each flag manifold has an associated homogeneous coordinate ring, which is a module for $\mathrm{GL}_{n}$, with a well understood decomposition into irreducible representations. One could think of the theory of "flag harmonics" to be concerned with the decomposition of these $\mathrm{GL}_{n}$ representations into irreducible subspaces for $\mathrm{O}_{n}$.

But there is more to the theory of spherical harmonics than just the irreducible decomposition. There is the action of the $\mathrm{O}_{n}$ invariant $r^{2}$, and even more directly implicated in the name, there is the understanding of the operator $\Delta$, and especially its kernel, the harmonic functions.

One can ask if there is some analog of this structure, in particular, analogs of $r^{2}$ and $\Delta$, for other flag manifolds. This is easy to answer for $r^{2}$. The invariants of $\mathrm{O}_{n}$ in representations of $\mathrm{GL}_{n}$ are well-understood; this is a special case of the Cartan-Helgason Theorem (see Theorem 4.1 on page 535 of [He]). In particular, for each Grassmann variety, the algebra of $\mathrm{O}_{n}$ invariants is generated by a single well-understood element that bears a clear relation to $r^{2}$.

It is rather less obvious, but it turns out that there also exist differential operators analogous to the Laplacian. Together, they generate an algebra of operators that commute with the action of $\mathrm{O}_{n}$. One can think of the theory
of spherical harmonics as a description of the action of this algebra, and most particularly, of the action of the Laplacian analog.

As we have seen above, in the classical theory of spherical harmonics, one has the remarkable circumstance that the Laplacian and $r^{2}$ generate a Lie algebra. The commutation relations in this Lie algebra let one give a precise description, not just of the kernel of $\Delta$, but of the action of $\Delta$ on all of $\mathcal{P}$. Although one cannot expect that the analogs of $r^{2}$ and $\Delta$ will live inside a finite-dimensional Lie algebra for more general flag manifolds, still one might hope to give a reasonably explicit description of the action of the $\Delta$ analog. In this paper, we present an example that provides some encouragement for such a hope.
2. Notation. The action of $\mathrm{GL}_{n}(\mathbb{R})$ on $\mathcal{P}=\mathcal{P}\left(\mathbb{R}^{n}\right)$ can be complexified to give a representation of $\mathrm{GL}_{n}(\mathbb{R})$ on the algebra $\mathcal{P}\left(\mathbb{C}^{n}\right)$ of polynomial functions on $\mathbb{C}^{n}$. This action can be extended to an action by the complexification of $\mathrm{GL}_{n}(\mathbb{R})$, which is the complex general linear group $\mathrm{GL}_{n}(\mathbb{C})$ consisting of all $n \times n$ invertible complex matrices. The complexification of $\mathrm{O}_{n}$ is the complex orthogonal group $\mathrm{O}_{n}(\mathbb{C})$. By a suitable change of coordinates on $\mathbb{C}^{n}$, we can assume that $\mathrm{O}_{n}(\mathbb{C})$ is the subgroup of $\mathrm{GL}_{n}(\mathbb{C})$ which preserves the symmetric bilinear form

$$
\left\langle\left(\begin{array}{c}
u_{1}  \tag{2.1}\\
\vdots \\
u_{n}
\end{array}\right),\left(\begin{array}{c}
v_{1} \\
\vdots \\
v_{n}
\end{array}\right)\right\rangle=\sum_{j=1}^{n} u_{j} v_{n+1-j}
$$

on $\mathbb{C}^{n}$. In the rest of the paper, we shall work with $\mathrm{GL}_{n}(\mathbb{C})$ and $\mathrm{O}_{n}(\mathbb{C})$. So from now on, $\mathrm{GL}_{n}$ and $\mathrm{O}_{n}$ shall stand for $\mathrm{GL}_{n}(\mathbb{C})$ and $\mathrm{O}_{n}(\mathbb{C})$ respectively.

We now introduce notation for the irreducible representations of $\mathrm{GL}_{n}$. Let $B_{n}=A_{n} U_{n}$ be the standard Borel subgroup of upper triangular matrices in $\mathrm{GL}_{n}$, where $A_{n}$ is the diagonal torus in $\mathrm{GL}_{n}$ and $U_{n}$ is the maximal unipotent subgroup consisting of all the upper triangular matrices with 1 's on the diagonal. Recall that a Young diagram $\lambda$ is an array of square boxes arranged in left-justified horizontal rows, with each row no longer than the one above it ([|F]). If $\lambda$ has at most $m$ rows, then we shall write it as

$$
\lambda=\left(\lambda_{1}, \ldots, \lambda_{m}\right)
$$

where for each $i, \lambda_{i}$ is the number of boxes in the $i$ th row of $\lambda$. We shall denote the number of rows in $\lambda$ by $d(\lambda)$, and call it the depth of $\lambda$. For later use, we let $\mathbf{1}_{k}$ be the Young diagram with only one column of $k$ boxes, i.e.

$$
\begin{equation*}
\mathbf{1}_{k}=(\overbrace{1, \ldots, 1}^{k}) . \tag{2.2}
\end{equation*}
$$

For a Young diagram $\lambda=\left(\lambda_{1}, \ldots, \lambda_{n}\right)$ with at most $n$ rows, let $\psi_{n}^{\lambda}$ : $A_{n} \rightarrow \mathbb{C}^{\times}$be the character given by

$$
\psi_{n}^{\lambda}\left[\operatorname{diag}\left(a_{1}, \ldots, a_{n}\right)\right]=a_{1}^{\lambda_{1}} \cdots a_{n}^{\lambda_{n}} .
$$

Here $\operatorname{diag}\left(a_{1}, \ldots, a_{n}\right)$ is the $n \times n$ diagonal matrix such that its diagonal entries are $a_{1}, \ldots, a_{n}$. Then $\psi_{n}^{\lambda}$ is a dominant weight for $\mathrm{GL}_{n}$ with respect to the Borel subgroup $B_{n}(\underline{G W})$, and we shall denote the irreducible representation of $\mathrm{GL}_{n}$ with highest weight $\psi_{n}^{\lambda}$ by $\rho_{n}^{\lambda}$. We shall abuse notation and say that the highest weight of $\rho_{n}^{\lambda}$ is $\lambda$. If $\lambda=m \mathbf{1}_{n}$, then we also say that the highest weight of $\rho_{n}^{\lambda}$ is $\operatorname{det}_{n}^{m}$.

The irreducible finite-dimensional representations of $\mathrm{O}_{n}$ are parameterized by Young diagrams $\lambda$ such that the sum of the lengths of the first two columns of $\lambda$ does not exceed $n$ ( $\mathrm{Wy},[\mathrm{GW},[\mathrm{Ho})$. For such a Young diagram $\lambda$, we shall denote the $\mathrm{O}_{n}$ representation associated with $\lambda$ by $\sigma_{n}^{\lambda}$. Specifically, $\sigma_{n}^{\lambda}$ is the irreducible representation of $\mathrm{O}_{n}$ generated by the $\mathrm{GL}_{n}$ highest weight vector in $\rho_{n}^{\lambda}$. See Section 3.6 of [HO for more details.

Let $\mathrm{SO}_{n}$ denote the subgroup of $\mathrm{O}_{n}$ consisting of all elements of $\mathrm{O}_{n}$ with determinant 1 , and let

$$
\begin{equation*}
A_{\mathrm{SO}_{n}}=A_{n} \cap \mathrm{SO}_{n}, \quad N_{n}=U_{n} \cap \mathrm{SO}_{n} . \tag{2.3}
\end{equation*}
$$

Explicitly,
$A_{\mathrm{SO}_{n}}= \begin{cases}\left\{\operatorname{diag}\left(a_{1}, \ldots, a_{m}, a_{m}^{-1}, \ldots, a_{1}^{-1}\right): a_{1}, \ldots, a_{m} \in \mathbb{C}^{\times}\right\}, & n=2 m, \\ \left\{\operatorname{diag}\left(a_{1}, \ldots, a_{m}, 1, a_{m}^{-1}, \ldots, a_{1}^{-1}\right): a_{1}, \ldots, a_{m} \in \mathbb{C}^{\times}\right\}, & n=2 m+1 .\end{cases}$
Let $\lambda$ be a Young diagram such that the sum of the lengths of its first two columns does not exceed $n$. If $2 d(\lambda) \neq n$, then the restriction of $\sigma_{n}^{\lambda}$ to $\mathrm{SO}_{n}$ is irreducible. If in addition $2 d(\lambda)<n$ and $\phi_{n}^{\lambda}: A_{\mathrm{SO}_{n}} \rightarrow \mathbb{C}^{\times}$is the restriction of the character $\psi_{n}^{\lambda}$ to $A_{\mathrm{SO}_{n}}$, then as an $\mathrm{SO}_{n}$ module, $\sigma_{n}^{\lambda}$ has highest weight $\phi_{n}^{\lambda}$. In this case, we shall abuse notation and say that $\lambda$ is the highest weight of $\sigma_{n}^{\lambda}$.
3. The $\mathrm{O}_{n}$ highest weight vectors on Grassmannians. Let $2 k<n$ and let $\mathbb{G}_{k}^{n}$ be the set of all $k$-dimensional subspaces of $\mathbb{C}^{n}$. It has a structure of a projective variety. Let $\mathcal{R}\left(\mathbb{G}_{k}^{n}\right)$ be the homogeneous coordinate ring ( ${ }^{1}$ ) of $\mathbb{G}_{k}^{n}$. It carries an action by $\mathrm{GL}_{n} \times \mathrm{GL}_{k}$, and $\mathrm{O}_{n}$ acts by re-

[^1]striction from $\mathrm{GL}_{n}$. In this section, we shall describe the $\mathrm{O}_{n} \times \mathrm{GL}_{k}$ module structure of $\mathcal{R}\left(\mathbb{G}_{k}^{n}\right)$.

Let $\mathrm{M}_{n k}=\mathrm{M}_{n k}(\mathbb{C})$ be the space of $n \times k$ complex matrices, and let $\mathcal{P}\left(\mathrm{M}_{n k}\right)$ be the algebra of polynomial functions on $\mathrm{M}_{n k}$, that is, each $p \in$ $\mathcal{P}\left(\mathrm{M}_{n k}\right)$ is of the form

$$
p(x)=\sum_{\alpha} a_{\alpha} x^{\alpha},
$$

where $x=\left(x_{i j}\right) \in \mathrm{M}_{n k}$, each $\alpha=\left(\alpha_{i j}\right)$ appearing in the sum is an $n \times k$ matrix of nonnegative integers, $a_{\alpha} \in \mathbb{C}$ and

$$
x^{\alpha}=\prod_{i, j} x_{i j}^{\alpha_{i j}} .
$$

We now define an action of $\mathrm{GL}_{n} \times \mathrm{GL}_{k}$ on $\mathrm{M}_{n k}$. For each $n \times n$ complex matrix $A$, there exists a unique $n \times n$ complex matrix $A^{\tau}$ such that

$$
\langle A u, v\rangle=\left\langle u, A^{\tau} v\right\rangle,
$$

where $\langle\cdot, \cdot\rangle$ is given in (2.1). For $g \in \mathrm{GL}_{n}, h \in \mathrm{GL}_{k}$, and $T \in \mathrm{M}_{n k}$, let

$$
\begin{equation*}
(g, h)(T)=\left(g^{-1}\right)^{\tau} T h^{-1}, \quad g \in \mathrm{GL}_{n}, h \in \mathrm{GL}_{k}, T \in \mathrm{M}_{n k} \tag{3.1}
\end{equation*}
$$

This action induces an action of $\mathrm{GL}_{n} \times \mathrm{GL}_{k}$ on $\mathcal{P}\left(\mathrm{M}_{n k}\right)$ in the usual way. We have used $\left(g^{-1}\right)^{\tau}$ in the action by $\mathrm{GL}_{n}$ so that this gives rise to a more symmetrical decomposition of $\mathcal{P}\left(\mathrm{M}_{n k}\right)$ into irreducible $\mathrm{GL}_{n} \times \mathrm{GL}_{k}$ representations (see (3.2) below). Moreover, $g \in \mathrm{O}_{n}$ if and only if $g=\left(g^{-1}\right)^{\tau}$.

Let $\mathrm{SL}_{k}$ be the subgroup of $\mathrm{GL}_{k}$ consisting of all elements of $\mathrm{GL}_{k}$ with determinant 1 . Then it is well known ([F]) that $\mathcal{R}\left(\mathbb{G}_{k}^{n}\right)$ can be identified with the algebra

$$
\mathfrak{A}=\mathcal{P}\left(\mathrm{M}_{n k}\right)^{\mathrm{SL}_{k}}
$$

of $\mathrm{SL}_{k}$ invariants in $\mathcal{P}\left(\mathrm{M}_{n k}\right)$. Thus in the remaining part of the paper, we shall replace $\mathcal{R}\left(\mathbb{G}_{k}^{n}\right)$ by $\mathfrak{A}$.

By the $\left(\mathrm{GL}_{n}, \mathrm{GL}_{k}\right)$-duality ( $(\mathrm{Ho})$ ), under the action by $\mathrm{GL}_{n} \times \mathrm{GL}_{k}$, we have the decomposition

$$
\begin{equation*}
\mathcal{P}\left(\mathrm{M}_{n k}\right)=\bigoplus_{d(\lambda) \leq k} \rho_{n}^{\lambda} \otimes \rho_{k}^{\lambda} . \tag{3.2}
\end{equation*}
$$

By extracting the $\mathrm{SL}_{k}$ invariants, we obtain

$$
\mathfrak{A}=\mathcal{P}\left(\mathrm{M}_{n k}\right)^{\mathrm{SL}_{k}}=\bigoplus_{d(\lambda) \leq k} \rho_{n}^{\lambda} \otimes\left(\rho_{k}^{\lambda}\right)^{\mathrm{SL}_{k}},
$$

where $\left(\rho_{k}^{\lambda}\right)^{\mathrm{SL}_{k}}$ denotes the space of $\mathrm{SL}_{k}$ invariant vectors in $\rho_{n}^{\lambda}$. Now $\left(\rho_{k}^{\lambda}\right)^{\mathrm{SL}_{k}}$ $\neq 0$ if and only if $\lambda=m \mathbf{1}_{k}$ for some nonnegative integer $m$, i.e. $\rho_{n}^{\lambda}$ is the one-dimensional space on which $\mathrm{GL}_{k}$ acts by the character $h \mapsto(\operatorname{det} h)^{m}$.

We shall write $\rho_{k}^{m \mathbf{1}_{k}}$ as $\operatorname{det}_{k}^{m}$. Then under the action by $\mathrm{GL}_{n} \times \mathrm{GL}_{k}$,

$$
\begin{equation*}
\mathfrak{A}=\bigoplus_{m=0}^{\infty} \mathfrak{A}_{m} \tag{3.3}
\end{equation*}
$$

where for each $m$,

$$
\mathfrak{A}_{m} \cong \rho_{n}^{m \mathbf{1}_{k}} \otimes \operatorname{det}_{k}^{m}
$$

Next, recall that $N_{n}$ is the standard maximal unipotent subgroup in $\mathrm{SO}_{n}$ given in (2.3). Let

$$
\mathfrak{A}^{N_{n}}=\mathcal{P}\left(\mathrm{M}_{n k}\right)^{N_{n} \times \mathrm{SL}_{k}}
$$

be the algebra of $N_{n}$ invariants in $\mathfrak{A}$. We now describe the generators of $\mathfrak{A}^{N_{n}}$. For $1 \leq i, j \leq k$ and $x \in \mathrm{M}_{n k}$, let

$$
r_{i j}^{2}(x)=\left\langle x_{i}, x_{j}\right\rangle
$$

where $\langle\cdot, \cdot\rangle$ is the symmetric bilinear form given in 2.1), and $x_{i}$ and $x_{j}$ are the $i$ th and $j$ th columns of $x$ respectively. Then the algebra $\mathcal{P}\left(\mathrm{M}_{n k}\right)^{\mathrm{O}_{n}}$ of $\mathrm{O}_{n}$ invariants in $\mathcal{P}\left(\mathrm{M}_{n k}\right)$ is a polynomial algebra on the generators $\left\{r_{i j}^{2}\right.$ : $1 \leq i \leq j \leq k\}([\mathrm{Ho}$, GW $])$. Define

$$
\left.\left.\gamma_{j}=\left\lvert\, \begin{array}{ccccc} 
& & & x_{11} & \cdots \\
x_{1 k} \\
& 0 & & \vdots & \\
& & & x_{j 1} & \cdots \\
x_{j k} \\
x_{11} & \cdots & x_{j 1} & r_{11}^{2} & \cdots \\
r_{1 k} \\
\vdots & & \vdots & \vdots & \\
x_{1 k} & \cdots & x_{j k} & r_{k 1}^{2} & \cdots
\end{array}\right.\right) \quad r_{k k}^{2} . \mid 0 \leq j \leq k-1\right)
$$

and

$$
\gamma_{k}=\left|\begin{array}{ccc}
x_{11} & \cdots & x_{1 k} \\
\vdots & & \vdots \\
x_{k 1} & \cdots & x_{k k}
\end{array}\right|
$$

Here, the vertical lines in $\gamma_{j}$ indicate determinant. These polynomials are joint $\mathrm{SO}_{n} \times \mathrm{GL}_{k}$ highest weight vectors with the following weights:

|  | $\mathrm{SO}_{n}$ weight | $\mathrm{GL}_{k}$ weight |
| :---: | :---: | :---: |
| $\gamma_{0}$ | 0 | $\operatorname{det}_{k}^{2}$ |
| $\gamma_{j}(1 \leq j \leq k-1)$ | $2 \mathbf{1}_{j}$ | $\operatorname{det}_{k}^{2}$ |
| $\gamma_{k}$ | $\mathbf{1}_{k}$ | $\operatorname{det}_{k}$ |

Proposition 3.1 ( $\lfloor\boxed{A T Z}]$ ). The algebra $\mathfrak{A}^{N_{n}}$ is a polynomial algebra on the generators $\gamma_{0}, \gamma_{1}, \ldots, \gamma_{k}$.

Let $\mathbf{a}=\left(a_{0}, a_{1}, \ldots, a_{k}\right) \in \mathbb{Z}_{\geq 0}^{k+1}$, and

$$
\gamma^{\mathbf{a}}=\gamma_{0}^{a_{0}} \gamma_{1}^{a_{1}} \cdots \gamma_{k}^{a_{k}}
$$

Then $\gamma^{\mathbf{a}}$ is an $\mathrm{SO}_{n} \times \mathrm{GL}_{k}$ highest weight vector, and it has $\mathrm{SO}_{n}$ weight $\lambda=\left(\lambda_{1}, \ldots, \lambda_{k}\right)$ and $\mathrm{GL}_{k}$ weight $\operatorname{det}_{k}^{m(\mathbf{a})}$ where

$$
\lambda_{j}=2 \sum_{i=j}^{k-1} a_{i}+a_{k} \quad(1 \leq j \leq k-1), \quad \lambda_{k}=a_{k}
$$

and

$$
\begin{equation*}
m(\mathbf{a})=2 \sum_{i=0}^{k-1} a_{i}+a_{k} \tag{3.4}
\end{equation*}
$$

Let $\mathfrak{A}_{m, \mathbf{a}}$ be the irreducible $\mathrm{O}_{n}$ module generated by $\gamma^{\mathbf{a}}$. Then as a representation of $\mathrm{O}_{n}, \mathfrak{A}_{m}$ admits the decomposition

$$
\mathfrak{A}_{m}=\bigoplus_{\mathbf{a}} \mathfrak{A}_{m, \mathbf{a}}
$$

where the sum is taken over all $\mathbf{a} \in \mathbb{Z}_{>0}^{k+1}$ such that $m(\mathbf{a})=m$. For each $\mathfrak{A}_{m, \mathbf{a}}$ which appears in the sum, its $\mathrm{SO}_{n}$ highest weight $\lambda$ is such that $\lambda_{1} \leq m$ and $\lambda_{j} \equiv m(\bmod 2)$ for all $j$. Conversely, every such $\lambda$ is the $\mathrm{SO}_{n}$ highest weight of a unique $\mathfrak{A}_{m, \mathbf{a}}$ in the sum. It follows that

$$
\begin{equation*}
\mathfrak{A}_{m} \cong \bigoplus_{\substack{\lambda_{1} \leq m \\ \lambda_{j} \equiv m(\bmod 2)}} \sigma_{n}^{\lambda} \tag{3.5}
\end{equation*}
$$

which is $\mathrm{O}_{n}$ multiplicity free: any two irreducible $\mathrm{O}_{n}$ submodules are nonisomorphic.
4. The map $\partial: \mathcal{P}\left(\mathrm{M}_{n k}\right) \rightarrow \mathcal{D}\left(\mathrm{M}_{n k}\right)$. Let $\mathcal{D}\left(\mathrm{M}_{n k}\right)$ be the space of constant-coefficient differential operators on $\mathcal{P}\left(\mathrm{M}_{n k}\right)$, that is, $\mathcal{D}\left(\mathrm{M}_{n k}\right)$ consists of the operators of the form

$$
\sum_{\alpha} a_{\alpha} \frac{\partial^{|\alpha|}}{\partial x^{\alpha}}
$$

where each $\alpha=\left(\alpha_{i j}\right)$ appearing in the sum is an $n \times k$ matrix of nonnegative integers, $a_{\alpha} \in \mathbb{C}$,

$$
|\alpha|=\sum_{i, j} \alpha_{i j}
$$

and

$$
\frac{\partial^{|\alpha|}}{\partial x^{\alpha}}=\frac{\partial^{|\alpha|}}{\partial x_{11}^{\alpha_{11}} \partial x_{12}^{\alpha_{12}} \cdots \partial x_{n k}^{\alpha_{n k}}}
$$

The algebra $\mathcal{D}\left(\mathrm{M}_{n k}\right)$ is naturally isomorphic to the symmetric algebra $S\left(\mathrm{M}_{n k}\right)$ on $\mathrm{M}_{n k}$. In fact, for each $u \in \mathrm{M}_{n k}$ and $f \in \mathcal{P}\left(\mathrm{M}_{n k}\right)$, we let

$$
D_{u}(f)(v)=\lim _{t \rightarrow 0} \frac{f(v+t u)-f(v)}{t}, \quad v \in \mathrm{M}_{n k}
$$

Then the map $D: \mathrm{M}_{n k} \rightarrow \mathcal{D}\left(\mathrm{M}_{n k}\right)$ is linear, and extends uniquely to an algebra isomorphism $D: S\left(\mathrm{M}_{n k}\right) \rightarrow \mathcal{D}\left(\mathrm{M}_{n k}\right)$. Let $\mathrm{GL}_{n}$ act on $\mathrm{M}_{n k}$ as in (3.1), that is,

$$
g \cdot v=\left(g^{-1}\right)^{\tau} v, \quad g \in \mathrm{GL}_{n}, v \in \mathrm{M}_{n k}
$$

We extend this action to an action by $\mathrm{GL}_{n}$ on $S\left(\mathrm{M}_{n k}\right)$ by algebra automorphisms, and this in turn induces an action on $\mathcal{D}\left(\mathrm{M}_{n k}\right)$ via the map $D$.

For $1 \leq i \leq n$ and $1 \leq j \leq k$, let

$$
\partial\left(x_{i j}\right)=\frac{\partial}{\partial x_{n+1-i, j}}
$$

Then $\partial$ extends uniquely to an algebra isomorphism

$$
\partial: \mathcal{P}\left(\mathrm{M}_{n k}\right) \rightarrow \mathcal{D}\left(\mathrm{M}_{n k}\right)
$$

Specifically, if $p=\sum_{\alpha} a_{\alpha} x^{\alpha} \in \mathcal{P}\left(\mathrm{M}_{n k}\right)$, then

$$
\begin{equation*}
\partial(p)=\sum_{\alpha} a_{\alpha} \frac{\partial^{\left|\alpha^{\prime}\right|}}{\partial x^{\alpha^{\prime}}} \tag{4.1}
\end{equation*}
$$

where $\alpha^{\prime}=\left(\alpha_{i j}^{\prime}\right)$ and $\alpha_{i j}^{\prime}=\alpha_{n+1-i, j}$ for all $1 \leq i \leq n$ and $1 \leq j \leq k$.
LEMMA 4.1. The map $\partial: \mathcal{P}\left(\mathrm{M}_{n k}\right) \rightarrow \mathcal{D}\left(\mathrm{M}_{n k}\right)$ is an $\mathrm{O}_{n}$ map, that is, it commutes with the action of $\mathrm{O}_{n}$ on $\mathcal{P}\left(\mathrm{M}_{n k}\right)$ and $\mathcal{D}\left(\mathrm{M}_{n k}\right)$ defined by the restriction from the actions of $\mathrm{GL}_{n}$.

Proof. Define for $u, v \in \mathrm{M}_{n k}$,

$$
(u, v)=\left[\partial^{-1}\left(D_{v}\right)\right](u)
$$

For $1 \leq j \leq k$, let $u_{j}$ (respectively $v_{j}$ ) be the $j$ th column of $u$ (respectively $v$ ). Then

$$
(u, v)=\sum_{j=1}^{k}\left\langle u_{j}, v_{j}\right\rangle
$$

where $\langle\cdot, \cdot\rangle$ is the symmetric form on $\mathbb{C}^{n}$ given in 2.1. Hence for each $n \times n$ complex matrix $A$, we have

$$
(A u, v)=\left(u, A^{\tau} v\right)
$$

Now for $g \in \mathrm{GL}_{n}$,

$$
\begin{aligned}
{\left[\partial^{-1}\left(g \cdot D_{v}\right)\right](u) } & =\left[\partial^{-1}\left(D_{g \cdot v}\right)\right](u)=(u, g \cdot v)=\left(g^{\tau} \cdot u, v\right) \\
& =\left[\partial^{-1}\left(D_{v}\right)\right]\left(g^{\tau} \cdot u\right)=\left[\left(g^{-1}\right)^{\tau} \cdot\left(\partial^{-1}\left(D_{v}\right)\right)\right](u)
\end{aligned}
$$

This shows that $\partial^{-1}\left(g . D_{v}\right)=\left(g^{-1}\right)^{\tau} .\left(\partial^{-1}\left(D_{v}\right)\right)$ for every $v \in \mathrm{M}_{n k}$. Since $\mathrm{GL}_{n}$ acts on $\mathcal{P}\left(\mathrm{M}_{n k}\right)$ and $\mathcal{D}\left(\mathrm{M}_{n k}\right)$ by algebra automorphisms and $\partial^{-1}$ is an algebra isomorphism, we have $\partial^{-1} g=\left(g^{-1}\right)^{\tau} \partial^{-1}$ on $\mathcal{D}\left(\mathrm{M}_{n k}\right)$, or equivalently, $g \partial=\partial\left(g^{-1}\right)^{\tau}$ on $\mathcal{P}\left(\mathrm{M}_{n k}\right)$. In particular, $g \partial=\partial g$ for all $g \in \mathrm{O}_{n}$. This proves the lemma.

For $1 \leq i, j \leq k$, let

$$
\Delta_{i j}=\partial\left(r_{i j}^{2}\right)=\sum_{a=1}^{n} \frac{\partial^{2}}{\partial x_{a, i} \partial x_{n+1-a, j}} .
$$

We also let

$$
L=\partial\left(\gamma_{0}\right)=\left|\begin{array}{ccc}
\Delta_{11} & \cdots & \Delta_{1 k} \\
\vdots & & \vdots \\
\Delta_{n 1} & \cdots & \Delta_{n k}
\end{array}\right|
$$

Corollary 4.2. The map $L: \mathcal{P}\left(\mathrm{M}_{n k}\right) \rightarrow \mathcal{P}\left(\mathrm{M}_{n k}\right)$ is an $\mathrm{O}_{n}$ map.
Proof. This is because $L$ is an $\mathrm{O}_{n}$ invariant in $\mathcal{D}\left(\mathrm{M}_{n k}\right)$.
Next, we shall define an inner product on $\mathcal{P}\left(\mathrm{M}_{n k}\right)$ such that multiplication by $\gamma_{0}$ and the operator $L$ are adjoints of each other. If $p=\sum_{\alpha} a_{\alpha} x^{\alpha} \in$ $\mathcal{P}\left(\mathrm{M}_{n k}\right)$, we let

$$
\widetilde{p}(x)=\sum_{\alpha} \overline{a_{\alpha}} x^{\alpha} .
$$

Here for each $\alpha, \overline{a_{\alpha}}$ is the complex conjugate of $a_{\alpha}$. Then for $p, q \in \mathcal{P}\left(\mathrm{M}_{n k}\right)$, define

$$
\begin{equation*}
\langle p, q\rangle=\{[\partial(p)](\widetilde{q})\}(0) . \tag{4.2}
\end{equation*}
$$

Explicitly, if $p(x)=\sum_{\alpha} a_{\alpha} x^{\alpha} \in \mathcal{P}\left(\mathrm{M}_{n k}\right)$ and $q(x)=\sum_{\beta} b_{\beta} x^{\beta} \in \mathcal{P}\left(\mathrm{M}_{n k}\right)$, then by 4.1),

$$
\langle p, q\rangle=\sum_{\alpha, \beta} a_{\alpha} \overline{b_{\beta}}\left[\frac{\partial^{\left|\alpha^{\prime}\right|}}{\partial x^{\alpha^{\prime}}}\left(x^{\beta}\right)\right](0)=\sum_{\alpha} \alpha^{\prime}!a_{\alpha} \overline{b_{\alpha^{\prime}}},
$$

where for each $\alpha=\left(\alpha_{i j}\right)$,

$$
\alpha^{\prime}!=\prod_{i, j} \alpha_{i j}^{\prime}!.
$$

From this it is easy to see that $\langle\cdot, \cdot\rangle$ is an inner product on $\mathcal{P}\left(\mathrm{M}_{n k}\right)$, and homogeneous polynomials with distinct total degrees are orthogonal. In particular, the sum (3.3) is an orthogonal sum.

Moreover, for $p, q, r \in \mathcal{P}\left(\mathrm{M}_{n k}\right)$,

$$
\langle p q, r\rangle=\{[\partial(q p)](\widetilde{r})\}(0)=\{[\partial(q) \partial(p)](\widetilde{r})\}(0)=\langle q,[\widetilde{\partial(p)](\widetilde{r})}\rangle=\langle q,[\partial(\widetilde{p})](r)\rangle .
$$

Hence the operator $\partial(\widetilde{p})$ is the adjoint to multiplication by $p$. In particular, since $\widetilde{\gamma_{0}}=\gamma_{0}, L=\partial\left(\widetilde{\gamma_{0}}\right)$ is the adjoint to multiplication by $\gamma_{0}$.
5. Generalized Laplacian on Grassmannians. From the discussion in Section 3, we see that $\gamma_{0}$ generates the subalgebra of $\mathrm{O}_{n}$ invariants in $\mathfrak{A}$, that is,

$$
\mathfrak{A}^{\mathrm{O}_{n}}=\mathbb{C}\left[\gamma_{0}\right] .
$$

Multiplication by $\gamma_{0}$ defines an injective $\mathrm{O}_{n}$ map on $\mathfrak{A}$. We shall abuse notation and denote this operator also by $\gamma_{0}$. Since $\gamma_{0}$ is also a GL ${ }_{k}$ eigenvector corresponding to $\operatorname{det}_{k}^{2}, \gamma_{0}$ defines an $\mathrm{O}_{n}$ map

$$
\gamma_{0}: \mathfrak{A}_{m} \rightarrow \mathfrak{A}_{m+2}
$$

so its image $\gamma_{0} \mathfrak{A}_{m}$ in $\mathfrak{A}_{m+2}$ is an $\mathrm{O}_{n}$ submodule isomorphic to $\mathfrak{A}_{m}$.
We recall that the "dual" operator $L=\partial\left(\gamma_{0}\right)$ is the adjoint of $\gamma_{0}$ with respect to the inner product defined in 4.2 . Let

$$
\mathcal{H}=\{p \in \mathfrak{A}: L(p)=0\}
$$

be the space of all "harmonic polynomials" in $\mathfrak{A}$. For each $m \geq 0$, the restriction of $L$ to $\mathfrak{A}_{m}$ also defines an $\mathrm{O}_{n}$ map

$$
L: \mathfrak{A}_{m} \rightarrow \mathfrak{A}_{m-2}
$$

In fact, if $p \in \mathfrak{A}_{\ell}, q \in \mathfrak{A}_{m}$ and $\ell+2 \neq m$, then

$$
0=\left\langle\gamma_{0} p, q\right\rangle=\langle p, L(q)\rangle
$$

This shows that $L(q)$ is orthogonal to $\mathfrak{A}_{\ell}$ for $\ell \neq m-2$, so that $L(q) \in \mathfrak{A}_{m-2}$.
Next, we let

$$
\mathcal{H}_{m}=\left\{f \in \mathfrak{A}_{m}: L(f)=0\right\} .
$$

Then we have

$$
\mathcal{H}=\bigoplus_{m=0}^{\infty} \mathcal{H}_{m}
$$

Proposition 5.1. For each $m \geq 2$,

$$
\mathcal{H}_{m}=\bigoplus_{\substack{\lambda_{1}=m \\ \lambda_{j} \equiv m(\bmod 2)}} \sigma_{n}^{\lambda} \quad \text { and } \quad \mathfrak{A}_{m}=\mathcal{H}_{m} \oplus \gamma_{0} \mathfrak{A}_{m-2}
$$

Proof. Let $p \in \mathfrak{A}_{m-2}$ and $q \in \mathfrak{A}_{m}$. Then $\left\langle\gamma_{0} p, q\right\rangle=\langle p, L(q)\rangle$. Thus if $q \in \mathcal{H}_{m}$, then

$$
\left\langle\gamma_{0} p, q\right\rangle=\langle p, 0\rangle=0
$$

so that

$$
q \in\left(\gamma_{0} \mathfrak{A}_{m}\right)^{\perp}
$$

the orthogonal complement of $\gamma_{0} \mathfrak{A}_{m}$ in $\mathfrak{A}_{m}$. It follows that $\mathcal{H}_{m} \subseteq\left(\gamma_{0} \mathfrak{A}_{m}\right)^{\perp}$.

Let $\sigma_{n}^{\lambda}$ be an $\mathrm{O}_{n}$ representation which occurs in $\mathfrak{A}_{m}$ with $\lambda_{1}=m$. Since $L: \mathfrak{A}_{m} \rightarrow \mathfrak{A}_{m-2}$ and $\sigma_{n}^{\lambda}$ does not occur in $\mathfrak{A}_{m-2}$,

$$
L\left(\sigma_{n}^{\lambda}\right)=0
$$

So $\sigma_{n}^{\lambda} \subseteq \mathcal{H}_{m}$. But these $\mathrm{O}_{n}$ representations $\sigma_{n}^{\lambda}$ together with those in $\gamma_{0} \mathfrak{A}_{m-2}$ have exhausted all the $\mathrm{O}_{n}$ representations in $\mathfrak{A}_{m}$. So the lemma follows.

Corollary 5.2. The space $\mathcal{H}^{N_{n}}$ of $N_{n}$ invariants in $\mathcal{H}$ is the subalgebra of $\mathfrak{A}$ generated by $\left\{\gamma_{1}, \ldots, \gamma_{k}\right\}$.

Theorem 5.3. For $m \geq 2$,

$$
\mathfrak{A}_{m}=\bigoplus_{j=0}^{[m / 2]} \gamma_{0}^{j} \mathcal{H}_{m-2 j} .
$$

Consequently,

$$
\mathfrak{A} \cong \mathcal{H} \otimes \mathfrak{A}^{\mathrm{O}_{n}} .
$$

Proof. The first assertion follows from (3.5) by induction on $m$ together with the observation that $\mathfrak{A}_{0}=\mathcal{H}_{0}$ and $\mathfrak{A}_{1}=\mathcal{H}_{1}$. The second assertion follows from the first.
6. Eigenvalues of $L \gamma_{0}$ in the case $k=2$. For each $m \geq 0$, we have the $\mathrm{O}_{n}$ maps

$$
\gamma_{0}: \mathfrak{A}_{m} \rightarrow \mathfrak{A}_{m+2} \quad \text { and } \quad L: \mathfrak{A}_{m+2} \rightarrow \mathfrak{A}_{m}
$$

By composing these two maps, we obtain the $\mathrm{O}_{n}$ map

$$
L \gamma_{0}: \mathfrak{A}_{m} \rightarrow \mathfrak{A}_{m} .
$$

Since $\mathfrak{A}_{m}$ is multiplicity free as a representation of $\mathrm{O}_{n}$, by Schur's Lemma, it acts by a scalar on each irreducible $\mathrm{O}_{n}$ submodule of $\mathfrak{A}_{m}$. Thus if $\mathfrak{A}_{m, \mathbf{a}}$ is an irreducible $\mathrm{O}_{n}$ submodule of $\mathfrak{A}_{m}$, then there exists a complex number $c(\mathbf{a})$ such that

$$
L \gamma_{0}(v)=c(\mathbf{a}) v, \quad \forall v \in \mathfrak{A}_{m, \mathbf{a}} .
$$

In particular, if we take $v=\gamma^{\mathbf{a}}$, then

$$
\begin{equation*}
L \gamma_{0}\left(\gamma^{\mathbf{a}}\right)=c(\mathbf{a}) \gamma^{\mathbf{a}} . \tag{6.1}
\end{equation*}
$$

It is easy to compute $c(\mathbf{a})$ when $k=1$, which has been discussed in the Introduction. We shall compute the scalar $c(\mathbf{a})$ in the case $k=2$ in this
section. In this case,

$$
\begin{aligned}
L & =\left|\begin{array}{ll}
\Delta_{11} & \Delta_{12} \\
\Delta_{12} & \Delta_{22}
\end{array}\right|, \\
\gamma_{0} & =\left|\begin{array}{ll}
r_{11}^{2} & r_{12}^{2} \\
r_{12}^{2} & r_{22}^{2}
\end{array}\right|, \quad \gamma_{1}=\left|\begin{array}{ccc}
0 & x_{11} & x_{12} \\
x_{11} & r_{11}^{2} & r_{12}^{2} \\
x_{12} & r_{12}^{2} & r_{22}^{2}
\end{array}\right|, \quad \gamma_{2}=\left|\begin{array}{ll}
x_{11} & x_{12} \\
x_{21} & x_{22}
\end{array}\right| .
\end{aligned}
$$

Theorem 6.1. If $k=2$ and $\mathbf{a}=\left(a_{0}, a_{1}, a_{2}\right) \in \mathbb{Z}_{\geq 0}^{3}$, then $c(\mathbf{a})=2\left(a_{0}+1\right)\left(2 a_{0}+2 a_{1}+3\right)\left(2 a_{0}+2 a_{1}+2 a_{2}+n-1\right)\left(2 a_{0}+4 a_{1}+2 a_{2}+n\right)$.

The theorem follows by induction on $a_{0}$ from Lemma 6.2 below. If $T_{1}$ and $T_{2}$ are two linear operators on $\mathfrak{A}$, then the commutator $\left[T_{1}, T_{2}\right]$ is the linear operator

$$
\left[T_{1}, T_{2}\right]=T_{1} T_{2}-T_{2} T_{1} .
$$

Lemma 6.2. If $k=2$ and $\mathbf{a}=\left(a_{0}, a_{1}, a_{2}\right) \in \mathbb{Z}_{\geq 0}^{3}$, then $\left[L, \gamma_{0}\right]\left(\gamma^{\mathbf{a}}\right)=2\left(4 a_{0}+4 a_{1}+2 a_{2}+n\right)$ $\times\left(8 a_{0}^{2}+4 a_{1}^{2}+16 a_{0} a_{1}+8 a_{0} a_{2}+4 a_{1} a_{2}+4 a_{0} n+2 a_{1} n+4 a_{1}+6 a_{2}+3 n-3\right) \gamma^{\mathbf{a}}$.

Before we prove Lemma 6.2, we need to introduce some notations. Let $\mathcal{P D}\left(\mathrm{M}_{n k}\right)$ be the algebra of polynomial-coefficient differential operators on $\mathcal{P}\left(\mathrm{M}_{n k}\right)$. For $1 \leq i, j \leq k$, let

$$
E_{i j}=\sum_{p=1}^{n} x_{p i} \frac{\partial}{\partial x_{p j}}+\delta_{i j} \frac{n}{2} .
$$

Let

$$
\begin{aligned}
\mathfrak{p}^{+} & =\operatorname{Span}\left\{r_{i j}^{2}: 1 \leq i \leq j \leq k\right\}, \\
\mathfrak{p}^{-} & =\operatorname{Span}\left\{\Delta_{i j}: 1 \leq i \leq j \leq k\right\}, \\
\mathfrak{k} & =\operatorname{Span}\left\{E_{i j}: 1 \leq i, j \leq k\right\} .
\end{aligned}
$$

Then $\mathfrak{k}$ is a Lie algebra isomorphic to $\mathfrak{g l}_{k}$, and

$$
\mathfrak{p}^{-} \oplus \mathfrak{k} \oplus \mathfrak{p}^{+}
$$

forms a Lie algebra isomorphic to the symplectic Lie algebra $\mathfrak{s p}_{2 k}$. We shall denote $\mathfrak{k}$ and $\mathfrak{p}^{-} \oplus \mathfrak{k} \oplus \mathfrak{p}^{+}$by $\mathfrak{g l}_{k}$ and $\mathfrak{s p}_{2 k}$ respectively. Then $\mathfrak{s p}_{2 k}$ generates the algebra $\mathcal{P} \mathcal{D}\left(\mathrm{M}_{n k}\right)^{\mathrm{O}_{n}}$ of operators in $\mathcal{P} \mathcal{D}\left(\mathrm{M}_{n k}\right)$ commuting with $\mathrm{O}_{n}$ ([H0], GW]). Consequently, $\mathcal{P D}\left(\mathrm{M}_{n k}\right)^{\mathrm{O}_{n}}$ is a homomorphic image of the universal enveloping algebra $\mathcal{U}\left(\mathfrak{s p}_{2 k}\right)$ of $\mathfrak{s p}_{2 k}$.

Proof of Lemma 6.2. Since both $L$ and $\gamma_{0}$ are elements of $\mathcal{P} \mathcal{D}\left(\mathrm{M}_{n 2}\right)^{\mathrm{O}_{n}}$, so is $\left[L, \gamma_{0}\right]$. We shall express $\left[L, \gamma_{0}\right]$ in the form

$$
\left[L, \gamma_{0}\right]=X+Y E_{12}+Z E_{21}
$$

where $X, Y, Z$ are elements $\mathcal{P} \mathcal{D}\left(\mathrm{M}_{n 2}\right)^{\mathrm{O}_{n}}$ such that the expression for $X$ does not involve the elements $E_{12}$ and $E_{21}$. Since $\mathrm{GL}_{2}$ acts on $\gamma^{\mathbf{a}}$ by a determinant character, the Lie algebra $\mathfrak{s l}_{2}$ of $\mathrm{SL}_{2}$ will annihilate it. In particular,

$$
E_{12}\left(\gamma^{\mathbf{a}}\right)=E_{21}\left(\gamma^{\mathbf{a}}\right)=0
$$

Consequently,

$$
\begin{equation*}
\left[L, \gamma_{0}\right]\left(\gamma^{\mathbf{a}}\right)=X\left(\gamma^{\mathbf{a}}\right)+Y E_{12}\left(\gamma^{\mathbf{a}}\right)+Z E_{21}\left(\gamma^{\mathbf{a}}\right)=X\left(\gamma^{\mathbf{a}}\right) \tag{6.2}
\end{equation*}
$$

Thus it suffices to compute $X\left(\gamma^{\mathbf{a}}\right)$.
We now determine the elements $X, Y$ and $Z$. We have

$$
\begin{align*}
{\left[L, \gamma_{0}\right] } & =\left[\left[\left.\begin{array}{ll}
\Delta_{11} & \Delta_{12} \\
\Delta_{12} & \Delta_{22}
\end{array} \right\rvert\,, \gamma_{0}\right]=\left|\begin{array}{ll}
{\left[\Delta_{11}, \gamma_{0}\right]} & \Delta_{12} \\
{\left[\Delta_{12}, \gamma_{0}\right]} & \Delta_{22}
\end{array}\right|+\left|\begin{array}{cc}
\Delta_{11} & {\left[\Delta_{12}, \gamma_{0}\right]} \\
\Delta_{12} & {\left[\Delta_{22}, \gamma_{0}\right]}
\end{array}\right|\right.  \tag{6.3}\\
& =\left[\Delta_{11}, \gamma_{0}\right] \Delta_{22}-\left[\Delta_{12}, \gamma_{0}\right] \Delta_{12}+\Delta_{11}\left[\Delta_{22}, \gamma_{0}\right]-\Delta_{12}\left[\Delta_{12}, \gamma_{0}\right]
\end{align*}
$$

We will need the following commutation relations in $\mathfrak{s p}_{4}$ in our computations:

$$
\begin{aligned}
{\left[\Delta_{a b}, r_{c d}^{2}\right] } & =\delta_{b c} E_{d a}+\delta_{a c} E_{d b}+\delta_{b d} E_{c a}+\delta_{a d} E_{c b} \\
{\left[E_{a b}, r_{c d}^{2}\right] } & =\delta_{b c} r_{a d}^{2}+\delta_{b d} r_{a c}^{2} \\
{\left[E_{a b}, \Delta_{c d}\right] } & =-\delta_{a c} \Delta_{b d}-\delta_{a d} \Delta_{c b}
\end{aligned}
$$

Using these formulas, we obtain

$$
\begin{aligned}
& {\left[\Delta_{11}, \gamma_{0}\right]=-2 r_{22}^{2}+4 r_{22}^{2} E_{11}-4 r_{12}^{2} E_{21}} \\
& {\left[\Delta_{22}, \gamma_{0}\right]=-2 r_{11}^{2}+4 r_{11}^{2} E_{22}-4 r_{12}^{2} E_{12}} \\
& {\left[\Delta_{12}, \gamma_{0}\right]=2 r_{12}^{2}-2 r_{12}^{2}\left(E_{11}+E_{22}\right)+2 r_{11}^{2} E_{21}+2 r_{22}^{2} E_{12}}
\end{aligned}
$$

Substituting these expressions into 6.3 and simplifying, we obtain

$$
\begin{aligned}
X= & 4\left(r_{22}^{2} \Delta_{22}+r_{12}^{2} \Delta_{12}\right) E_{11}+4\left(r_{11}^{2} \Delta_{11}+r_{12}^{2} \Delta_{12}\right) E_{22} \\
& -14 E_{11}+2 E_{22}+16 E_{11} E_{22}+2\left(E_{11}+E_{22}\right)^{2} \\
Y= & -4 r_{22}^{2} \Delta_{12}-4 r_{12}^{2} \Delta_{11}-16 E_{21} \\
Z= & -4 r_{12}^{2} \Delta_{22}-4 r_{11}^{2} \Delta_{12}
\end{aligned}
$$

It follows from $(6.2)$ that

$$
\begin{aligned}
{\left[L, \gamma_{0}\right]\left(\gamma^{\mathbf{a}}\right)=} & X\left(\gamma^{\mathbf{a}}\right) \\
= & 4\left(r_{22}^{2} \Delta_{22}+r_{12}^{2} \Delta_{12}\right) E_{11}\left(\gamma^{\mathbf{a}}\right)+4\left(r_{11}^{2} \Delta_{11}+r_{12}^{2} \Delta_{12}\right) E_{22}\left(\gamma^{\mathbf{a}}\right) \\
& +\left\{-14 E_{11}+2 E_{22}+16 E_{11} E_{22}+2\left(E_{11}+E_{22}\right)^{2}\right\}\left(\gamma^{\mathbf{a}}\right)
\end{aligned}
$$

Now

$$
E_{11}\left(\gamma^{\mathbf{a}}\right)=E_{22}\left(\gamma^{\mathbf{a}}\right)=\lambda \gamma^{\mathbf{a}} \quad \text { where } \quad \lambda=2 a_{0}+2 a_{1}+a_{2}+n / 2
$$

So

$$
\begin{align*}
{\left[L, \gamma_{0}\right]\left(\gamma^{\mathbf{a}}\right) } & =4 \lambda T\left(\gamma^{\mathbf{a}}\right)+\left\{-14 \lambda+2 \lambda+16 \lambda^{2}+2(2 \lambda)^{2}\right\} \gamma^{\mathbf{a}}  \tag{6.4}\\
& =4 \lambda T\left(\gamma^{\mathbf{a}}\right)+12 \lambda(2 \lambda-1) \gamma^{\mathbf{a}}
\end{align*}
$$

where $T$ is the operator

$$
T=r_{11}^{2} \Delta_{11}+r_{22}^{2} \Delta_{22}+2 r_{12}^{2} \Delta_{12}
$$

We note that the operator $T$ corresponds to the trivial representation of $\mathrm{GL}_{2}$ in the tensor product $S^{2}\left(\mathbb{C}^{2}\right) \otimes S^{2}\left(\mathbb{C}^{2 *}\right)$. So it will act as a scalar on $\gamma^{\mathbf{a}}$. We now compute this scalar.

Routine calculations show

$$
\begin{aligned}
\Delta_{11}\left(\gamma^{\mathbf{a}}\right)= & 2 a_{0}\left(2 a_{0}+4 a_{1}+2 a_{2}+n-3\right) r_{22}^{2} \gamma_{0}^{a_{0}-1} \gamma_{1}^{a_{1}} \gamma_{2}^{a_{2}} \\
& -2 a_{1}\left(2 a_{1}+2 a_{2}+n-4\right) x_{12}^{2} \gamma_{0}^{a_{0}} \gamma_{1}^{a_{1}-1} \gamma_{2}^{a_{2}}, \\
\Delta_{22}\left(\gamma^{\mathbf{a}}\right)= & 2 a_{0}\left(2 a_{0}+4 a_{1}+2 a_{2}+n-3\right) r_{11}^{2} \gamma_{0}^{a_{0}-1} \gamma_{1}^{a_{1}} \gamma^{a_{2}} \\
& -2 a_{1}\left(2 a_{1}+2 a_{2}+n-4\right) x_{11}^{2} \gamma_{0}^{a_{0}} \gamma_{1}^{a_{1}-1} \gamma_{2}^{a_{2}}, \\
\Delta_{12}\left(\gamma^{\mathbf{a}}\right)= & -2 a_{0}\left(2 a_{0}+4 a_{1}+2 a_{2}+n-3\right) r_{12}^{2} \gamma_{0}^{a_{0}-1} \gamma_{1}^{a_{1}} \gamma_{2}^{a_{2}} \\
& +2 a_{1}\left(2 a_{1}+2 a_{2}+n-4\right) x_{11} x_{12} \gamma_{0}^{a_{0}} \gamma_{1}^{a_{1}-1} \gamma_{2}^{a_{2}} .
\end{aligned}
$$

Using these formulas, we obtain

$$
\begin{aligned}
T\left(\gamma^{\mathbf{a}}\right)= & 2 a_{0}\left(2 a_{0}+4 a_{1}+2 a_{2}+n-3\right)\left\{r_{11}^{2} r_{22}^{2}+r_{22}^{2} r_{11}^{2}-2\left(r_{12}^{2}\right)^{2}\right\} \gamma_{0}^{a_{0}-1} \gamma_{1}^{a_{1}} \gamma_{2}^{a_{2}} \\
& +2 k\left(2 a_{1}+2 a_{2}+n-4\right)\left\{-x_{12}^{2} r_{11}^{2}-x_{11}^{2} r_{22}^{2}+2 x_{11} x_{12} r_{12}^{2}\right\} \gamma_{0}^{a_{0}} \gamma_{1}^{k-1} \gamma_{2}^{a_{2}} \\
= & 2 a_{0}\left(2 a_{0}+4 a_{1}+2 a_{2}+n-3\right)\left\{2 \gamma_{0}\right\} \gamma_{0}^{a_{0}-1} \gamma_{1}^{k} \gamma_{2}^{a_{2}} \\
& +2 a_{1}\left(2 a_{1}+2 a_{2}+n-4\right)\left\{\gamma_{1}\right\} \gamma_{0}^{a_{0}} \gamma_{1}^{k-1} \gamma_{2}^{a_{2}} \\
=\{ & \left.4 a_{0}\left(2 a_{0}+4 a_{1}+2 a_{2}+n-3\right)+2 a_{1}\left(2 a_{1}+2 a_{2}+n-4\right)\right\} \gamma^{\mathbf{a}} .
\end{aligned}
$$

This together with (6.4) gives

$$
\begin{aligned}
{\left[L, \gamma_{0}\right]\left(\gamma^{\mathbf{a}}\right)=} & 4 \lambda T\left(\gamma^{\mathbf{a}}\right)+12 \lambda(2 \lambda-1) \gamma^{\mathbf{a}} \\
= & 4 \lambda\left\{4 a_{0}\left(2 a_{0}+4 a_{1}+2 a_{2}+n-3\right)+2 a_{1}\left(2 a_{1}+2 a_{2}+n-4\right)\right\} \gamma^{\mathbf{a}} \\
& +12 \lambda(2 \lambda-1) \gamma^{\mathbf{a}} \\
= & 2\left(4 a_{0}+4 a_{1}+2 a_{2}+n\right)\left(8 a_{0}^{2}+4 a_{1}^{2}+16 a_{0} a_{1}+8 a_{0} a_{2}\right. \\
& \left.+4 a_{1} a_{2}+4 a_{0} n+2 a_{1} n+4 a_{1}+6 a_{2}+3 n-3\right) \gamma^{\mathbf{a}}
\end{aligned}
$$

This proves the lemma.
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[^1]:    $\left.{ }^{1}{ }^{1}\right)$ If $G$ is a reductive algebraic group, then a flag manifold for $G$ is a coset space $G / P$, where $P$ is a parabolic subgroup of $G$. The homogeneous coordinate ring for $G / P$ is the ring $\mathcal{R}\left(G / P^{(2)}\right)$ of regular functions on the variety $G / P^{(2)}$, where $P^{(2)}$ is the commutator subgroup of $P$. The variety $G / P$ is a projective variety, and therefore has no nonconstant regular functions. However, the variety $G / P^{(2)}$ is a torus bundle over $G / P$, and is quasiaffine, so it has a large collection of regular functions, which can be thought of as (sums of) sections of appropriate line bundles over $G / P$.

