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# ASYMPTOTICS OF PARABOLIC EQUATIONS WITH POSSIBLE BLOW-UP 

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#### Abstract

We describe the long-time behaviour of solutions of parabolic equations in the case when some solutions may blow up in a finite or infinite time. This is done by providing a maximal compact invariant set attracting any initial data for which the corresponding solution does not blow up. The abstract result is applied to the FrankKamenetskii equation and the $N$-dimensional Navier-Stokes system with small external force.


1. Introduction. In this paper we study the asymptotic behaviour of parabolic equations when some solutions may blow up in a finite or infinite time. We consider $X^{\alpha}$ solutions as in [C-D1] and earlier in [HE] with the modification given in $[\mathrm{MI}]$. We make use of the theory of semilinear abstract parabolic equations given e.g. in [HE], [HA] or [CZ].

The situation that for some initial data the corresponding solution blows up often occurs in physical applications. For a detailed mathematical description of this phenomenon we refer the reader to $[G-V]$. Here we only mention a particular problem with a parameter $\lambda>0$,

$$
\left\{\begin{array}{l}
u_{t}=\Delta u+\lambda e^{u}, \quad t>0 \\
u(0)=u_{0}
\end{array}\right.
$$

which comes from combustion theory and is known under the name of solid fuel ignition model, exponential reaction-diffusion equation or FrankKamenetskii equation, the latter name being used from now on in this paper. Among many other results concerning this problem (see e.g. [GE], [J-L], [B-E]) it was shown in [FU] that for special initial data $u_{0}$ the corresponding solution blows up. On the other hand, there also exist parabolic problems for which global solvability for all initial data remains unknown. A typical example is the famous $N$-dimensional Navier-Stokes system, $N \geq 3$, that has now been investigated for more than a century.

[^0]The above mentioned specific circumstances make it impossible to describe the asymptotics by using the notion of a global attractor in the large space of initial data as in [HA], [LA], [C-D1]. Instead, we consider the maximal compact invariant set attracting each initial condition for which the corresponding solution does not blow up. We first describe all these ideas abstractly in Section 2 considering a Cauchy problem for a semilinear sectorial equation in a Banach space. Next in Section 3 we discuss two particular examples: the one-dimensional version of the Frank-Kamenetskii equation and the $N$-dimensional Navier-Stokes system with small external force.
2. Abstract parabolic problems with possible blow-up. Consider an abstract parabolic problem

$$
\left\{\begin{array}{l}
u_{t}+A u=F(u), \quad t>0  \tag{2.1}\\
u(0)=u_{0}
\end{array}\right.
$$

where $X$ is a Banach space and $A: X \supset \operatorname{dom}(A) \rightarrow X$ is a positive sectorial operator with compact resolvent. Moreover, assume that $F: X^{\alpha} \rightarrow X(\alpha \in$ $[0,1)$ is fixed from now on) is Lipschitz continuous on bounded subsets of the fractional power space $X^{\alpha}=\operatorname{dom}\left(A^{\alpha}\right)(c f .[H E],[A M])$.

Under the above assumptions the theory of semilinear parabolic equations given e.g. in [HE] ensures for each $u_{0} \in X^{\alpha}$ the existence of a unique local $X^{\alpha}$ solution defined on a maximal interval of existence $\left[0, \tau_{u_{0}}\right)$, where $\tau_{u_{0}} \leq \infty$. Thus we know that

$$
u \in C\left(\left[0, \tau_{u_{0}}\right), X^{\alpha}\right) \cap C\left(\left(0, \tau_{u_{0}}\right), X^{1}\right) \cap C^{1}\left(\left(0, \tau_{u_{0}}\right), X\right)
$$

and (2.1) is satisfied in $X$. Moreover, we have either $\tau_{u_{0}}=\infty$, or $\tau_{u_{0}}<\infty$ and

$$
\limsup _{t}\left\|u\left(t, u_{0}\right)\right\|_{X^{\alpha}}=\infty
$$

Since the problem (2.1) is autonomous, the uniqueness of solutions allows us to construct a local semiflow on $X^{\alpha}$. We thus have

$$
u\left(s, u\left(t, u_{0}\right)\right)=u\left(s+t, u_{0}\right), \quad u_{0} \in X^{\alpha}, s, t \geq 0, s+t<\tau_{u_{0}}
$$

and the solutions of (2.1) are continuous with respect to their initial data on compact time intervals (cf. [HE, Theorem 3.4.1] or [C-D1, Proposition 2.3.2]).

For our further investigations let us define a metric space

$$
\begin{equation*}
V=\left\{u_{0} \in X^{\alpha}: \sup _{t \in\left[0, \tau_{u_{0}}\right)}\left\|u\left(t, u_{0}\right)\right\|_{X^{\alpha}}<\infty\right\} \tag{2.2}
\end{equation*}
$$

and assume that $V$ is nonvoid.
Consider a $C^{0}$-semigroup $T(t): V \rightarrow V$ defined by

$$
\begin{equation*}
T(t) u_{0}=u\left(t, u_{0}\right), \quad t \geq 0, u_{0} \in V \tag{2.3}
\end{equation*}
$$

Note that we do not know in advance whether $V$ is a closed subset of $X^{\alpha}$. Thus it is unknown if $V$, which is the natural phase space for the problem (2.1), is a complete metric space or not. Therefore the assumption of the compact resolvent does not necessarily imply the compactness of the semigroup $\{T(t): t \geq 0\}$ on $V$. Nevertheless, we shall show below that this semigroup is asymptotically smooth in the sense of [HA].

For $B \subset V$ we denote its orbit by $\gamma^{+}(B)=\bigcup_{t \geq 0} T(t) B$, while its $\omega$-limit set is given by

$$
\omega(B)=\bigcap_{t \geq 0} \operatorname{cl}_{X^{\alpha}} T(t) \gamma^{+}(B)
$$

We also abbreviate $\gamma^{+}(\{v\})$ and $\omega(\{v\})$ to $\gamma^{+}(v)$ and $\omega(v)$, respectively.
REMARK 2.1. If $B \subset V$ and $\gamma^{+}(B)$ is bounded, then $\mathrm{cl}_{X^{\alpha}} \gamma^{+}(B) \subset V$ and $T(t) B$ is precompact in $X^{\alpha}$ for any $t>0$.

Proof. Let $v_{0} \in \operatorname{cl}_{X^{\alpha}} \gamma^{+}(B)$. Then there exist $v_{n} \in B$ and $t_{n} \geq 0$ such that $T\left(t_{n}\right) v_{n} \rightarrow v_{0}$ in $X^{\alpha}$. Since

$$
\forall_{s \geq 0} \quad\left\|u\left(s, T\left(t_{n}\right) v_{n}\right)\right\|_{X^{\alpha}}=\left\|u\left(s+t_{n}, v_{n}\right)\right\|_{X^{\alpha}} \leq R_{\gamma^{+}(B)}
$$

the norm $\left\|u\left(s, v_{0}\right)\right\|_{X^{\alpha}}$ cannot blow up so that $v_{0} \in V$. Also, if $\gamma^{+}(B)$ is bounded in $X^{\alpha}$, then $T(t) B$ with $t>0$ is in fact bounded in $X^{\alpha+\varepsilon}$, which, via compactness of the embeddings $X^{\beta} \subset X^{\alpha}, \beta>\alpha$, ensures that it is precompact in $X^{\alpha}$.

Proposition 2.2. The $C^{0}$-semigroup $\{T(t): t \geq 0\}$ on $V$ is asymptotically smooth, i.e. each nonvoid closed (in $V$ ) bounded and positively invariant set $W \subset V$ contains a nonvoid compact subset $\omega(W)$ which attracts $W$.

Proof. Since $\gamma^{+}(W) \subset W$, we infer from Remark 2.1 that $\mathrm{cl}_{X^{\alpha}} T(1) \gamma^{+}(W)$ is a compact subset of $V$. Thus

$$
\omega(W)=\bigcap_{t \geq 0} \operatorname{cl}_{X^{\alpha}} T(t) \gamma^{+}(W) \subset \operatorname{cl}_{X^{\alpha}} T(1) \gamma^{+}(W) \subset V
$$

is compact and nonvoid as the intersection of a centered family of closed sets in a compact space. Also $\omega(W) \subset W$, since by our assumptions on $W$ and Remark 2.1 we have

$$
\omega(W) \subset \operatorname{cl}_{X^{\alpha}} W \cap V=\operatorname{cl}_{V} W=W
$$

We now prove that $\omega(W)$ attracts $W$, that is,

$$
d(T(t) W, \omega(W)) \underset{t \rightarrow \infty}{\longrightarrow} 0
$$

where $d$ denotes the Hausdorff semidistance. Contrary to our claim suppose that

$$
\begin{equation*}
\exists_{\varepsilon>0} \exists_{t_{n} \rightarrow \infty} \forall_{n \in \mathbb{N}} \exists_{x_{n} \in W} \forall_{v \in \omega(W)} \quad\left\|T\left(t_{n}\right) x_{n}-v\right\|_{X^{\alpha}}>\varepsilon \tag{2.4}
\end{equation*}
$$

Since almost all elements of the sequence $\left\{T\left(t_{n}\right) x_{n}\right\}$ belong to the compact set $\mathrm{cl}_{X^{\alpha}} T(1) \gamma^{+}(W)$ it contains a subsequence convergent to an element of $\omega(W)$, which contradicts (2.4).

Remark 2.3. It is worth noting that Remark 2.1 implies immediately that if $u_{0} \in V$, then $\mathrm{cl}_{X^{\alpha}} \gamma^{+}\left(u_{0}\right)$ is a compact positively invariant subset of $V$ and the $\omega$-limit set $\omega\left(u_{0}\right)$ is a nonvoid compact connected and invariant subset of $V$ which attracts $u_{0}$ (cf. [LA, Theorem 2.1]). More generally, since by Proposition 2.2 the semigroup is asymptotically smooth, we infer that for $\emptyset \neq B \subset V$ with $\gamma^{+}(B)$ bounded, the $\omega$-limit set $\omega(B)$ is a nonvoid compact and invariant subset of $V$ attracting $B$ (see [C-D1, Proposition 1.1.1]).

Under an additional assumption of the point dissipativeness of the semigroup $\{T(t): t \geq 0\}$ we now prove the existence of a nonvoid compact and invariant set which attracts each point of $V$. Indeed, if $B_{0}$ is a nonvoid bounded subset of $V$ attracting points of $V$, then any $\varepsilon$-neighbourhood $\mathcal{N}\left(B_{0}\right)=\bigcup_{u \in B_{0}} B(u, \varepsilon)$ of $B_{0}$ in $V$ absorbs each point of $V$. Consequently,

$$
\widetilde{B}_{0}=\left\{v \in \mathcal{N}\left(B_{0}\right): \gamma^{+}(v) \subset \mathcal{N}\left(B_{0}\right)\right\}
$$

is a nonvoid bounded positively invariant subset of $V$ absorbing each point of $V$. From this and Remark 2.3 it follows that $\omega\left(\widetilde{B}_{0}\right)$ is a nonvoid compact invariant subset of $V$ attracting each point of $V$.

The required dissipativeness property can be easily controlled if there exists a Lyapunov function. Recall that by a Lyapunov function on $V$ we mean a continuous function $\mathcal{L}: V \rightarrow \mathbb{R}$ such that for any $u_{0} \in V$,
(i) the function $t \mapsto \mathcal{L}\left(u\left(t, u_{0}\right)\right)$ is nonincreasing in $(0, \infty)$,
(ii) if $\mathcal{L}\left(u\left(\cdot, u_{0}\right)\right) \equiv \mathcal{L}\left(u_{0}\right)$, then $u_{0} \in \mathcal{E}$,
where $\mathcal{E}$ denotes the set of all stationary solutions of (2.1). Recall also that the existence of a Lyapunov function on $V$ implies that $\omega\left(u_{0}\right) \subset \mathcal{E}$ for each $u_{0} \in V$. Therefore, if there exists a Lyapunov function on the metric space $V$ given in (2.2), then the set $\mathcal{E}$ is nonvoid, and if it is bounded, it is also compact and attracts each point of $V$. Thus we get

Corollary 2.4. Suppose that $\{T(t): t \geq 0\}$ is defined on $V$ given in (2.2). Assume further that $\{T(t): t \geq 0\}$ is point dissipative (for example, there exists a Lyapunov function on $V$ and $\mathcal{E}$ is bounded). Then for any $u_{0} \in X^{\alpha}$ the $X^{\alpha}$ solution $u\left(\cdot, u_{0}\right)$ of (2.1) either blows up (in a finite or infinite time) or stays bounded and approaches a nonvoid compact and invariant set.

We have shown above that the semigroup is asymptotically smooth and we have also stated natural conditions ensuring its point dissipativeness. However, these two properties do not guarantee the existence of a compact global attractor in $V$. It would exist if we knew the semigroup on $V$ was
compact (cf. [C-D1, Corollary 1.1.6]), which, as we have already observed, may not be the case, or if the orbits of bounded sets were bounded (cf. [C-D1, Theorem 1.1.2]). Unfortunately, the latter condition may be difficult to check in specific examples. As will be seen in Section 3 it is much easier to examine the boundedness of the set of all (hypothetical) bounded complete orbits of points.

Following [LA], we recall that a complete trajectory of a point $v \in V$ for the semigroup $\{T(t): t \geq 0\}$ is the curve $\phi: \mathbb{R} \rightarrow V$ satisfying the following conditions:
(i) $\phi(0)=v$,
(ii) $T(t) \phi(s)=\phi(s+t), s \in \mathbb{R}, t \geq 0$.

If $\mathcal{S}$ denotes the set of all points in $V$ for which there exists at least one bounded complete trajectory for the semigroup $\{T(t): t \geq 0\}$, then $T(t) \mathcal{S}=$ $\mathcal{S}$ for each $t>0$. Also the following result holds.

Theorem 2.5. Suppose that the semigroup $\{T(t): t \geq 0\}$ is defined by (2.3) on the nonvoid metric space $V$ given in (2.2). Then $\mathcal{S}$ is a nonvoid invariant subset of $V$ which attracts each subset of $V$ with bounded orbit. If $\mathcal{S}$ is bounded, then it is a compact and maximal bounded invariant subset of $V$. If, additionally, the orbits of bounded subsets of $V$ are bounded, then $\mathcal{S}$ is a compact global attractor in $V$.

Proof. As a consequence of Remark 2.3, $\mathcal{S}$ is nonvoid whenever $V$ is nonvoid. Moreover, $\omega(B) \subset \mathcal{S}$ for any $\emptyset \neq B \subset V$ such that $\gamma^{+}(B)$ is bounded. It is next sufficient to note that if $\mathcal{S}$ is bounded, then-since it is also invariant- $\mathrm{cl}_{X^{\alpha}} \mathcal{S}$ is a compact subset of $V$ (see Remark 2.1). The proof is thus complete.

Corollary 2.6. Suppose $A$ is a sectorial operator in a Banach space $X$ and $A$ has compact resolvent. Assume that $F: X^{\alpha} \rightarrow X$, with $\alpha \in[0,1)$ fixed, is Lipschitz continuous on bounded subsets of $X^{\alpha}$. Let all (hypothetical) bounded complete orbits of points be uniformly bounded in $X^{\alpha}$. Then for any $u_{0} \in X^{\alpha}$ the $X^{\alpha}$ solution $u\left(\cdot, u_{0}\right)$ of (2.1) either blows up (in a finite or infinite time) or stays bounded and approaches a maximal compact invariant set.

We recall that similarly to the case of an $\omega$-limit set, the $\alpha$-limit set of $u_{0} \in \mathcal{S}$ along a bounded complete trajectory $\phi$ of $u_{0}$,

$$
\alpha_{\phi}\left(u_{0}\right)=\bigcap_{t \leq 0} \operatorname{cl}_{X^{\alpha}} \bigcup_{s \leq t}\{\phi(s)\},
$$

is a nonvoid compact subset of $V$. If, in addition, there exists a Lyapunov function on $V$, then $\alpha_{\phi}\left(u_{0}\right) \subset \mathcal{E}$. In the latter case abstract conditions for the boundedness of $\mathcal{S}$ can be formulated.

Proposition 2.7. Assume that there exists a Lyapunov function $\mathcal{L}$ on $V$. Then the following conditions are equivalent:
(a) $\mathcal{S}$ is a bounded subset of $V$,
(b) $\mathcal{E}$ is a bounded subset of $V$ and one of the equivalent conditions holds:
(i) if $v_{n} \in \mathcal{S}$ and $\left\|v_{n}\right\|_{X^{\alpha}} \rightarrow \infty$, then $\left|\mathcal{L}\left(v_{n}\right)\right| \rightarrow \infty$ (cf. [HA, Definition 3.8.1]),
(ii) if $B \subset \mathcal{S}$ and $\mathcal{L}(B)$ is a bounded subset of $\mathbb{R}$, then $B$ is bounded.

Proof. We shall prove that (b) implies (a). Therefore we assume (b) and suppose contrary to our claim that there exist $v_{n}=\phi_{n}\left(t_{n}\right)$, where $t_{n} \in \mathbb{R}$, $\phi_{n}$ are bounded complete trajectories, say of some $u_{n}$, and $\left\|v_{n}\right\|_{X^{\alpha}} \rightarrow \infty$ as $n \rightarrow \infty$. We choose $\alpha_{n} \in \alpha_{\phi_{n}}\left(u_{n}\right)$ and $\omega_{n} \in \omega\left(u_{n}\right)$, both in $\mathcal{E}$. Since the Lyapunov function $\mathcal{L}$ is nonincreasing along each complete trajectory, we see that

$$
\max \{\mathcal{L}(e): e \in \mathcal{E}\} \geq \mathcal{L}\left(\alpha_{n}\right) \geq \mathcal{L}\left(v_{n}\right) \geq \mathcal{L}\left(\omega_{n}\right) \geq \min \{\mathcal{L}(e): e \in \mathcal{E}\}
$$

where the maximum and minimum exist due to the compactness of $\mathcal{E}$. But this is impossible, because of (i). This shows that $\mathcal{S}$ is bounded.

## 3. Examples

Example 3.1. Consider the Dirichlet problem for the Frank-Kamenetskii equation

$$
\left\{\begin{array}{l}
u_{t}=\Delta u+\lambda e^{u}, \quad x \in \Omega, t>0  \tag{3.1}\\
u(x, t)=0, \quad x \in \partial \Omega, t>0 \\
u(x, 0)=u_{0}(x), \quad x \in \Omega
\end{array}\right.
$$

where $\lambda>0$ is a parameter and $\Omega=B(0,1) \subset \mathbb{R}^{N}$. This problem occurs in models of thermal explosions, especially in the description of thermal self-ignition of a chemically active mixture contained in some vessel. We refer the reader to $[\mathrm{FK}],[\mathrm{GE}, \S 15]$ and $[\mathrm{B}-\mathrm{E}]$ for more details.

Rewriting the problem (3.1) in an abstract setting we consider

$$
\left\{\begin{array}{l}
u_{t}+A u=F(u), \quad t>0  \tag{3.2}\\
u(0)=u_{0}
\end{array}\right.
$$

in the Hilbert space $X=L^{2}(\Omega)$, where $A=-\Delta_{D}: X \supset \operatorname{dom}(A) \rightarrow X$ with $\operatorname{dom}(A)=H^{2}(\Omega) \cap H_{0}^{1}(\Omega)$. It is well known that this operator is a positive sectorial operator in $X$ with compact resolvent. Although the Frank-Kamenetskii equation is also interesting and complex for dimensions $3 \leq N \leq 9$ (especially because of an infinite number of stationary solutions for $\lambda=2(N-2)$, see e.g. [J-L], [B-E, Theorem 2.19], [F-P], [N-S]), we restrict our attention to $N=1$. Nevertheless, we still use the general notation.

Fix $3 / 4<\alpha<1$ so that $X^{\alpha} \subset C^{1}(\bar{\Omega})$. Evidently, if $u \in X^{\alpha}$, then $u \in$ $C^{1}(\bar{\Omega}), F(u)=\lambda e^{u} \in C^{1}(\bar{\Omega})$, and $F: X^{\alpha} \rightarrow X$ is Lipschitz continuous on bounded sets. Consequently, (3.1) generates a local semiflow of $X^{\alpha}$ solutions. Denoting then by $V$ the set of all initial data for which the solution stays bounded in $X^{\alpha}$ we see that

$$
\begin{equation*}
\mathcal{L}(u)=\frac{1}{2} \int_{\Omega}|\nabla u|^{2} d x-\int_{\Omega} \lambda e^{u} d x \tag{3.3}
\end{equation*}
$$

is a Lyapunov function on $V$.
Stationary solutions and their Morse indices. If $w \in \operatorname{dom}(A)$ is a stationary solution, then $w$ satisfies

$$
\left\{\begin{array}{l}
-\Delta w=\lambda e^{w}, \quad x \in \Omega  \tag{3.4}\\
w(x)=0, \quad x \in \partial \Omega
\end{array}\right.
$$

which is also known in the literature under the names of the Emden-Fowler equation or Gelfand problem. From the regularity theory of elliptic operators (cf. [TR, Theorem 5.4.1]) it follows that the stationary solutions are smooth and in particular they belong to $C^{2}([-1,1])$. Also, all solutions of (3.4) are positive, and thus radially symmetric by the result of $[\mathrm{G}-\mathrm{N}-\mathrm{N}]$. We recall (see [B-E, Theorem 2.19]) that there exists $\lambda^{*}>0, \lambda^{*} \approx 0.878$, such that
(a) for each $\lambda \in\left(0, \lambda^{*}\right)$ there are two solutions,
(b) for $\lambda=\lambda^{*}$ there is a unique solution,
(c) for $\lambda>\lambda^{*}$ there are no solutions.

Moreover, each solution $w$ has to satisfy

$$
\begin{equation*}
w(x)=w(0)-2 \ln \cosh \left(\frac{1}{2} \sqrt{2 \lambda e^{w(0)}} x\right), \quad-1 \leq x \leq 1 \tag{3.5}
\end{equation*}
$$

If $\lambda \in\left(0, \lambda^{*}\right)$, then from $[\mathrm{FU}]$ we infer that there exists the minimal solution. Let us denote it by $w^{+}$and the maximal solution by $w^{-}$. We know that there exists $\gamma>0$ such that

$$
\gamma \varrho(x) \leq w^{-}(x)-w^{+}(x), \quad x \in \bar{\Omega}
$$

where $\varrho(x)$ is the distance from $x$ to $\partial \Omega$. Consequently, the curve shown in Figure 1 describes the set of solution curves

$$
\{(\lambda, w) \in(0, \infty) \times C(\bar{\Omega}):(\lambda, w) \text { satisfies }(3.4)\}
$$

We remark that, as shown in [FU, Theorem 6] (see also the refinement in [FI, Remark 2.5]), if $\lambda \in\left(0, \lambda^{*}\right), w^{-}(x) \leq u_{0}(x)$ for $x \in \bar{\Omega}$ and $u_{0} \not \equiv w^{-}$, then the solution $u$ of (3.1) blows up in a finite time.

Consider now the linearization of (3.4) at $w$,

$$
\left\{\begin{array}{l}
\Delta v+\lambda e^{w} v=0, \quad x \in \Omega  \tag{3.6}\\
v(x)=0, \quad x \in \partial \Omega
\end{array}\right.
$$



Fig. 1. Dependence between $w(0)$ and $\lambda$
and observe that $\sigma\left(-\Delta-\lambda e^{w}\right)$ consists only of real eigenvalues. We say that $w$ is a hyperbolic stationary solution if $0 \notin \sigma\left(-\Delta-\lambda e^{w}\right)$. Furthermore, the number of negative eigenvalues of $-\Delta-\lambda e^{w}$ is called the Morse index $\operatorname{ind}(w)$ of the stationary solution $w$. Although $0 \in \sigma\left(-\Delta-\lambda^{*} e^{w}\right)$, it is known that if $\lambda \in\left(0, \lambda^{*}\right)$, then $w^{-}$and $w^{+}$are both hyperbolic stationary solutions. Additionally, we have $\operatorname{ind}\left(w^{+}\right)=0$ and $\operatorname{ind}\left(w^{-}\right)=1$ (for details see [C-R, Proposition 2.15] and [N-S, Section 2]).

Unstable manifold of $w^{-}$and description of $\mathcal{S}$. Let $\lambda \in\left(0, \lambda^{*}\right)$. Note that $w^{-}$is a hyperbolic fixed point in the sense of [C-C-H, p. 357]. Then the unstable set $W^{\mathrm{u}}\left(w^{-}\right)$is a $C^{1}$ submanifold of $X^{\alpha}$ with

$$
\operatorname{dim} W^{\mathrm{u}}\left(w^{-}\right)=\operatorname{ind}\left(w^{-}\right)=1
$$

(see [HE, Theorem 6.1.9], [C-C-H, Appendix C]) and by [B-F1, Theorem 2.1] we have

$$
\begin{equation*}
\forall_{v \in W^{\mathrm{u}}\left(w^{-}\right)} \quad z\left(v-w^{-}\right)<\operatorname{dim} W^{\mathrm{u}}\left(w^{-}\right)=1 \tag{3.7}
\end{equation*}
$$

where $z(g)$ denotes the number of sign changes of a continuous function $g$.
The existence of a Lyapunov function excludes the existence of (nonconstant) homoclinic orbits. Thus we restrict our attention to heteroclinic orbits. Let $\phi$ be a (hypothetical) complete trajectory of $u_{0} \in X^{\alpha}$ such that $\phi(t) \rightarrow w^{-}$as $t \rightarrow-\infty$ and $\phi(t) \rightarrow w^{+}$as $t \rightarrow \infty$ in $X^{\alpha} \subset C^{1}(\bar{\Omega})$. Since the complete trajectory $\phi$ does not blow up, from (3.7) we obtain

$$
\forall_{t \in \mathbb{R}} \forall_{x \in \bar{\Omega}} \quad \phi(t)(x) \leq w^{-}(x)
$$

If we show that

$$
\begin{equation*}
\exists_{t^{*}<0} \forall_{t \leq t^{*}} \forall_{x \in \bar{\Omega}} \quad w^{+}(x) \leq \phi(t)(x) \tag{3.8}
\end{equation*}
$$

by monotonicity we will then have

$$
\forall_{t \in \mathbb{R}} \forall_{x \in \bar{\Omega}} \quad w^{+}(x) \leq \phi(t)(x) \leq w^{-}(x)
$$

Hence for the boundedness of $\mathcal{S}$ it is sufficient to prove (3.8).
Contrary to our claim, assume that there exist $t_{n} \rightarrow-\infty, y_{n} \in \bar{\Omega}$ such that

$$
\begin{equation*}
\phi\left(t_{n}\right)\left(y_{n}\right)<w^{+}\left(y_{n}\right) \tag{3.9}
\end{equation*}
$$

Then there exist $x_{n} \rightarrow x_{0}$ with $x_{0} \in \bar{\Omega}$ and

$$
\begin{equation*}
\phi\left(t_{n}\right)\left(x_{n}\right)=w^{+}\left(x_{n}\right), \quad n \in \mathbb{N} . \tag{3.10}
\end{equation*}
$$

Indeed, there exists $n_{0} \in \mathbb{N}$ such that for $n \geq n_{0}$ we have $\phi\left(t_{n}\right)(0)>w^{+}(0)$. Since (3.9) holds, the Darboux property ensures the existence of $x_{n} \in \bar{\Omega}$ such that $\phi\left(t_{n}\right)\left(x_{n}\right)=w^{+}\left(x_{n}\right)$. By the compactness of $\bar{\Omega}$ we may choose a convergent subsequence of $\left\{x_{n}\right\}$, still denoted by $\left\{x_{n}\right\}$, which is as required.

We consider two cases. If $x_{0} \in(-1,1)$, then $\phi\left(t_{n}\right)\left(x_{n}\right)$ tends to $w^{-}\left(x_{0}\right)$ and $w^{-}\left(x_{0}\right)=w^{+}\left(x_{0}\right)$, which is impossible. Therefore $x_{0} \in\{-1,1\}$. From (3.5) we get

$$
\left(w^{+}\right)^{\prime}(1)>\left(w^{-}\right)^{\prime}(1), \quad\left(w^{+}\right)^{\prime}(-1)<\left(w^{-}\right)^{\prime}(-1)
$$

Let us consider $x_{0}=1$. By (3.10) and the mean value theorem we have

$$
\frac{w^{+}\left(x_{n}\right)-w^{+}(1)}{x_{n}-1}=\frac{\phi\left(t_{n}\right)\left(x_{n}\right)-\phi\left(t_{n}\right)(1)}{x_{n}-1}=\left[\phi\left(t_{n}\right)\right]^{\prime}\left(\xi_{n}\right) .
$$

Since the left hand side tends to $\left(w^{+}\right)^{\prime}(1)$ and the right hand side to $\left(w^{-}\right)^{\prime}(1)$, we get $\left(w^{+}\right)^{\prime}(1)=\left(w^{-}\right)^{\prime}(1)$, a contradiction. The same reasoning applies to $x_{0}=-1$. This ends the proof of the boundedness of $\mathcal{S}$.

Observe that the solution semigroup for $u \in X^{\alpha}$ satisfying $w^{+}(x) \leq$ $u(x) \leq w^{-}(x)$ for all $x \in \bar{\Omega}$ is compact. Hence it has a compact connected global attractor. In particular there must exist a heteroclinic orbit connecting $w^{-}$to $w^{+}$.

In fact there exists a unique heteroclinic orbit connecting these two equilibria. This can be established using the argument in [B-F3, Lemma 3.5]. Thus we obtain the description of $\mathcal{S}=\left\{w^{-}, w^{+}, \phi(\mathbb{R})\right\}$, where $\phi$ is the only complete trajectory connecting $w^{-}$to $w^{+}$.

We conclude that $\mathcal{S}$ is a maximal compact invariant set attracting any subset of $V$ with bounded orbit. In particular, if $u_{0} \in X^{\alpha}$, then the corresponding $X^{\alpha}$ solution either blows up or stays bounded and approaches a maximal compact invariant set.

Example 3.2. Consider the $N$-dimensional Navier-Stokes system for incompressible viscous fluid flow subject to a small perturbation. Following [C-D2, Section 4], we shall show that in this case there exists a Lyapunov function. We consider the problem

$$
\left\{\begin{array}{l}
u_{t}=\nu \Delta u-\nabla p-(u \cdot \nabla) u+f, \quad t>0, x \in \Omega  \tag{3.11}\\
\operatorname{div} u=0, \quad t>0, x \in \Omega \\
u(t, x)=0, \quad t>0, x \in \partial \Omega \\
u(0, x)=u_{0}(x), \quad x \in \Omega
\end{array}\right.
$$

where $N \geq 2, \nu>0$ is a viscosity constant and $\Omega$ is a bounded domain in $\mathbb{R}^{N}$ with boundary $\partial \Omega$ of class $C^{2+\varepsilon}$.

For any $f \in\left[L^{p}(\Omega)\right]^{N}, p>N$, the system can be viewed as an abstract Cauchy problem

$$
\left\{\begin{array}{l}
u_{t}+A u=F(u), \quad t>0  \tag{3.12}\\
u(0)=u_{0}
\end{array}\right.
$$

in the space

$$
X=\operatorname{cl}_{\left[L^{p}(\Omega)\right]^{N}}\left\{\phi \in\left[C_{0}^{\infty}(\Omega)\right]^{N}: \operatorname{div} \phi=0 \text { in } \Omega\right\}
$$

using the projector $P$ given by the decomposition of $\left[L^{p}(\Omega)\right]^{N}$ into the spaces of divergence free vector fields and scalar function gradients (see $[\mathrm{F}-\mathrm{M}]$ and $[\mathrm{G}-\mathrm{M}])$. Namely, we define $A=-\nu P \Delta: X \supset \operatorname{dom}(A) \rightarrow X$ with

$$
\operatorname{dom}(A)=X \cap\left\{\phi \in\left[W^{2, p}(\Omega)\right]^{N}: \phi=0 \text { on } \partial \Omega\right\}
$$

which is a sectorial operator with compact resolvent, and $F: X^{\alpha} \rightarrow X$ by

$$
\begin{equation*}
F(u)=-P(u \cdot \nabla) u+P f, \quad u \in X^{\alpha} \tag{3.13}
\end{equation*}
$$

Restricting further $\alpha$ to the interval $[1 / 2,1)$ we observe that $F$ in (3.13) is well defined and is Lipschitz continuous on bounded subsets of $X^{\alpha}$.

Recall that for sufficiently small $f \in\left[L^{p}(\Omega)\right]^{N}$ (especially if the external force $f$ is zero)
(3.14) there exists a stationary solution $w \in \operatorname{dom}(A)$ of the Navier-Stokes system such that $\|w\|_{\left[W^{1, \infty}(\Omega)\right]^{N}}<\nu / C_{\Omega}^{2}$,
where $C_{\Omega}$ is the constant in the Poincaré inequality.
Lyapunov function on $V$ and description of $\mathcal{S}$. Assuming (3.14) we define $V$ as in (2.2) and consider the functional

$$
\begin{equation*}
\mathcal{L}(u)=\frac{1}{2}\|u-w\|_{\left[L^{2}(\Omega)\right]^{N}}^{2}, \quad u \in V \tag{3.15}
\end{equation*}
$$

We shall show that $\mathcal{L}$ is a Lyapunov function on $V$. Since $p>N \geq 2$, it follows that $\mathcal{L}$ is continuous on $V$. Fix $u_{0} \in V$. Letting $u(t)=u\left(t, u_{0}\right), t \geq 0$, we have

$$
(u-w)_{t}=-A(u-w)-P((u-w) \cdot \nabla) w-P(u \cdot \nabla)(u-w)
$$

for $t>0$. From $[\mathrm{F}-\mathrm{M}]$ it follows that

$$
P v=P_{2} v, \quad v \in\left[L^{p}(\Omega)\right]^{N}
$$

where $P_{2}$ is a selfadjoint bounded projection operator on $\left[L^{2}(\Omega)\right]^{N}$. Hence for $v_{1}, v_{2}, v_{3} \in \operatorname{dom}(A)$ we have

$$
\begin{aligned}
\left\langle P\left(v_{1} \cdot \nabla\right) v_{2}, v_{3}\right\rangle_{\left[L^{2}(\Omega)\right]^{N}} & =\left\langle P_{2}\left(v_{1} \cdot \nabla\right) v_{2}, v_{3}\right\rangle_{\left[L^{2}(\Omega)\right]^{N}} \\
& =\left\langle\left(v_{1} \cdot \nabla\right) v_{2}, v_{3}\right\rangle_{\left[L^{2}(\Omega)\right]^{N}}
\end{aligned}
$$

and

$$
\left\langle A v_{1}, v_{1}\right\rangle_{\left[L^{2}(\Omega)\right]^{N}}=-\nu\left\langle\Delta v_{1}, v_{1}\right\rangle_{\left[L^{2}(\Omega)\right]^{N}} .
$$

Multiplying (3.16) by $u-w$ in $\left[L^{2}(\Omega)\right]^{N}$ we obtain

$$
\begin{align*}
\frac{1}{2} \frac{d}{d t}\|u-w\|_{\left[L^{2}(\Omega)\right]^{N}}^{2} \leq & -\nu \sum_{i=1}^{N}\left\|\nabla\left(u_{i}-w_{i}\right)\right\|_{L^{2}(\Omega)}^{2}  \tag{3.17}\\
& +\|w\|_{\left[W^{1, \infty}(\Omega)\right]^{N}}\|u-w\|_{\left[L^{2}(\Omega)\right]^{N}}^{2} \\
\leq & \left(-\nu / C_{\Omega}^{2}+\|w\|_{\left[W^{1, \infty}(\Omega)\right]^{N}}^{2}\|u-w\|_{\left[L^{2}(\Omega)\right]^{N}}^{2}\right.
\end{align*}
$$

The above inequality proves that $t \mapsto \mathcal{L}\left(u\left(t, u_{0}\right)\right)$ is nonincreasing for $t>0$. Moreover, there exists $a>0$ such that

$$
\|u-w\|_{\left[L^{2}(\Omega)\right]^{N}} \leq\left\|u_{0}-w\right\|_{\left[L^{2}(\Omega)\right]^{N}} e^{-a t}, \quad t \geq 0 .
$$

This shows in particular that

$$
\begin{equation*}
\lim _{t \rightarrow \infty} u\left(t, u_{0}\right)=w \quad \text { in }\left[L^{2}(\Omega)\right]^{N} . \tag{3.18}
\end{equation*}
$$

Our task will be, however, to justify the convergence in $X^{\alpha}$. Meanwhile, assume that $\mathcal{L}\left(u\left(\cdot, u_{0}\right)\right) \equiv \mathcal{L}\left(u_{0}\right)$. Then as in (3.17) we have

$$
0=\frac{d}{d t} \mathcal{L}\left(u\left(t, u_{0}\right)\right) \leq\left(-\nu / C_{\Omega}^{2}+\|w\|_{\left.\left[W^{1, \infty}(\Omega)\right]^{N}\right)}\left\|u\left(t, u_{0}\right)-w\right\|_{\left[L^{2}(\Omega)\right]^{N}}^{2} \leq 0 .\right.
$$

This implies $u_{0}=w \in \mathcal{E}$ and consequently $\mathcal{E}=\{w\}$.
Suppose that $\phi$ is a (hypothetical) bounded complete trajectory of $u_{0}$. Since $\alpha_{\phi}\left(u_{0}\right)=\{w\}$ and $\omega\left(u_{0}\right)=\{w\}$, we have

$$
\mathcal{L}(w) \geq \mathcal{L}(\phi(s)) \geq \mathcal{L}(w), \quad s \in \mathbb{R}
$$

so that $\phi(\mathbb{R})=\{w\}$. Hence we obtain $\mathcal{S}=\{w\}$.
We thus infer that if (3.14) holds, then for any $u_{0} \in X^{\alpha}, \alpha \in[1 / 2,1)$, the $X^{\alpha}$ solution $u\left(\cdot, u_{0}\right)$ of the Navier-Stokes system either blows up in $X^{\alpha}$ or stays bounded and approaches in $X^{\alpha}$ the maximal compact invariant set $\{w\}$.

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