

The norms and singular numbers of polynomials of the classical Volterra operator in $L_2(0, 1)$

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Abstract. The spectral problem $(s^2I - \phi(V)^*\phi(V))f = 0$ for an arbitrary complex polynomial ϕ of the classical Volterra operator V in $L_2(0, 1)$ is considered. An equivalent boundary value problem for a differential equation of order $2n$, $n = \deg(\phi)$, is constructed. In the case $\phi(z) = 1 + az$ the singular numbers are explicitly described in terms of roots of a transcendental equation, their localization and asymptotic behavior is investigated, and an explicit formula for the $\|I + aV\|_2$ is given. For all $a \neq 0$ this norm turns out to be greater than 1.

1. Introduction. For any compact linear operator A in a Hilbert space the singular numbers $s_k(A)$ are the distances from A to the set of all operators of rank less than or equal to $k - 1$, $k \geq 1$. Their squares are the eigenvalues of the compact selfadjoint nonnegative operator A^*A counted according to their multiplicities (see e.g. [1]). In particular, $s_1(A) = \|A\|$. The latter has been used by Halmos [2] to calculate the L_2 -norm $\|\cdot\|_2$ of the classical Volterra operator

$$(Vf)(x) = \int_0^x f(t) dt, \quad 0 \leq x \leq 1.$$

Actually, the Halmos calculation yields

$$(1.1) \quad s_k(V) = \frac{2}{(2k-1)\pi}$$

for all $k \geq 1$, in particular,

$$(1.2) \quad \|V\|_2 = 2/\pi.$$

The point is that the spectral problem

$$(V^*Vf)(x) = s^2f(x), \quad 0 \leq x \leq 1,$$

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is equivalent to the boundary value problem

$$f''(x) + \lambda f(x) = 0, \quad f'(0) = 0, \quad f(1) = 0,$$

where $\lambda = 1/s^2$. This yields $f(x) = \cos \sqrt{\lambda}x$ (under the normalization $f(0) = 1$) and $\cos \sqrt{\lambda} = 1$, which immediately implies (1.1).

A similar equivalence for the powers V^n was established by Thorpe [4]. The corresponding boundary value problem is

$$(-1)^n f^{(2n)} = \lambda f, \quad f^{(l)}(0) = f^{(n+l)}(1) = 0, \quad 0 \leq l \leq n-1.$$

In the present paper we generalize these results to the arbitrary complex polynomials

$$(1.3) \quad \phi(V) = \sum_{i=0}^n a_i V^i, \quad a_n \neq 0, \quad n \geq 1.$$

Note that the operator $\phi(V)$ is not compact if $a_0 \neq 0$, but in any case $\phi(V)^* \phi(V) = |a_0|^2 I + K$ where I is the identity operator and K is a compact self-adjoint operator. This suggests defining the *singular numbers* of $\phi(V)$ as the nonnegative square roots of the eigenvalues of $\phi(V)^* \phi(V)$. In fact, they are positive, since the operator $\phi(V)$ is injective. Indeed, it is invertible if $a_0 \neq 0$, otherwise, it is of the form $V^l \psi(V)$, $l \geq 1$, with $\psi(V)$ invertible. Actually, the singleton $\{a_0\}$ is the spectrum of $\phi(V)$.

The singular numbers of $\phi(V)$ constitute a countable set S_ϕ converging to $|a_0|$. We have

$$(1.4) \quad \sup S_\phi = \|\phi(V)\|_2 \geq |a_0|.$$

In Section 2 we consider the case $n = 1$ and obtain an explicit formula for the singular numbers, in particular, for $\|I + aV\|_2$ in terms of roots of a transcendental equation that comes from a boundary value problem. We describe the localization of these roots in much detail. The general case $n \geq 1$ is considered in Section 3 where we construct a boundary value problem equivalent to the spectral problem in question. In Section 4 the problem of equality in (1.4) is discussed.

2. The case $n = 1$. Let $\phi(V) = a_0 I + a_1 V$, $a_1 \neq 0$. Since the case $a_0 = 0$ trivially reduces to that of [2], one can assume $a_0 \neq 0$. Without loss of generality one can set $a_0 = 1$ and then denote a_1 by a , for short. Thus, we will consider $\phi(V) = I + aV$ with $a \neq 0$.

Our spectral problem

$$(I + \bar{a}V^*)(I + aV)f = s^2 f$$

can be rewritten as

$$(2.1) \quad (s^2 - 1)f - \bar{a}V^* f - aVf - |a|^2 V^* V f = 0.$$

We proceed from this integral equation to a differential equation by applying the operator $D = d/dx$ twice. Note that $DV = I$, while $DV^* = -I$ since

$$(V^*f)(x) = \int_x^1 f(t) dt, \quad 0 \leq x \leq 1.$$

Thus, (2.1) yields

$$(2.2) \quad (s^2 - 1)f' - (a - \bar{a})f + |a|^2Vf = 0$$

and then

$$(2.3) \quad (s^2 - 1)f'' - (a - \bar{a})f' + |a|^2f = 0.$$

Now we insert $x = 0$ and $x = 1$ into (2.1) and (2.2) taking into account that

$$(Vf)_{x=0} = (V^*f)_{x=1} = 0,$$

and

$$(Vf)_{x=1} = (V^*f)_{x=0} = J_1 \equiv \int_0^1 f dt,$$

hence,

$$(V^*Vf)_{x=1} = 0, \quad (V^*Vf)_{x=0} = J_2 \equiv \int_0^1 (Vf) dt.$$

In this way we obtain four linear homogeneous equations for the six values: $f(0)$, $f'(0)$, $f(1)$, $f'(1)$, J_1 , J_2 . Eliminating J_1 and J_2 we get two boundary conditions for the differential equation (2.3). Let us emphasize that in [2] and [4] no integrals remain after substitution of $x = 0$ and $x = 1$ into the corresponding integral equation and its derivatives. In our case this is true only for one of the four equalities, namely, we get

$$(2.4) \quad (s^2 - 1)f'(0) - (a - \bar{a})f(0) = 0$$

when putting $x = 0$ in (2.2). However,

$$(s^2 - 1)f(1) - aJ_1 = 0, \quad (s^2 - 1)f'(1) - (a - \bar{a})f(1) + |a|^2J_1 = 0$$

for $x = 1$ in (2.1) and (2.2), respectively. Eliminating J_1 we obtain

$$(2.5) \quad (s^2 - 1)f'(1) + (\bar{a}s^2 - a)f(1) = 0.$$

LEMMA 2.1. *For every s the integral equation (2.1) is equivalent to the differential equation (2.3) with the boundary conditions (2.4) and (2.5).*

Proof. We already know that each solution f to the equation (2.1) satisfies (2.3)–(2.5). In the converse direction we start with f satisfying (2.3)–(2.5) and set

$$g = ((I + \bar{a}V^*)(I + aV) - s^2I)f.$$

We have to show that $g = 0$. To this end we note that $g'' = 0$ by (2.3), hence $g(x)$ is a linear function of x . Furthermore,

$$g'(0) = (s^2 - 1)f'(0) - (a - \bar{a})f(0),$$

hence $g'(0) = 0$ by (2.4). Thus, $g(x)$ is a constant, so

$$\bar{a}g(1) = \bar{a}g(1) + g'(1) = (s^2 - 1)f'(1) - (\bar{a}s^2 - a)f(1)$$

by (2.5). Since $a \neq 0$, we get $g(1) = 0$. ■

The singular numbers in question are just those s for which the boundary value problem (2.3)–(2.5) has a solution $f \neq 0$. It is easy to see that $s \neq 1$. Indeed, otherwise from (2.3) it follows that

$$(2.6) \quad (a - \bar{a})f' - |a|^2 f = 0.$$

If a is real then $f = 0$ at once. If a is not real then f satisfies the linear differential equation (2.6) and, in addition, $f(0) = 0$ by (2.4). Hence, $f = 0$ again.

Now we are in a position to prove the following theorem.

THEOREM 2.2. *Let $a = \alpha + i\beta$. The singular numbers of the operator $I + aV$ are given by the formula*

$$(2.7) \quad s = \sqrt{\frac{\alpha^2 + \Delta}{\alpha^2 + \beta^2}}$$

where α and β are the real and the imaginary parts of a and Δ runs over all real roots of the equation

$$(2.8) \quad \sqrt{\Delta} \cot \frac{\sqrt{\Delta}(\alpha^2 + \beta^2)}{\Delta - \beta^2} = \alpha.$$

Though the value $\sqrt{\Delta}$ is determined only up to the factor ± 1 , the right-hand side of (2.8) is uniquely determined since the function $\cot(\cdot)$ is odd. Also note that if Δ is a root of the equation (2.8) then

$$\Delta \neq \beta^2, \quad \frac{\sqrt{\Delta}(\alpha^2 + \beta^2)}{\Delta - \beta^2} \neq m\pi$$

for all integers m , in particular, $\Delta \neq 0$. However, $\Delta = 0$ becomes admissible by passing to the limit as $\Delta \rightarrow 0$. The limit equality is

$$\frac{\beta^2}{\alpha^2 + \beta^2} = -\alpha,$$

or equivalently,

$$(2.9) \quad \beta^2 = -\frac{\alpha^3}{\alpha + 1}$$

where automatically $-1 < \alpha < 0$ and $\beta \neq 0$, so a is not real. Under the relation (2.9) Theorem 2.2 extends by including $\Delta = 0$ into (2.7), so we have

the special singular number

$$(2.10) \quad s^0 = \frac{|\alpha|}{\sqrt{\alpha^2 + \beta^2}} = \sqrt{1 + \alpha}$$

in this situation.

For $\Delta < 0$ the equation (2.8) can be rewritten as

$$(2.11) \quad \sqrt{|\Delta|} \coth \frac{\sqrt{|\Delta|}(\alpha^2 + \beta^2)}{|\Delta| + \beta^2} = -\alpha.$$

Proof of Theorem 2.2. In our current notation the equation (2.3) is

$$(2.12) \quad (s^2 - 1)f'' - 2i\beta f' + (\alpha^2 + \beta^2)f = 0,$$

and the boundary conditions (2.4) and (2.5) are

$$(2.13) \quad (s^2 - 1)f'(0) - 2i\beta f(0) = 0$$

and

$$(2.14) \quad (s^2 - 1)f'(1) + ((s^2 - 1)\alpha - (s^2 + 1)i\beta)f(1) = 0.$$

The characteristic equation for the differential equation (2.12) is

$$(2.15) \quad (s^2 - 1)r^2 - 2i\beta r + (\alpha^2 + \beta^2) = 0.$$

Its roots are

$$(2.16) \quad r_1 = i \frac{\beta + \sqrt{\Delta}}{s^2 - 1}, \quad r_2 = i \frac{\beta - \sqrt{\Delta}}{s^2 - 1},$$

where

$$(2.17) \quad \Delta = \beta^2 + (s^2 - 1)(\alpha^2 + \beta^2).$$

The latter is equivalent to (2.7). We have to prove that (2.8) (including the equality $\Delta = 0$ in the case (2.9)) is equivalent to the nontrivial solvability of the boundary value problem (2.12)–(2.14). Let us start with $\Delta \neq 0$, i.e. $r_1 \neq r_2$.

The general solution to (2.12) is

$$f(x) = c_1 e^{r_1 x} + c_2 e^{r_2 x},$$

where c_1 and c_2 are arbitrary constants. Hence,

$$f(0) = c_1 + c_2, \quad f'(0) = c_1 r_1 + c_2 r_2$$

and

$$f(1) = c_1 e^{r_1} + c_2 e^{r_2}, \quad f'(1) = c_1 r_1 e^{r_1} + c_2 r_2 e^{r_2}.$$

By substitution into (2.13) and (2.14) we obtain a system of linear homogeneous equations for c_1 and c_2 . This system has a nontrivial solution if and only if its determinant is equal to zero. This equality reduces to

$$((s^2 - 1)r_1 - \alpha - i\beta)e^{r_1} = ((s^2 - 1)r_2 - \alpha - i\beta)e^{r_2}$$

by an elementary calculation taking (2.15) into account. Using (2.16) we obtain

$$(2.18) \quad (\alpha - i\sqrt{\Delta})e^{r_1} = (\alpha + i\sqrt{\Delta})e^{r_2}.$$

Note that $\alpha - i\sqrt{\Delta} \neq 0$, otherwise, $\alpha^2 + \Delta = 0$, which implies $s = 0$ by (2.17). Therefore, one can rewrite (2.18) as

$$e^{r_1 - r_2} = \frac{\alpha + i\sqrt{\Delta}}{\alpha - i\sqrt{\Delta}},$$

or equivalently,

$$\frac{e^{r_1 - r_2} + 1}{e^{r_1 - r_2} - 1} = \frac{\alpha}{i\sqrt{\Delta}}.$$

This is nothing but the equation (2.8) since

$$r_1 - r_2 = \frac{2i\sqrt{\Delta}}{s^2 - 1} = \frac{2i\sqrt{\Delta}(\alpha^2 + \beta^2)}{\Delta - \beta^2}$$

by (2.16) and (2.17).

In the case $\Delta = 0$ the only root of (2.15) is

$$r = \frac{i\beta}{s^2 - 1} = -\frac{i(\alpha^2 + \beta^2)}{\beta}$$

since

$$s^2 - 1 = -\frac{\beta^2}{\alpha^2 + \beta^2}$$

by (2.17). (Note that $\beta \neq 0$ since $s \neq 1$.) The general solution to (2.12) is now of the form

$$f(x) = (c_1 + c_2x)e^{rx}.$$

Accordingly,

$$f(0) = c_1, \quad f'(0) = c_1r_1 + c_2r_2$$

and

$$f(1) = c_1e^{r_1} + c_2e^{r_2}, \quad f'(1) = c_1r_1e^{r_1} + c_2r_2e^{r_2}.$$

It is easy to check that the determinant of the corresponding linear system for c_1 and c_2 vanishes if and only if the relation (2.9) is valid. As we know, the latter is the limit form of (2.8) as $\Delta \rightarrow 0$. ■

Now we investigate the equation (2.8) with unknown $\Delta \neq 0$ (written as (2.11) for $\Delta < 0$). We start with $\beta = 0$. For the new unknown

$$(2.19) \quad \xi = \alpha^2 / \sqrt{|\Delta|} > 0$$

we have the equation

$$(2.20) \quad \coth \xi = -\xi/\alpha$$

if $\Delta < 0$, but

$$(2.21) \quad \cot \xi = \xi/\alpha$$

if $\Delta > 0$. By (2.7) and (2.19) the corresponding singular number is

$$(2.22) \quad s = \sqrt{1 + \frac{\varepsilon\alpha^2}{\xi^2}}$$

where ξ is a root of (2.20) or (2.21) and $\varepsilon = -1$ or 1 , respectively. By these equations one can rewrite (2.22) as

$$(2.23) \quad s = 1/\cosh \xi \quad \text{or} \quad s = 1/|\cos \xi|,$$

respectively.

The equation (2.20) has no positive roots if $\alpha > 0$, but if $\alpha < 0$ then (2.20) has exactly one positive root. Indeed, in the latter case the function $\coth \xi + \xi/\alpha$ on $(0, \infty)$ monotonically decreases from $+\infty$ to $-\infty$.

In contrast, the equation (2.21) has infinitely many positive roots for any $\alpha \neq 0$: there is exactly one root ξ_k of (2.21) in each interval $((k-1)\pi, k\pi)$, $k \geq 1$. Accordingly,

$$(2.24) \quad s_k = \sqrt{1 + \frac{\alpha^2}{\xi_k^2}} = \frac{1}{|\cos \xi_k|}, \quad k \geq 1.$$

Now let $\beta \neq 0$. Then we introduce the new unknown

$$(2.25) \quad \xi = \sqrt{|\Delta|}/|\beta| > 0,$$

instead of Δ , and the new real parameters

$$(2.26) \quad \gamma = \frac{\alpha^2 + \beta^2}{|\beta|}, \quad \delta = \frac{\alpha}{|\beta|},$$

instead of α and β . In this setting we have

$$(2.27) \quad \xi \coth \frac{\gamma\xi}{\xi^2 + 1} = -\delta$$

if $\Delta < 0$, but

$$(2.28) \quad \xi \cot \frac{\gamma\xi}{\xi^2 - 1} = \delta$$

if $\Delta > 0$. The corresponding singular number is

$$(2.29) \quad s = \sqrt{\frac{\delta^2 + \varepsilon\xi^2}{\delta^2 + 1}}$$

where ξ is a root of (2.27) or (2.28) and ε is defined as in (2.22).

Since $\gamma > 0$, the equation (2.27) has no positive roots if $\delta \geq 0$. Let $\delta < 0$, i.e. $\alpha < 0$ by (2.26). Then all positive roots of (2.27) (if they exist) are less

than $|\delta|$. With this restriction (2.27) is equivalent to $g(\xi) = 0$ where

$$g(\xi) = \frac{2\gamma\xi}{\xi^2 + 1} - \log \frac{|\delta| + \xi}{|\delta| - \xi}.$$

The derivative $g'(\xi)$ is

$$\gamma(\xi^2 - 1)(\xi^2 - \delta^2) - |\delta|(\xi^2 + 1)^2,$$

up to a positive factor. This biquadratic polynomial has at most two positive roots, so the same is true of $g'(\xi)$ and hence of $g(x)$ because of Rolle's theorem and $g(0+) = 0$. Thus, *the equation (2.27) has at most two positive roots if $\delta < 0$.*

It remains to investigate the positive roots of the equation (2.28). With $\delta = 0$ they are

$$(2.30) \quad \xi_k^0 = \frac{\gamma + \varepsilon_k \sqrt{\gamma^2 + (2k - 1)^2 \pi^2}}{(2k - 1)\pi}$$

where k runs over all integers and $\varepsilon_k = \text{sign}(2k - 1)$. It is easy to see that $\xi_{k+1}^0 < \xi_k^0$ for $k \neq 0$ but $\xi_1^0 > \xi_0^0$. Moreover, $\xi_1^0 > 1$, while $\xi_0^0 < 1$. (It is useful to note that $\xi_{-(k-1)}^0 \xi_k^0 = 1$.) Obviously, the roots ξ_k^0 and ξ_{-k}^0 tend to 1 as $k \rightarrow +\infty$.

Now let $\delta \neq 0$. Then with $\xi \neq \xi_k^0$ ($k = 0, \pm 1, \pm 2, \dots$), (2.28) is equivalent to $h(\xi) = 0$ where

$$h(\xi) = h_0(\xi) - \arctan\left(\frac{\xi}{\delta}\right) - q(\xi)\pi, \quad h_0(\xi) = \frac{\gamma\xi}{\xi^2 - 1}$$

and an integer coefficient $q(\xi)$ is determined by the inequality

$$-\frac{\pi}{2} < h_0(\xi) - q(\xi)\pi < \frac{\pi}{2}.$$

The function $h_0(\xi)$ monotonically decreases from 0 to $-\infty$ on the interval $(0, 1)$ and from $+\infty$ to 0 on the interval $(1, \infty)$. Since

$$h_0(\xi_k^0) = \frac{(2k - 1)\pi}{2},$$

we have $q(\xi) = k$ for $\xi_{k+1}^0 < \xi < \xi_k^0$, $k \neq 0$, and $q(\xi) = 0$ for $\xi < \xi_0^0$ or $\xi > \xi_1^0$. Thus,

$$h(\xi) = h_0(\xi) - \arctan\left(\frac{\xi}{\delta}\right) - k\pi$$

on the interval $J_k = (\xi_{k+1}^0, \xi_k^0)$, $k \neq 0$, and

$$h(\xi) = h_0(\xi) - \arctan\left(\frac{\xi}{\delta}\right)$$

on the intervals $J_0 = (0, \xi_0^0)$ and $J_\infty = (\xi_1^0, \infty)$. On each of these intervals

the function $h(\xi)$ is continuous. Moreover, we have the one-sided limit values

$$h(\xi_k^0-) = -\left(\frac{\pi}{2} + \arctan\left(\frac{\xi_k^0}{\delta}\right)\right) < 0, \quad h(\xi_k^0+) = \frac{\pi}{2} - \arctan\left(\frac{\xi_k^0}{\delta}\right) > 0$$

for all k , and

$$(2.31) \quad h(0+) = 0, \quad h(+\infty) = -\frac{\pi}{2} \operatorname{sign} \delta.$$

Therefore, the equation (2.28) has at least one root in each interval J_k , $k \neq 0, \infty$. For $k = \infty$ this is true if $\delta > 0$. For $k = 0$ it is true if $\delta < 0$ and $|\delta|\gamma < 1$. Indeed, in this case $h'(0) > 0$, so $h(\xi) > 0$ for small $\xi > 0$.

In general, the derivative $h'(\xi)$ is

$$\gamma(\xi^2 + 1)(\xi^2 + \delta^2) + \delta(\xi^2 - 1)^2,$$

up to a negative factor. If $\delta > 0$ then $h'(\xi) < 0$ for all $\xi > 0$. In this case the equation (2.28) has exactly one root $\xi_k \in J_k, k \neq 0$, including J_∞ , while J_0 does not contain roots at all.

If $\delta < 0$ then $h'(\xi)$ has at most two positive roots. These roots can lie either in a J_k , and then (2.28) has no more than three roots there, or they belong to some different J_k and J_l , and then the number of roots in each of them does not exceed 2. Any other interval J_m contains exactly one root. The number of roots in J_0 and in J_∞ does not exceed 2 because of (2.31).

Now we denote by ξ_k a root of $h(\xi)$ in the interval J_k . The sequences $(\xi_k : k \geq 1)$ and $(\xi_{-k} : k \geq 0)$ monotonically tend to 1 from above and from below, respectively. By (2.29) with $\varepsilon = 1$ the same is true for $(s_k : k \geq 1)$ and $(s_{-k} : k \geq 0)$. All these singular numbers are greater than $|\delta|/\sqrt{\delta^2 + 1}$. This lower bound is just the special singular number s^0 (see (2.10)) if $\Delta = 0$ is a root of the equation (2.8), i.e. if (2.9) is valid. There are at most two singular numbers below this bound since they appear only if $\Delta < 0$.

THEOREM 2.3. *Let $a = \alpha + i\beta \neq 0$. If $\beta = 0$, i.e. $a = \alpha \in \mathbb{R} \setminus \{0\}$, then*

$$(2.32) \quad \|I + aV\|_2 = \sqrt{1 + \frac{\alpha^2}{\xi_{\min}^2}} = \frac{1}{|\cos \xi_{\min}|}$$

where ξ_{\min} is the smallest positive root of the equation (2.21).

If $\beta \neq 0$, then

$$(2.33) \quad \|I + aV\|_2 = \sqrt{\frac{\delta^2 + \xi_{\max}^2}{\delta^2 + 1}}$$

where ξ_{\max} is the greatest root of the equation (2.28) with δ and γ defined by (2.26).

Proof. Combine formula (1.4) with (2.24) in the case $\beta = 0$ and with (2.29) in the case $\beta \neq 0, \varepsilon = 1$. ■

COROLLARY 2.4. *If $\alpha = 0$, i.e. $a = i\beta$, where $\beta \in \mathbb{R} \setminus \{0\}$, then*

$$(2.34) \quad \|I + aV\|_2 = \frac{|\beta|}{\pi} + \sqrt{1 + \frac{\beta^2}{\pi^2}}.$$

Indeed, in this case we have (2.33) with $\delta = 0$ and $\xi_{\max} = \xi_1^0$ coming from (2.30) with $\gamma = |\beta|$.

As $\beta \rightarrow \infty$, formula (2.34) leads to Halmos's formula (1.2). Similarly, (1.1) follows from (2.29) and (2.30).

COROLLARY 2.5. *The inequality*

$$(2.35) \quad \|I + aV\|_2 > 1$$

holds for all complex $a \neq 0$.

Proof. This follows from (2.32) if $\beta = 0$ and from (2.33) if $\beta \neq 0$ since $\xi_{\max} > 1$ in either case. ■

Another way to get the inequality (2.35) (except for $a \leq 0$) is to recall that the operator $I + aV$ is not power bounded if a is not real nonpositive (see [5]). With $a < 0$ this operator becomes power bounded [5], but its norm remains greater than 1 by Corollary 2.5.

3. The general case. For an arbitrary polynomial ϕ of degree $n \geq 1$ we proceed from the integral equation

$$(3.1) \quad (s^2I - \phi(V)^*\phi(V))f = 0$$

to

$$(3.2) \quad D^{2n}(s^2I - \phi(V)^*\phi(V))f = 0,$$

by differentiation of order $2n$. In more detail, if $\phi(V)$ is of the form (1.3) then the equation (3.1) can be rewritten as

$$\left((s^2 - |a_0|^2)I - \sum \bar{a}_i a_k (V^*)^i V^k \right) f = 0$$

with the summation over $0 \leq i, k \leq m$, $i + k \geq 1$. Accordingly, (3.2) takes the form

$$(3.3) \quad (s^2 - |a_0|^2)f^{(2n)} - \sum_{j=1}^{2n} c_j f^{(2n-j)} = 0$$

where

$$(3.4) \quad c_j = \sum_{i+k=j} (-1)^i \bar{a}_i a_k.$$

Note that $c_{2n} = (-1)^n |a_n|^2 \neq 0$, so (3.3) cannot be an identity. This is a linear homogeneous differential equation with constant coefficients. The

formal leading coefficient is $s^2 - |a_0|^2$, so the order of the equation (3.3) is $2n$ if $s \neq |a_0|$. The other coefficients do not depend on s .

Now we substitute $x = 0$ and $x = 1$ into the intermediate integro-differential equations

$$(3.5) \quad D^i(s^2I - \phi(V)^*\phi(V))f = 0, \quad 0 \leq i \leq 2n - 1,$$

including (3.1). This yields $4n$ linear homogeneous equations for $4n$ boundary values

$$f(0), f'(0), \dots, f^{(2n-1)}(0), f(1), f'(1), \dots, f^{(2n-1)}(1)$$

and the integrals

$$((V^*)^i V^k f)(0), \quad (V^k f)(1), \quad 1 \leq i, k \leq n.$$

The total number of these integrals is $n(n + 2)$, but each one is a linear combination of $2n$ power moments

$$M_l = \int_0^1 t^l f(t) dt, \quad 0 \leq l \leq 2n - 1.$$

Indeed, from the classical formula

$$(V^k f)(x) = \frac{1}{(k - 1)!} \int_0^x (x - t)^{k-1} f(t) dt$$

and its version

$$((V^*)^i f)(x) = \frac{1}{(i - 1)!} \int_x^1 (t - x)^{i-1} f(t) dt$$

it follows that

$$(V^k f)(1) = \frac{1}{(k - 1)!} \int_0^1 (1 - t)^{k-1} f(t) dt$$

and

$$((V^*)^i V^k f)(0) = \frac{1}{(i - 1)!(k - 1)!} \int_0^1 f(t) dt \int_t^1 (s - t)^{k-1} s^{i-1} ds.$$

All these integrals are linear combinations of M_l , $0 \leq l \leq k - 1$, and M_{i+k-1} . The orders of these moments are less than $2n$.

Now let F be the $4n$ -column consisting of the values $f^{(m)}(0)$ and $f^{(m)}(1)$, $0 \leq m \leq 2n - 1$, and let M be the $2n$ -column consisting of the values M_l , $0 \leq l \leq 2n - 1$. Then we have an equation

$$(3.6) \quad AF + BM = 0$$

where A and B are matrices of sizes $4n \times 4n$ and $4n \times 2n$, respectively.

To eliminate M from (3.6) we take any $4n$ -row w such that $wB = 0$. Then we obtain a boundary condition $(wA)F = 0$. In fact, only linearly

independent systems of boundary conditions are interesting. Accordingly, we choose a basis $\{w_j : 1 \leq j \leq r\}$ in the space of solutions of the equation $wB = 0$ and get

$$(3.7) \quad (w_j A)F = 0, \quad 1 \leq j \leq r.$$

Note that $r \geq 2n$ since the vector equation $wB = 0$ can be rewritten as a system of $2n$ linear homogeneous scalar equations with $4n$ unknowns.

THEOREM 3.1. *The integral equation (3.1) is equivalent to the differential equation (3.2) with the boundary conditions (3.7), where $r = 2n$ and $\{w_j\}$ is any linearly independent system of $2n$ solutions of the equation $wB = 0$.*

Proof. We only need to prove that every solution f of the boundary value problem under consideration satisfies (3.1). In other words, we have to prove that $g = 0$ for

$$g = (s^2 I - \phi(V)^* \phi(V))f.$$

By (3.2) we have $D^{2n}g = 0$, therefore, g is a polynomial of degree less than or equal to $2n - 1$. The column G consisting of the values $g^{(m)}(0)$ and $g^{(m)}(1)$, $0 \leq m \leq 2n - 1$, satisfies the boundary conditions

$$(3.8) \quad w_j G = 0, \quad 1 \leq j \leq 2n,$$

since $G = AF + BM$ and $w_j AF = 0$ by (3.7), while $w_j B = 0$ by the choice of w_j .

The upper and the lower halves of the column G are G_0 and G_1 whose entries are $g^{(m)}(0)$ and $g^{(m)}(1)$, $0 \leq m \leq 2n - 1$, respectively. The Taylor expansion at $x = 0$ shows that $G_1 = CG_0$ where C is a $2n \times 2n$ -matrix. The matrix C is invertible by the Taylor expansion at $x = 1$. The equations (3.8) can be rewritten as

$$(3.9) \quad (u_j + v_j C)G_0 = 0, \quad 1 \leq j \leq 2n,$$

where u_j and v_j are the left and the right halves of the row w_j . From (3.9) it follows that $G_0 = 0$, and then $g = 0$ since the rows $u_j + v_j C$, $1 \leq j \leq 2n$, are linearly independent. The latter is true since the rows w_j , $1 \leq j \leq 2n$, are linearly independent and the matrix C is invertible. ■

By Theorem 3.1 the singular numbers of the operator $\phi(V)$ are just those s for which the boundary linear functionals $F \mapsto (w_j A)F$, $1 \leq j \leq 2n$, on the space of solutions of the differential equation (3.2) are linearly dependent. Equivalently, the latter means that

$$(3.10) \quad \det(w_j A F_i) = 0$$

where $\{f_i : 1 \leq i \leq 2n\}$ is an arbitrary basis of that space (a “fundamental system” of solutions of (3.2)). Thus, *the singular numbers are just the roots of the equation (3.10)*. The unknown s is contained in the matrix A and

in the columns F_i corresponding to the solutions f_i , $1 \leq i \leq 2n$. Actually, the matrix A is lower triangular and all its diagonal entries are equal to $s^2 - |a_0|^2$. The other entries of A do not depend on s .

Theorem 3.1 also implies that *for every singular number s the eigenspace of the operator $\phi(V)^*\phi(V)$ corresponding to the eigenvalue s^2 is of dimension $\leq 2n$.*

4. Appendix. An unsolved problem. According to Corollary 2.5 the only operator of L_2 -norm 1 in the family $\{I + aV : a \in \mathbb{C}\}$ is I .

PROBLEM. *Describe the class \mathcal{I} of functions $\phi(z)$ such that $\phi(0) = 1$ and $\|\phi(V)\|_2 = 1$.*

In this context $\phi(z)$ can be any function analytic in a neighborhood of $z = 0$, $\phi(0) = 1$. Since the spectrum of $\phi(V)$ is $\{1\}$ we have $\|\phi(V)\|_2 \geq 1$ a priori.

The class \mathcal{I} is not empty: the function $\mathbf{1}(z) \equiv 1$ belongs to \mathcal{I} . In total, \mathcal{I} is a convex multiplicative semigroup. On the other hand, if $\phi \in \mathcal{I}$ and $\phi \neq \mathbf{1}$ then $\mathbf{1}/\phi \notin \mathcal{I}$, i.e. $\|\phi(V)^{-1}\|_2 > 1$. Indeed, otherwise, $\phi(V)$ is an isometry, hence $\phi(V) = I$ by the classical Gelfand theorem.

For $\phi \in \mathcal{I}$ the operator $\phi(V)$ is power bounded. From Theorem 1.1 of [3] it follows that *if $\phi \in \mathcal{I} \setminus \{\mathbf{1}\}$ then $\phi'(0)$ is real negative*. This necessary condition is not sufficient even for linear functions, as we already know. Moreover, we do not know any nontrivial polynomial $\phi \in \mathcal{I}$. We conjecture that *there are no such polynomials*. However, there are some functional examples.

PROPOSITION 4.1. *If all roots of a polynomial ϕ are real negative and $\phi(0) = 1$ then $\mathbf{1}/\phi \in \mathcal{I}$.*

Proof. Halmos proved (see [2, Problem 150]) that $\|(I + V)^{-1}\|_2 = 1$. This means that the function $(1 + z)^{-1}$ belongs to \mathcal{I} . The same proof shows that the function $(1 + az)^{-1}$ belongs to \mathcal{I} for every $a > 0$. The general case is settled by the decomposition of $\phi(z)$ into factors of the form $1 + az$ where $-a$ runs over the reciprocals of the roots of ϕ . ■

REMARK 4.2. If a is not real nonnegative then $\phi(z) = (1 + az)^{-1}$ does not belong to \mathcal{I} since $\phi'(0) = -a$. In other words, $\|(I + aV)^{-1}\|_2 > 1$ for all complex a , except for $a \geq 0$.

Halmos's proof cited above is based on the inequality $\text{Re}(Vf, f) \geq 0$ that means that the operator $-V$ is dissipative. Hence, if $a \leq 0$ then $\|\exp(aV)\|_2 \leq 1$, so eventually $\|\exp(aV)\|_2 = 1$, thus $\exp(az) \in \mathcal{I}$. The sufficient condition $a \leq 0$ is also necessary since $a = \phi'(0)$.

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References

- [1] I. C. Gohberg and M. G. Kreĭn, *Introduction to Theory of Linear Nonselfadjoint Operators*, Transl. Math. Monogr. 18, Amer. Math. Soc., Providence, RI, 1969.
- [2] P. R. Halmos, *A Hilbert Space Problem Book*, Van Nostrand, Princeton, 1967.
- [3] Yu. Lyubich, *The power boundedness and resolvent conditions for functions of the classical Volterra operator*, *Studia Math.* 196 (2010), 41–63.
- [4] B. Thorpe, *The norm of powers of the indefinite integral operator in $(0, 1)$* , *Bull. London Math. Soc.* 30 (1998), 543–548.
- [5] D. Tsedenbayar, *On the power boundedness of certain Volterra operator pencils*, *Studia Math.* 156 (2003), 59–66.

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