

Dynamics of Schwarz reflections: mating polynomials with groups

Sabyasachi Mukherjee

(Joint work with Seung-Yeop Lee, Mikhail Lyubich, and Nikolai Makarov)

Stony Brook University

Będlewo, March 2018

Motivation from Physics: 2D Coulomb Gas Ensembles (Wiegmann-Zabrodin)

- Consider N electrons placed in the complex plane at points $\{z_j\}_{j=1}^N$,

Motivation from Physics: 2D Coulomb Gas Ensembles (Wiegmann-Zabrodin)

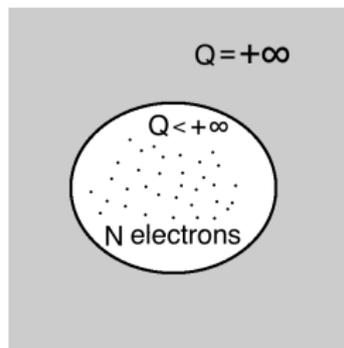
- Consider N electrons placed in the complex plane at points $\{z_j\}_{j=1}^N$, influenced by a strong external (magnetic/electrostatic) field with uniform density.

Motivation from Physics: 2D Coulomb Gas Ensembles (Wiegmann-Zabrodin)

- Consider N electrons placed in the complex plane at points $\{z_j\}_{j=1}^N$, influenced by a strong external (magnetic/electrostatic) field with uniform density. Let the potential of the external field be $NQ : \mathbb{C} \rightarrow \mathbb{R} \cup \{+\infty\}$.

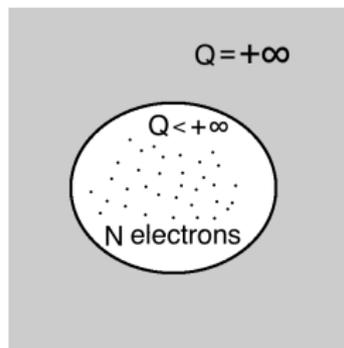
Motivation from Physics: 2D Coulomb Gas Ensembles (Wiegmann-Zabrodin)

- Consider N electrons placed in the complex plane at points $\{z_j\}_{j=1}^N$, influenced by a strong external (magnetic/electrostatic) field with uniform density. Let the potential of the external field be $NQ : \mathbb{C} \rightarrow \mathbb{R} \cup \{+\infty\}$.



Motivation from Physics: 2D Coulomb Gas Ensembles (Wiegmann-Zabrodin)

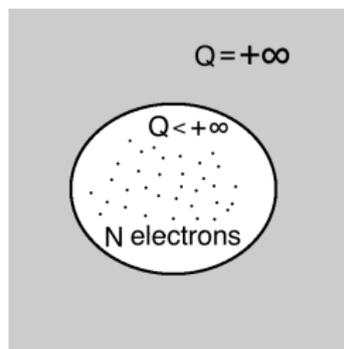
- Consider N electrons placed in the complex plane at points $\{z_j\}_{j=1}^N$, influenced by a strong external (magnetic/electrostatic) field with uniform density. Let the potential of the external field be $NQ : \mathbb{C} \rightarrow \mathbb{R} \cup \{+\infty\}$.



- The combined energy resulting from particle interaction and external potential is:

Motivation from Physics: 2D Coulomb Gas Ensembles (Wiegmann-Zabrodin)

- Consider N electrons placed in the complex plane at points $\{z_j\}_{j=1}^N$, influenced by a strong external (magnetic/electrostatic) field with uniform density. Let the potential of the external field be $NQ : \mathbb{C} \rightarrow \mathbb{R} \cup \{+\infty\}$.



- The combined energy resulting from particle interaction and external potential is:
$$\mathcal{E}_Q(z_1, \dots, z_N) = \sum_{i \neq j} \ln |z_i - z_j|^{-1} + N \sum_{j=1}^N Q(z_j).$$

Large N Behavior of 2D Coulomb Gas Ensembles

- It is more likely to find configurations of electrons with low energy.

Large N Behavior of 2D Coulomb Gas Ensembles

- It is more likely to find configurations of electrons with low energy. This leads to the following joint density of states

$$\frac{\exp(-\mathcal{E}_Q(z_1, \dots, z_N))}{Z_N} dVol_{2N} \in \text{Prob}(\mathbb{C}^N).$$

Large N Behavior of 2D Coulomb Gas Ensembles

- It is more likely to find configurations of electrons with low energy. This leads to the following joint density of states

$$\frac{\exp(-\mathcal{E}_Q(z_1, \dots, z_N))}{Z_N} dVol_{2N} \in \text{Prob}(\mathbb{C}^N).$$

- We are interested in the limiting behavior of the point process as the number of electrons grows to infinity.

Large N Behavior of 2D Coulomb Gas Ensembles

- It is more likely to find configurations of electrons with low energy. This leads to the following joint density of states

$$\frac{\exp(-\mathcal{E}_Q(z_1, \dots, z_N))}{Z_N} dVol_{2N} \in \text{Prob}(\mathbb{C}^N).$$

- We are interested in the limiting behavior of the point process as the number of electrons grows to infinity.
- In the limit, the electrons condensate on a compact set T , and they are distributed according to the normalized area measure of T . (Wiegmann, Zabrodin, Elbau, Felder, Hedenmalm, Makarov, and others.)

Large N Behavior of 2D Coulomb Gas Ensembles

- It is more likely to find configurations of electrons with low energy. This leads to the following joint density of states

$$\frac{\exp(-\mathcal{E}_Q(z_1, \dots, z_N))}{Z_N} dVol_{2N} \in \text{Prob}(\mathbb{C}^N).$$

- We are interested in the limiting behavior of the point process as the number of electrons grows to infinity.
- In the limit, the electrons condensate on a compact set T , and they are distributed according to the normalized area measure of T . (Wiegmann, Zabrodin, Elbau, Felder, Hedenmalm, Makarov, and others.)
- In various physically interesting cases, the external potential Q satisfies some algebraic properties. The compact set T (on which electrons condensate) is then called an *algebraic droplet* of Q .

Large N Behavior of 2D Coulomb Gas Ensembles

- It is more likely to find configurations of electrons with low energy. This leads to the following joint density of states

$$\frac{\exp(-\mathcal{E}_Q(z_1, \dots, z_N))}{Z_N} dVol_{2N} \in \text{Prob}(\mathbb{C}^N).$$

- We are interested in the limiting behavior of the point process as the number of electrons grows to infinity.
- In the limit, the electrons condensate on a compact set T , and they are distributed according to the normalized area measure of T . (Wiegmann, Zabrodin, Elbau, Felder, Hedenmalm, Makarov, and others.)
- In various physically interesting cases, the external potential Q satisfies some algebraic properties. The compact set T (on which electrons condensate) is then called an *algebraic droplet* of Q .
- In these situations, the complementary components of the droplet T admit global reflection maps. (Lee-Makarov)

Definition (Quadrature Domains)

A domain $\Omega \subsetneq \hat{\mathbb{C}}$ with $\infty \notin \partial\Omega$ and $\text{int}(\bar{\Omega}) = \Omega$ is called a *quadrature domain* if there exists a continuous function $\sigma : \bar{\Omega} \rightarrow \hat{\mathbb{C}}$ satisfying the following two properties:

- 1 $\sigma = \text{id}$ on $\partial\Omega$.
- 2 σ is anti-meromorphic on Ω .

Definition (Quadrature Domains)

A domain $\Omega \subsetneq \hat{\mathbb{C}}$ with $\infty \notin \partial\Omega$ and $\text{int}(\bar{\Omega}) = \Omega$ is called a *quadrature domain* if there exists a continuous function $\sigma : \bar{\Omega} \rightarrow \hat{\mathbb{C}}$ satisfying the following two properties:

- 1 $\sigma = \text{id}$ on $\partial\Omega$.
 - 2 σ is anti-meromorphic on Ω .
- The map σ is called the *Schwarz reflection map* of Ω .

From Algebraic Droplets to Quadrature Domains

Definition (Quadrature Domains)

A domain $\Omega \subsetneq \hat{\mathbb{C}}$ with $\infty \notin \partial\Omega$ and $\text{int}(\bar{\Omega}) = \Omega$ is called a *quadrature domain* if there exists a continuous function $\sigma : \bar{\Omega} \rightarrow \hat{\mathbb{C}}$ satisfying the following two properties:

- 1 $\sigma = \text{id}$ on $\partial\Omega$.
 - 2 σ is anti-meromorphic on Ω .
- The map σ is called the *Schwarz reflection map* of Ω .
 - Thus, the study of algebraic droplets is related to the study of quadrature domains.

Definition (Quadrature Domains)

A domain $\Omega \subsetneq \hat{\mathbb{C}}$ with $\infty \notin \partial\Omega$ and $\text{int}(\overline{\Omega}) = \Omega$ is called a *quadrature domain* if there exists a continuous function $\sigma : \overline{\Omega} \rightarrow \hat{\mathbb{C}}$ satisfying the following two properties:

- 1 $\sigma = \text{id}$ on $\partial\Omega$.
 - 2 σ is anti-meromorphic on Ω .
- The map σ is called the *Schwarz reflection map* of Ω .
 - Thus, the study of algebraic droplets is related to the study of quadrature domains. Lee and Makarov used iteration of Schwarz reflection maps to study the topology of quadrature domains.

From Algebraic Droplets to Quadrature Domains

Definition (Quadrature Domains)

A domain $\Omega \subsetneq \hat{\mathbb{C}}$ with $\infty \notin \partial\Omega$ and $\text{int}(\bar{\Omega}) = \Omega$ is called a *quadrature domain* if there exists a continuous function $\sigma : \bar{\Omega} \rightarrow \hat{\mathbb{C}}$ satisfying the following two properties:

- 1 $\sigma = \text{id}$ on $\partial\Omega$.
 - 2 σ is anti-meromorphic on Ω .
- The map σ is called the *Schwarz reflection map* of Ω .
 - Thus, the study of algebraic droplets is related to the study of quadrature domains. Lee and Makarov used iteration of Schwarz reflection maps to study the topology of quadrature domains.
 - Iteration of Schwarz reflections displays features of (anti-)rational maps as well as groups.

Simply Connected Quadrature Domains

Proposition

A simply connected domain $\Omega \subsetneq \hat{\mathbb{C}}$ with $\infty \notin \partial\Omega$ and $\text{int}(\overline{\Omega}) = \Omega$ is a quadrature domain if and only if the Riemann map $\phi : \mathbb{D} \rightarrow \Omega$ is rational.

Simply Connected Quadrature Domains

Proposition

A simply connected domain $\Omega \subsetneq \hat{\mathbb{C}}$ with $\infty \notin \partial\Omega$ and $\text{int}(\overline{\Omega}) = \Omega$ is a quadrature domain if and only if the Riemann map $\phi : \mathbb{D} \rightarrow \Omega$ is rational.

$$\begin{array}{ccc} \overline{\mathbb{D}} & \xrightarrow{\phi} & \overline{\Omega} \\ \downarrow 1/\bar{z} & & \downarrow \sigma \\ \hat{\mathbb{C}} \setminus \mathbb{D} & \xrightarrow{\phi} & \hat{\mathbb{C}} \end{array}$$

Simply Connected Quadrature Domains

Proposition

A simply connected domain $\Omega \subsetneq \hat{\mathbb{C}}$ with $\infty \notin \partial\Omega$ and $\text{int}(\overline{\Omega}) = \Omega$ is a quadrature domain if and only if the Riemann map $\phi : \mathbb{D} \rightarrow \Omega$ is rational.

$$\begin{array}{ccc} \overline{\mathbb{D}} & \xrightarrow{\phi} & \overline{\Omega} \\ \downarrow 1/\bar{z} & & \downarrow \sigma \\ \hat{\mathbb{C}} \setminus \mathbb{D} & \xrightarrow{\phi} & \hat{\mathbb{C}} \end{array}$$

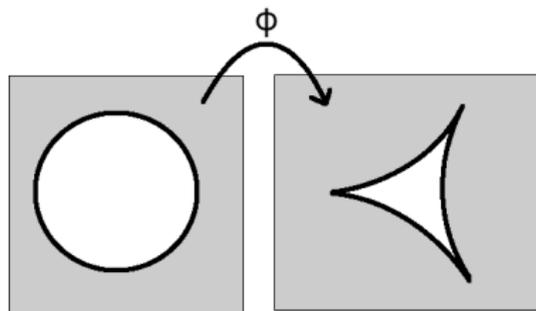
- The rational map ϕ semi-conjugates the reflection map $1/\bar{z}$ of \mathbb{D} to the Schwarz reflection map σ of Ω .

The Deltoid as an Algebraic Droplet

- For the cubic potential $Q(z) = |z|^2 - \Re(z^3)$, the deltoid appears as a “maximal” algebraic droplet.

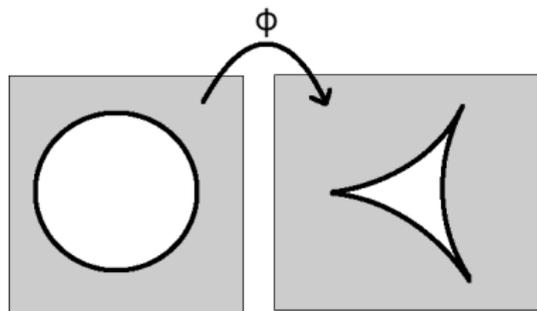
The Deltoid as an Algebraic Droplet

- For the cubic potential $Q(z) = |z|^2 - \Re(z^3)$, the deltoid appears as a “maximal” algebraic droplet.
- The complement of the deltoid has a Riemann map $\phi(z) = z + \frac{1}{2z^2}$, so it is a quadrature domain.



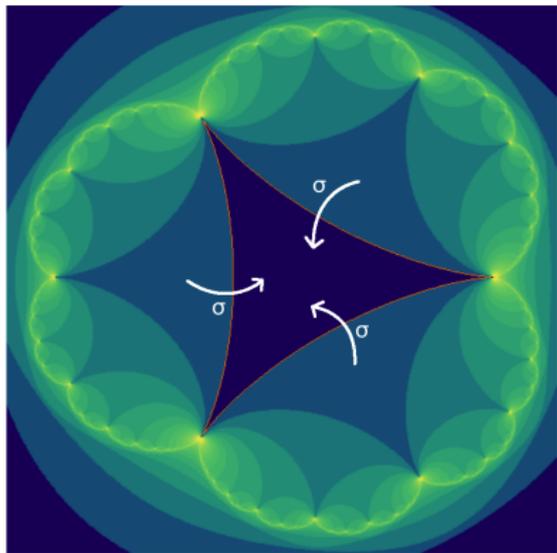
The Deltoid as an Algebraic Droplet

- For the cubic potential $Q(z) = |z|^2 - \Re(z^3)$, the deltoid appears as a “maximal” algebraic droplet.
- The complement of the deltoid has a Riemann map $\phi(z) = z + \frac{1}{2z^2}$, so it is a quadrature domain.

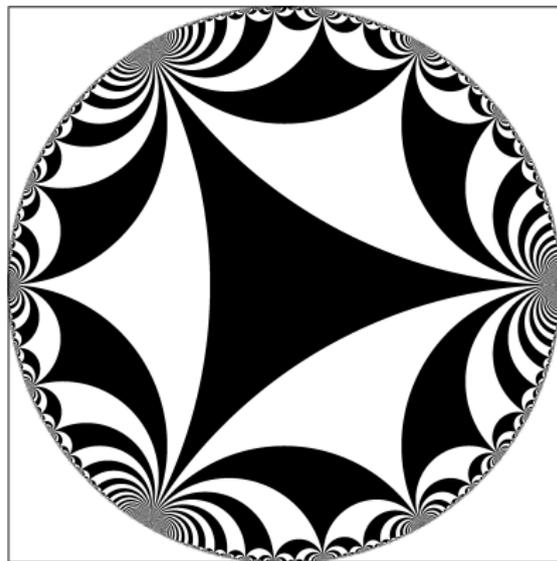
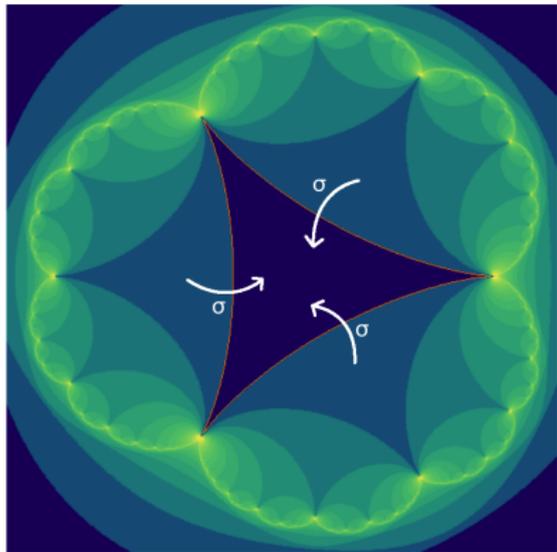


- The corresponding Schwarz reflection map is unicritical, and has a super-attracting fixed point at ∞ .

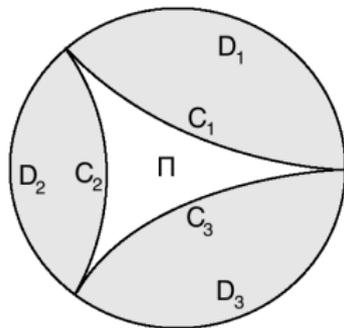
Dynamics of Deltoid Reflection



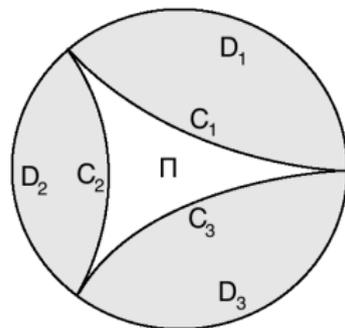
Dynamics of Deltoid Reflection



Deltoid Reflection as a Mating

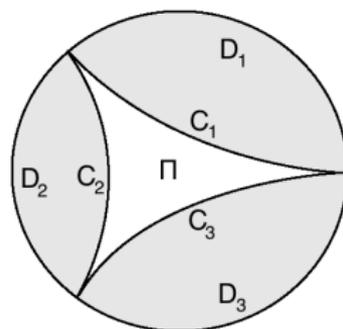


Deltoid Reflection as a Mating



- Denote reflection w.r.t. C_j by ρ_j .

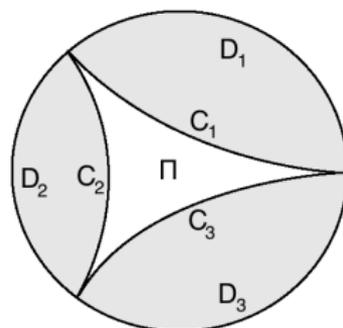
Deltoid Reflection as a Mating



- Denote reflection w.r.t. C_j by ρ_j .
- Define the reflection map $\rho : \mathbb{D} \setminus \Pi \rightarrow \mathbb{D}$ as:

$$z \mapsto \begin{cases} \rho_1(z) & \text{if } z \in D_1, \\ \rho_2(z) & \text{if } z \in D_2, \\ \rho_3(z) & \text{if } z \in D_3. \end{cases}$$

Deltoid Reflection as a Mating

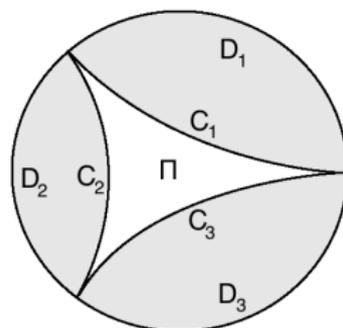


- Denote reflection w.r.t. C_j by ρ_j .
- Define the reflection map $\rho : \mathbb{D} \setminus \Pi \rightarrow \mathbb{D}$ as:

$$z \mapsto \begin{cases} \rho_1(z) & \text{if } z \in D_1, \\ \rho_2(z) & \text{if } z \in D_2, \\ \rho_3(z) & \text{if } z \in D_3. \end{cases}$$

- The map ρ extends to an orientation-reversing double cover of \mathbb{S}^1 .

Deltoid Reflection as a Mating



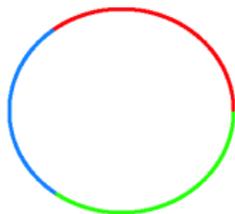
- Denote reflection w.r.t. C_j by ρ_j .
- Define the reflection map $\rho : \mathbb{D} \setminus \Pi \rightarrow \mathbb{D}$ as:

$$z \mapsto \begin{cases} \rho_1(z) & \text{if } z \in D_1, \\ \rho_2(z) & \text{if } z \in D_2, \\ \rho_3(z) & \text{if } z \in D_3. \end{cases}$$

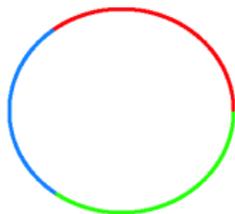
- The map ρ extends to an orientation-reversing double cover of \mathbb{S}^1 .
- The dynamics of the deltoid reflection map is a “mating” of ρ (on the tiling set) and \bar{z}^2 (on the non-escaping set).

- The orientation-reversing double coverings ρ and \bar{z}^2 (of \mathbb{S}^1) admit a common Markov partition.

- The orientation-reversing double coverings ρ and \bar{z}^2 (of \mathbb{S}^1) admit a common Markov partition.



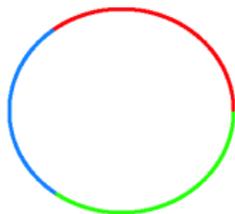
- The orientation-reversing double coverings ρ and \bar{z}^2 (of \mathbb{S}^1) admit a common Markov partition.



- Moreover, they have the same transition matrix

$$M := \begin{bmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{bmatrix}.$$

- The orientation-reversing double coverings ρ and \bar{z}^2 (of \mathbb{S}^1) admit a common Markov partition.

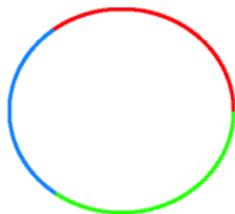


- Moreover, they have the same transition matrix

$$M := \begin{bmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{bmatrix}.$$

- Consequently, ρ and \bar{z}^2 are topologically conjugate by a circle homeomorphism \mathcal{H} .

- The orientation-reversing double coverings ρ and \bar{z}^2 (of \mathbb{S}^1) admit a common Markov partition.



- Moreover, they have the same transition matrix

$$M := \begin{bmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{bmatrix}.$$

- Consequently, ρ and \bar{z}^2 are topologically conjugate by a circle homeomorphism \mathcal{H} .
- \mathcal{H} "welds" the external class of quadratic antiholomorphic polynomials and that of the ideal triangle group.

The Circle and Cardioid Family

- Let \heartsuit be the principal hyperbolic component of the Mandelbrot set. Note that \heartsuit is a quadrature domain.

The Circle and Cardioid Family

- Let \heartsuit be the principal hyperbolic component of the Mandelbrot set. Note that \heartsuit is a quadrature domain.
- For any $a \in \mathbb{C} \setminus (-\infty, -1/12)$, the smallest disk $B(a, r_a)$ containing \heartsuit touches $\partial\heartsuit$ at a unique point.

The Circle and Cardioid Family

- Let \heartsuit be the principal hyperbolic component of the Mandelbrot set. Note that \heartsuit is a quadrature domain.
- For any $a \in \mathbb{C} \setminus (-\infty, -1/12)$, the smallest disk $B(a, r_a)$ containing \heartsuit touches $\partial\heartsuit$ at a unique point.
- Let $\Omega_a := \heartsuit \cup \overline{B}(a, r_a)^c$. We now define our dynamical system $F_a : \overline{\Omega}_a \rightarrow \hat{\mathbb{C}}$ as,

$$w \mapsto \begin{cases} \sigma(w) & \text{if } w \in \overline{\heartsuit}, \\ \sigma_a(w) & \text{if } w \in B(a, r_a)^c, \end{cases}$$

where σ is the Schwarz reflection of \heartsuit , and σ_a is reflection with respect to the circle $|w - a| = r_a$.

The Circle and Cardioid Family

- Let \heartsuit be the principal hyperbolic component of the Mandelbrot set. Note that \heartsuit is a quadrature domain.
- For any $a \in \mathbb{C} \setminus (-\infty, -1/12)$, the smallest disk $B(a, r_a)$ containing \heartsuit touches $\partial\heartsuit$ at a unique point.
- Let $\Omega_a := \heartsuit \cup \overline{B}(a, r_a)^c$. We now define our dynamical system $F_a : \overline{\Omega}_a \rightarrow \hat{\mathbb{C}}$ as,

$$w \mapsto \begin{cases} \sigma(w) & \text{if } w \in \overline{\heartsuit}, \\ \sigma_a(w) & \text{if } w \in B(a, r_a)^c, \end{cases}$$

where σ is the Schwarz reflection of \heartsuit , and σ_a is reflection with respect to the circle $|w - a| = r_a$.

- 0 is the only critical point of F_a .

The Circle and Cardioid Family

- Let \heartsuit be the principal hyperbolic component of the Mandelbrot set. Note that \heartsuit is a quadrature domain.
- For any $a \in \mathbb{C} \setminus (-\infty, -1/12)$, the smallest disk $B(a, r_a)$ containing \heartsuit touches $\partial\heartsuit$ at a unique point.
- Let $\Omega_a := \heartsuit \cup \overline{B}(a, r_a)^c$. We now define our dynamical system $F_a : \overline{\Omega}_a \rightarrow \hat{\mathbb{C}}$ as,

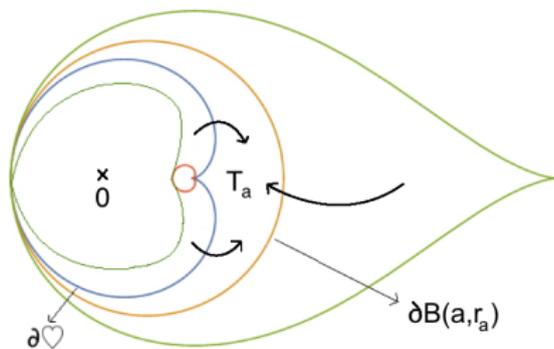
$$w \mapsto \begin{cases} \sigma(w) & \text{if } w \in \overline{\heartsuit}, \\ \sigma_a(w) & \text{if } w \in B(a, r_a)^c, \end{cases}$$

where σ is the Schwarz reflection of \heartsuit , and σ_a is reflection with respect to the circle $|w - a| = r_a$.

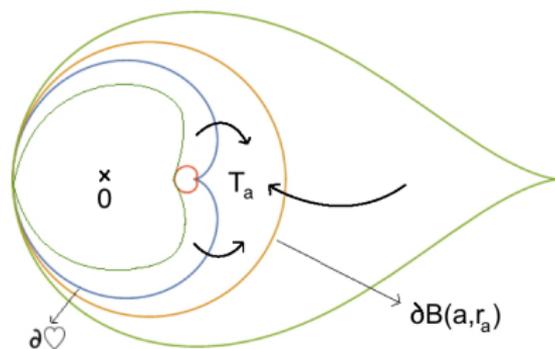
- 0 is the only critical point of F_a .
- We will call this family of maps \mathcal{S} ; i.e.

$$\mathcal{S} := \{F_a : \overline{\Omega}_a \rightarrow \hat{\mathbb{C}} : a \in \mathbb{C} \setminus (-\infty, -1/12)\}.$$

The Circle and Cardioid Family

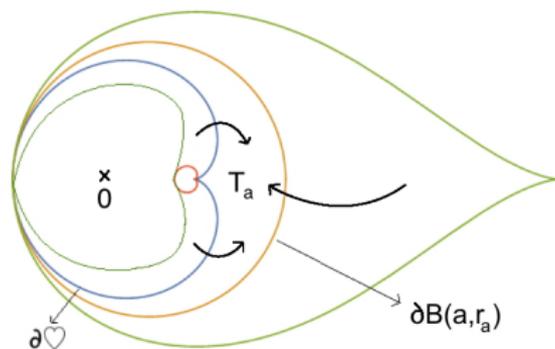


The Circle and Cardioid Family



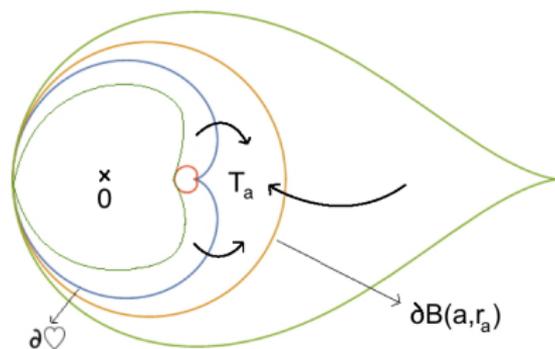
- The *tiling set* T_a^∞ of F_a is defined as the set of points in $\overline{\Omega}_a$ that eventually escape to T_a ; i.e. $T_a^\infty = \bigcup_{k=0}^{\infty} F_a^{-k}(T_a^0)$.

The Circle and Cardioid Family



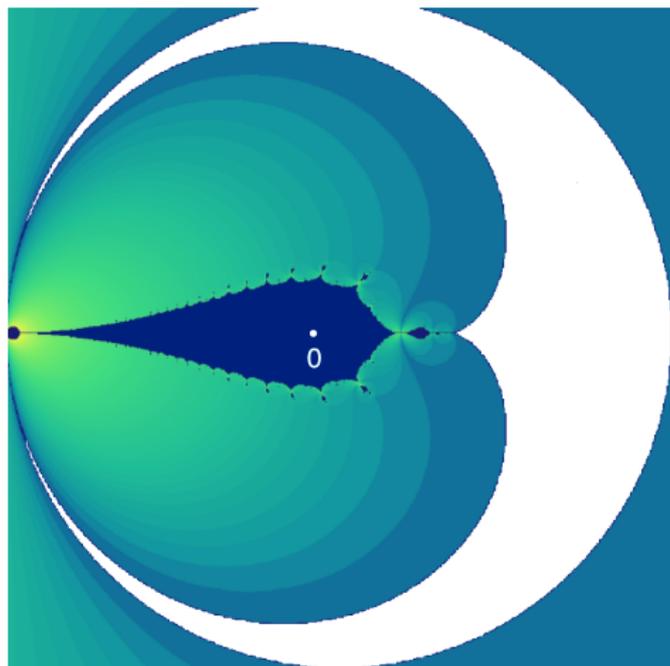
- The *tiling set* T_a^∞ of F_a is defined as the set of points in $\overline{\Omega}_a$ that eventually escape to T_a ; i.e. $T_a^\infty = \bigcup_{k=0}^{\infty} F_a^{-k}(T_a^0)$.
- The *non-escaping set* K_a of F_a is the complement $\hat{\mathbb{C}} \setminus T_a^\infty$. Connected components of $\text{int}(K_a)$ are called *Fatou components* of F_a .

The Circle and Cardioid Family

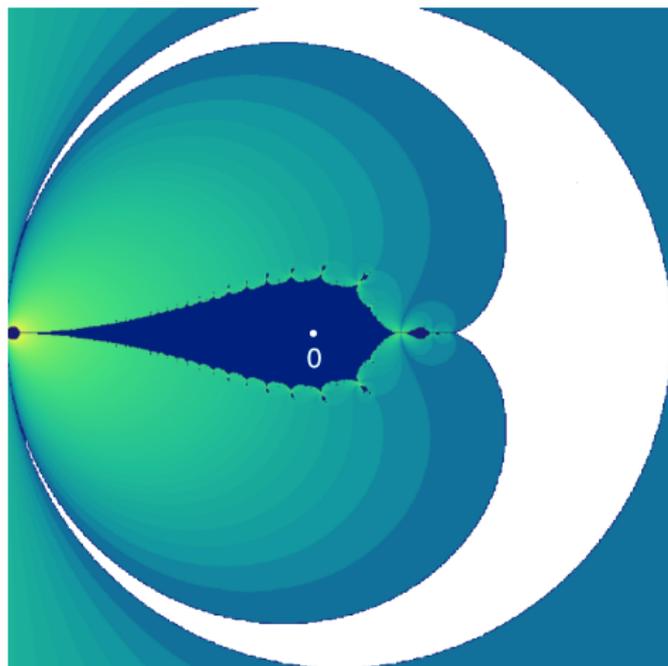


- The *tiling set* T_a^∞ of F_a is defined as the set of points in $\overline{\Omega}_a$ that eventually escape to T_a ; i.e. $T_a^\infty = \bigcup_{k=0}^{\infty} F_a^{-k}(T_a^0)$.
- The *non-escaping set* K_a of F_a is the complement $\hat{\mathbb{C}} \setminus T_a^\infty$. Connected components of $\text{int}(K_a)$ are called *Fatou components* of F_a .
- The boundary of T_a^∞ is called the *limit set* of F_a , and is denoted by Γ_a .

Dynamical Plane of the Basilica Map

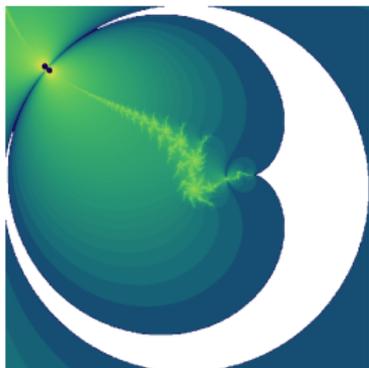


Dynamical Plane of the Basilica Map

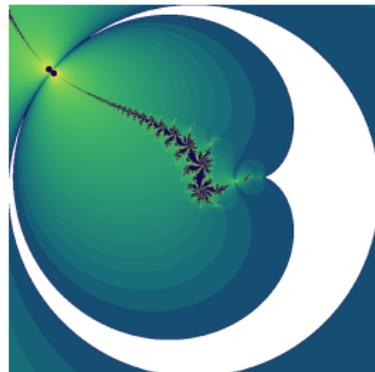
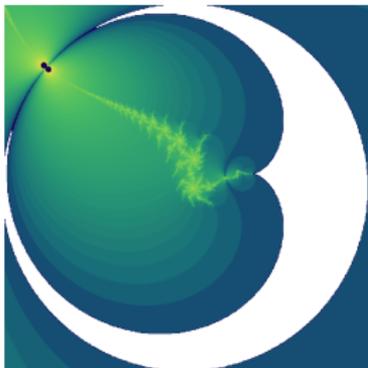


- $0 \mapsto \infty \mapsto 0$; the "Basilica" map.

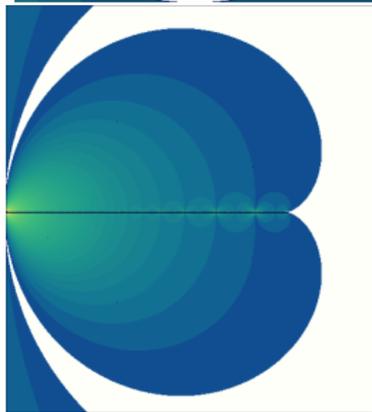
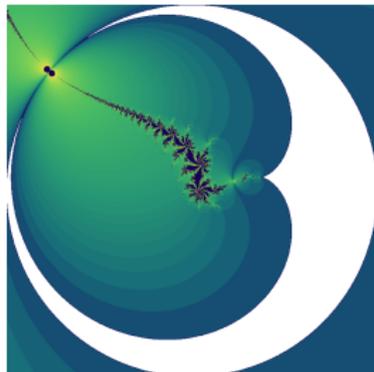
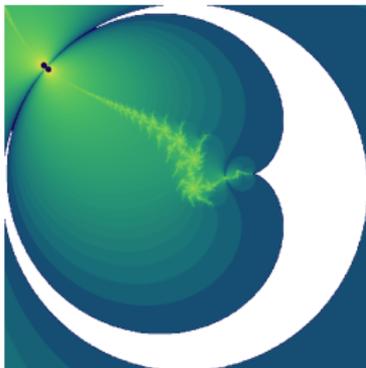
Various Dynamical Planes



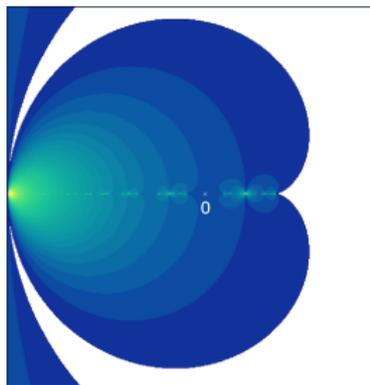
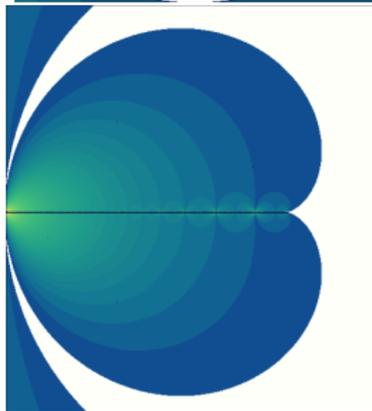
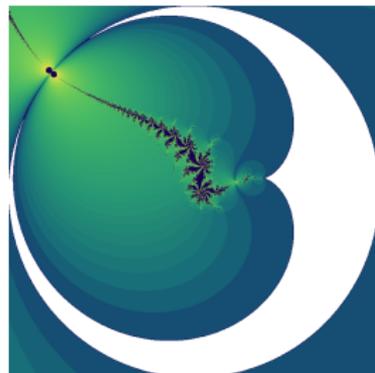
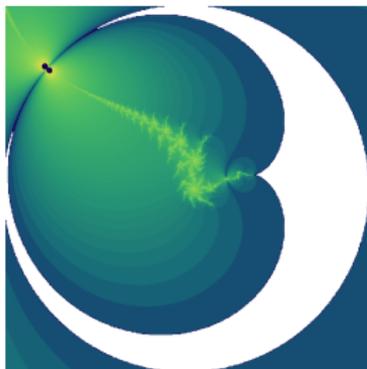
Various Dynamical Planes



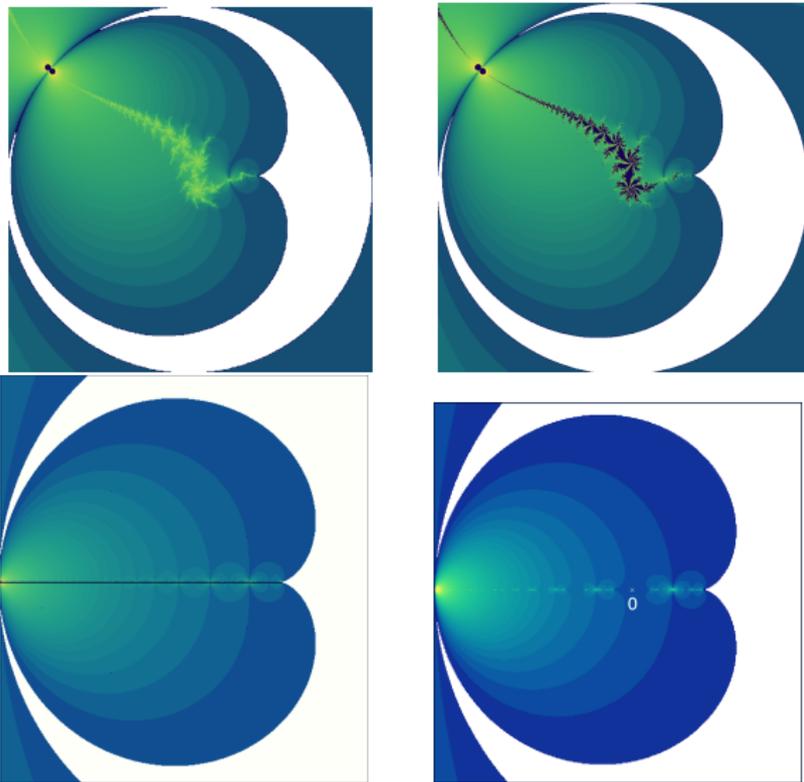
Various Dynamical Planes



Various Dynamical Planes



Various Dynamical Planes



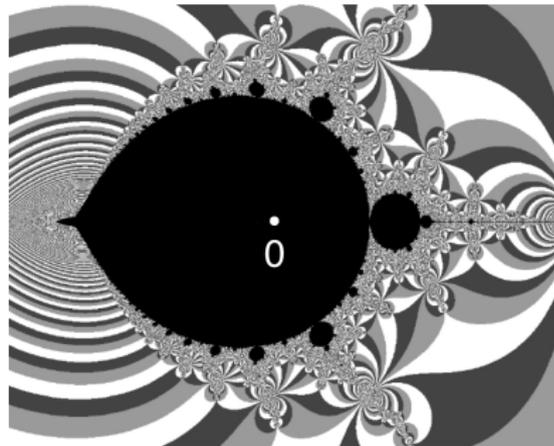
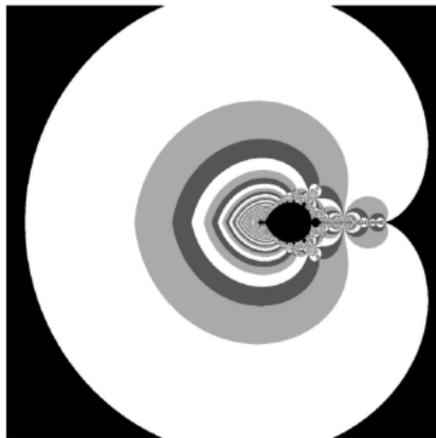
- “Full tessellation structure” iff tiles are unramified; i.e. $0 \notin T_a^\infty$.

The Connectedness Locus $\mathcal{C}(\mathcal{S})$

- $\mathcal{C}(\mathcal{S}) = \{a \in \mathbb{C} \setminus (-\infty, -1/12) : K_a \text{ is connected}\}$
 $= \{a \in \mathbb{C} \setminus (-\infty, -1/12) : 0 \notin T_a^\infty\}.$

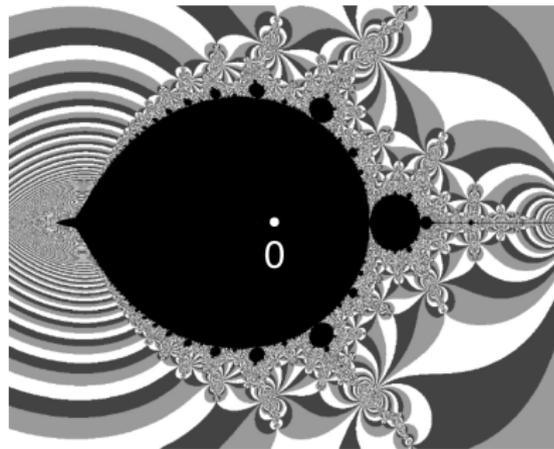
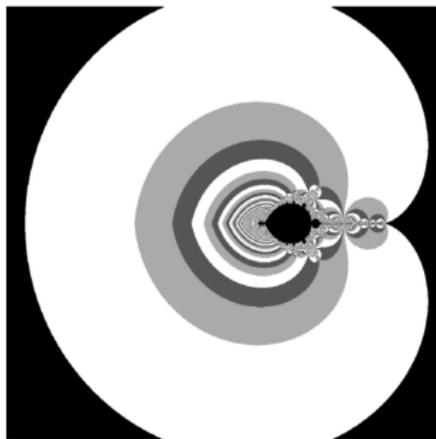
The Connectedness Locus $\mathcal{C}(\mathcal{S})$

- $\mathcal{C}(\mathcal{S}) = \{a \in \mathbb{C} \setminus (-\infty, -1/12) : K_a \text{ is connected}\}$
 $= \{a \in \mathbb{C} \setminus (-\infty, -1/12) : 0 \notin T_a^\infty\}$.



The Connectedness Locus $\mathcal{C}(\mathcal{S})$

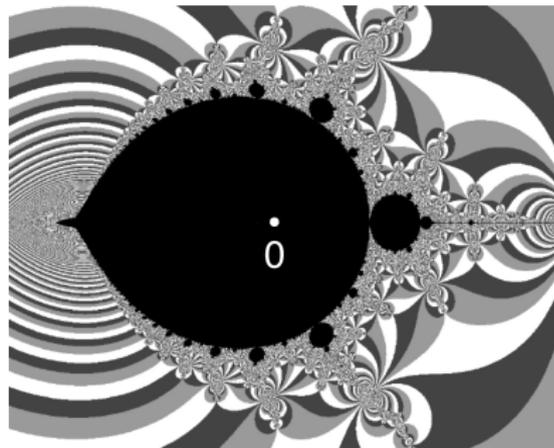
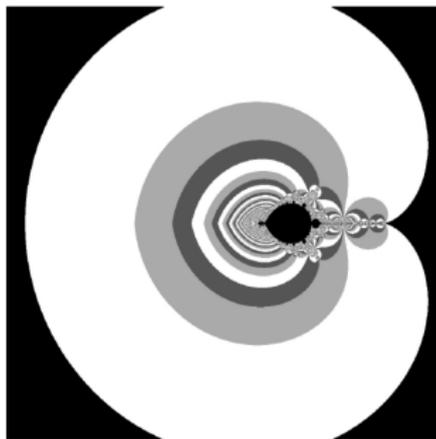
- $\mathcal{C}(\mathcal{S}) = \{a \in \mathbb{C} \setminus (-\infty, -1/12) : K_a \text{ is connected}\}$
 $= \{a \in \mathbb{C} \setminus (-\infty, -1/12) : 0 \notin T_a^\infty\}$.



- For maps in $\mathcal{C}(\mathcal{S})$, the dynamics on the tiling set is conformally conjugate to the reflection map ρ . This is our analogue of Böttcher coordinates.

The Connectedness Locus $\mathcal{C}(\mathcal{S})$

- $\mathcal{C}(\mathcal{S}) = \{a \in \mathbb{C} \setminus (-\infty, -1/12) : K_a \text{ is connected}\}$
 $= \{a \in \mathbb{C} \setminus (-\infty, -1/12) : 0 \notin T_a^\infty\}$.



- For maps in $\mathcal{C}(\mathcal{S})$, the dynamics on the tiling set is conformally conjugate to the reflection map ρ . This is our analogue of Böttcher coordinates.
- For maps outside $\mathcal{C}(\mathcal{S})$, the “Böttcher coordinate” is defined on low generation tiles.

External Dynamics and Lamination Model of Limit Sets

- If F_a is post-critically finite, then the limit set Γ_a is locally connected.

External Dynamics and Lamination Model of Limit Sets

- If F_a is post-critically finite, then the limit set Γ_a is locally connected.
- Hence, the external conjugacy extends continuously to \mathbb{S}^1 .

External Dynamics and Lamination Model of Limit Sets

- If F_a is post-critically finite, then the limit set Γ_a is locally connected.
- Hence, the external conjugacy extends continuously to \mathbb{S}^1 .

$$\begin{array}{ccc} \Gamma_a & \xrightarrow{\text{Böttcher}} & (\mathbb{R}/\mathbb{Z})/\sim \\ \downarrow F_a & & \downarrow \rho \\ \Gamma_a & \xrightarrow{\text{Böttcher}} & (\mathbb{R}/\mathbb{Z})/\sim \end{array}$$

External Dynamics and Lamination Model of Limit Sets

- If F_a is post-critically finite, then the limit set Γ_a is locally connected.
- Hence, the external conjugacy extends continuously to \mathbb{S}^1 .

$$\begin{array}{ccc} \Gamma_a & \xrightarrow{\text{Böttcher}} & (\mathbb{R}/\mathbb{Z})/\sim \\ \downarrow F_a & & \downarrow \rho \\ \Gamma_a & \xrightarrow{\text{Böttcher}} & (\mathbb{R}/\mathbb{Z})/\sim \end{array}$$

- This yields a pinched disk model of the non-escaping set K_a via a ρ -invariant geodesic lamination of \mathbb{D} .

External Dynamics and Lamination Model of Limit Sets

- If F_a is post-critically finite, then the limit set Γ_a is locally connected.
- Hence, the external conjugacy extends continuously to \mathbb{S}^1 .

$$\begin{array}{ccc} \Gamma_a & \xrightarrow{\text{Böttcher}} & (\mathbb{R}/\mathbb{Z})/\sim \\ \downarrow F_a & & \downarrow \rho \\ \Gamma_a & \xrightarrow{\text{Böttcher}} & (\mathbb{R}/\mathbb{Z})/\sim \end{array}$$

- This yields a pinched disk model of the non-escaping set K_a via a ρ -invariant geodesic lamination of \mathbb{D} .
- **Rigidity: PCF maps are uniquely determined by laminations.** The proof is based on the “Pullback Argument”, and involves an analysis of the boundary behavior of conformal maps near cusps and double points.

Bijection between PCF Parameters

Theorem

There exists a natural combinatorial bijection between the PCF parameters of \mathcal{S} and PCF anti-quadratic polynomials such that the laminations of the corresponding maps are related by the circle homeomorphism \mathcal{H} .

Bijection between PCF Parameters

Theorem

There exists a natural combinatorial bijection between the PCF parameters of \mathcal{S} and PCF anti-quadratic polynomials such that the laminations of the corresponding maps are related by the circle homeomorphism \mathcal{H} .

- The circle homeomorphism \mathcal{H} pushes forward ρ -invariant laminations to “formal rational laminations”. This defines a map from PCF Schwarz maps to PCF anti-quadratic polynomials (anti-holomorphic version of Kiwi’s theorem).

Bijection between PCF Parameters

Theorem

There exists a natural combinatorial bijection between the PCF parameters of S and PCF anti-quadratic polynomials such that the laminations of the corresponding maps are related by the circle homeomorphism \mathcal{H} .

- The circle homeomorphism \mathcal{H} pushes forward ρ -invariant laminations to “formal rational laminations”. This defines a map from PCF Schwarz maps to PCF anti-quadratic polynomials (anti-holomorphic version of Kiwi’s theorem).
- Injectivity: Rigidity of PCF maps.

Bijection between PCF Parameters

Theorem

There exists a natural combinatorial bijection between the PCF parameters of \mathcal{S} and PCF anti-quadratic polynomials such that the laminations of the corresponding maps are related by the circle homeomorphism \mathcal{H} .

- The circle homeomorphism \mathcal{H} pushes forward ρ -invariant laminations to “formal rational laminations”. This defines a map from PCF Schwarz maps to PCF anti-quadratic polynomials (anti-holomorphic version of Kiwi’s theorem).
- Injectivity: Rigidity of PCF maps.
- Surjectivity: Finding PCF maps in \mathcal{S} with prescribed laminations. We achieve this “from outside”; i.e. as limit points of suitable sequences of parameter tiles.

Theorem

Every PCF map F_a is a mating of the PCF quadratic anti-holomorphic polynomial $f_{\chi(a)}$ and the reflection map ρ . More precisely, F_a is topologically conjugate to $f_{\chi(a)}$ on its non-escaping set, and conformally conjugate to ρ on its tiling set.

Theorem

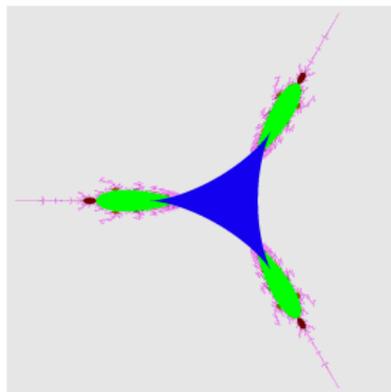
Every PCF map F_a is a mating of the PCF quadratic anti-holomorphic polynomial $f_{\chi(a)}$ and the reflection map ρ . More precisely, F_a is topologically conjugate to $f_{\chi(a)}$ on its non-escaping set, and conformally conjugate to ρ on its tiling set. The “welding” map is a factor of \mathcal{H} .

A Quick Tour of the Tricorn

- The *Tricorn* is the connectedness locus of anti-quadratic polynomials $\bar{z}^2 + c$.

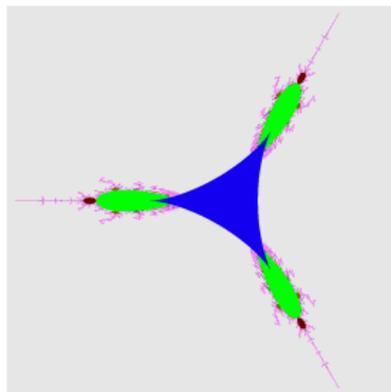
A Quick Tour of the Tricorn

- The *Tricorn* is the connectedness locus of anti-quadratic polynomials $\bar{z}^2 + c$.



A Quick Tour of the Tricorn

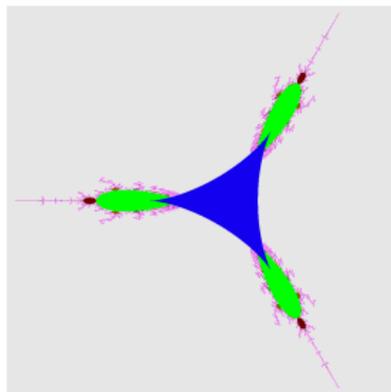
- The *Tricorn* is the connectedness locus of anti-quadratic polynomials $\bar{z}^2 + c$.



- The tricorn is not locally connected. (Hubbard-Schleicher)

A Quick Tour of the Tricorn

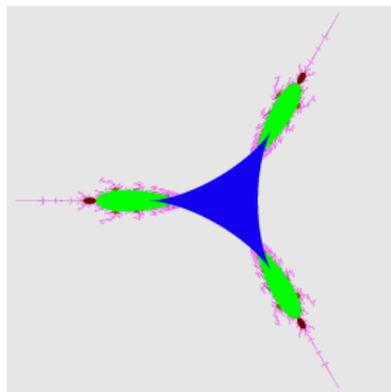
- The *Tricorn* is the connectedness locus of anti-quadratic polynomials $\bar{z}^2 + c$.



- The tricorn is not locally connected. (Hubbard-Schleicher)
- Many rational parameter rays non-trivially accumulate.

A Quick Tour of the Tricorn

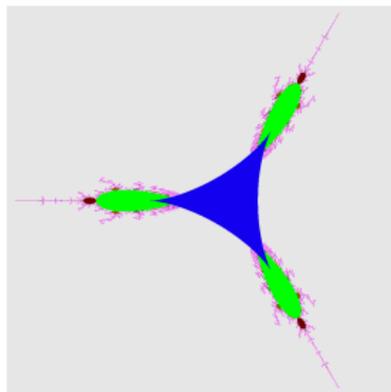
- The *Tricorn* is the connectedness locus of anti-quadratic polynomials $\bar{z}^2 + c$.



- The tricorn is not locally connected. (Hubbard-Schleicher)
- Many rational parameter rays non-trivially accumulate. Misiurewicz parameters are not dense on the boundary. (Inou-M)

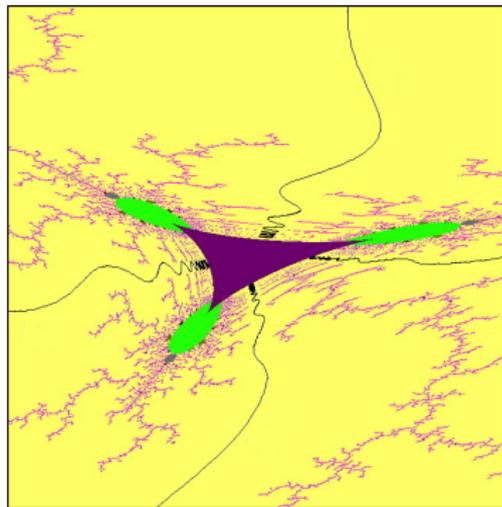
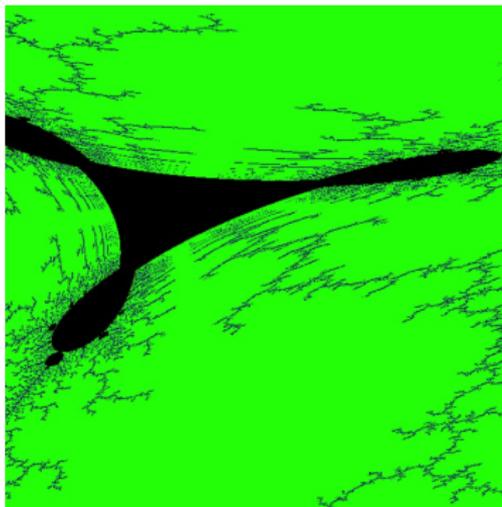
A Quick Tour of the Tricorn

- The *Tricorn* is the connectedness locus of anti-quadratic polynomials $\bar{z}^2 + c$.



- The tricorn is not locally connected. (Hubbard-Schleicher)
- Many rational parameter rays non-trivially accumulate. Misiurewicz parameters are not dense on the boundary. (Inou-M)
- “Baby tricorns” are not homeomorphic to the original tricorn. (Inou-M)

Topology of the Tricorn



Homeomorphism between Models

Theorem

The lamination model of $\mathcal{C}(\mathcal{F})$ is homeomorphic to that of the basilica limb of the Tricorn.

Theorem

The lamination model of $\mathcal{C}(\mathcal{F})$ is homeomorphic to that of the basilica limb of the Tricorn.

- Landing/accumulation patterns of parameter rays precisely correspond under \mathcal{H} .

Theorem

The lamination model of $\mathcal{C}(\mathcal{F})$ is homeomorphic to that of the basilica limb of the Tricorn.

- Landing/accumulation patterns of parameter rays precisely correspond under \mathcal{H} .
- This descends to a homeomorphism between the pinched disk models of connectedness loci.

Thank you!