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Weak bounds for the maximal function in weighted Orlicz spaces

by

RICHARD J. BAGBY (Las Cruces, N.Mex.)

Abstract. For Φ a Young function and $wE = \int_E w(x) dx$ a nontrivial doubling measure on E^n , it is shown that the Hardy-Littlewood maximal function Mf satisfies the weak-type inequality

$$w\{x: Mf(x) > \lambda\} \leq \int \Phi(c_1|f(x)|/\lambda)w(x) dx$$

if and only if the conjugate Young function Φ^* satisfies

$$\int_{Q} \Phi^{\sim} (\varepsilon w Q/|Q|w(x)) w(x) dx \le wQ$$

for all cubes Q, with some fixed $\varepsilon > 0$. When Φ^{\sim} is submultiplicative, it is shown that the last inequality can always be strengthened by replacing Φ^{\sim} by $(\Phi^{\sim})^{1+\delta}$ for some $\delta > 0$. In some weighted Orlicz classes (L' spaces, for example), this can be used to prove the equivalence of the weak inequality for Mf and the strong inequality

$$\int \Phi(Mf(x))w(x) dx \leq \int \Phi(c_2|f(x)|)w(x)dx,$$

but we show the equivalence is generally false.

1. Introduction. In [6], B. Muckenhoupt found a remarkably simple characterization of the class A_p of weight functions w such that the Hardy-Littlewood maximal function is bounded in $L^p(E^n, w(x) dx)$. For $1 , he proved that <math>w \in A_p$ if and only if there is a constant c such that

$$\left(\int\limits_{Q}w(x)\,dx\right)\left(\int\limits_{Q}w(x)^{-1/(p-1)}\,dx\right)^{p-1}\leqslant c|Q|^{p}$$

for all cubes Q in E^n ; this inequality is known as the A_p condition. Routine arguments show that the A_p condition is equivalent to a weak-type bound

$$\int\limits_{\{x:\,Mf(x)>\lambda\}}w(x)\,dx\leqslant c\lambda^{-p}\int|f(x)|^p\,dx$$

for arbitrary functions f and for all $\lambda > 0$, but in general an inequality of this type is strictly weaker than continuity in $L^p(E^n, w(x)dx)$. However, for the maximal function such an inequality does imply continuity in $L^q(E^n, w(x)dx)$ for all q > p. Muckenhoupt proved that whenever w satisfies the A_p condition, it must also satisfy the $A_{p-\varepsilon}$ condition for sufficiently small $\varepsilon > 0$, thereby proving that the maximal function is continuous in $L^p(E^n, w(x)dx)$.

Kerman and Torchinsky [4] then extended Muckenhoupt's results by considering weighted Orlicz spaces. They defined the class A_{Φ} of weight

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functions to be those for which the maximal function satisfies a bound

$$\int \Phi(Mf(x))w(x)\,dx \leqslant \int \Phi(c|f(x)|)w(x)\,dx$$

and proved that $w \in A_{\Phi}$ if and only if $w \in A_p$ for 1/p the upper index of Φ . Chung, Hunt, and Kurtz [2] found a refinement of Muckenhoupt's result by considering weighted Lorentz space norms. They found that weights satisfying the A(p, 1) condition

$$\int_{Q} w(x) \, dx / \int_{Q \cap E} w(x) \, dx \le c(|Q|/|Q \cap E|)^{p}$$

for arbitrary measurable sets E and cubes Q satisfy the same weak-type inequality as A_p weights provided f is a characteristic function, but not for general $f \in L^p(E^n, wdx)$. Consequently, A(p, 1) is a proper subset of A_p which also contains A_q for all q > p.

Here we present a significant further refinement of Muckenhoupt's work by considering weak-type inequalities for the Hardy-Littlewood maximal function in weighted Orlicz spaces. We define the class B_{Φ} to be those weight functions w satisfying

$$\int_{\{x: Mf(x) > \lambda\}} w(x) dx \leq \int \Phi(c|f(x)|/\lambda) w(x) dx$$

for all $\lambda > 0$, $f \in L^{\Phi}(E^n, w(x)dx)$. We prove that B_{Φ} can be characterized by a natural analogue of the A_p condition, and prove a theorem generalizing Muckenhoupt's result that $w \in A_p$ implies $w \in A_{p-\varepsilon}$. However, this last result requires a hypothesis ruling out choices of Φ such as $\Phi(t) = ct^p (1 + \log_+ t)^k$ with k > 0. In fact, our characterization of B_{Φ} shows that these are distinct proper subclasses of A_p containing A_q for all q > p.

Our characterization of B_{Φ} is simple enough that for a given weight w, we may by able to find a nearly optimal Φ for which $w \in B_{\Phi}$. By using the techniques of J. D. Parsons and the author [1], we can then obtain extremely delicate weighted bounds for the maximal function. For example, $w \in B_{\Phi}$ with $\Phi(t) = t^2 (\log_+ t)^2$ does not imply $w \in A_2$, but it does imply that the Hardy-Littlewood maximal function is bounded from $L^2(\log L)^3(wdx)$ to $L^2_{loc}(wdx)$. Such results were not previously obtainable.

2. Preliminaries. Let w be a fixed nonnegative, measurable function on E^n . For $A \subset E^n$ Lebesgue measurable, we shall write |A| for the Lebesgue measure of A and

$$wA = \int_A w(x) \, dx.$$

We shall generally use the letter Q for a cube in E^n , by which we mean the product of n intervals $[a_i, a_i + s]$ with $0 < s < \infty$. For Q such a cube, we write

 Q^* for the cube having the same center as Q, but edges three times as long. We then say that w defines a doubling measure if there is a constant c such that

$$wQ^* \leq cwQ$$
 for all cubes Q .

We use the version of the Hardy-Littlewood maximal function on E^n defined by

$$Mf(x) = \sup\{|Q|^{-1} \int_{Q} |f(y)| dy: x \in Q\}.$$

The simple result below allows us to bound the distribution function of Mf with respect to doubling Borel measures. The proof is quite well known.

(2.1) PROPOSITION. For K an arbitrary compact subset of $\{x: Mf(x) > \lambda\}$, it is possible to choose cubes Q_1, \ldots, Q_m with pairwise disjoint interiors such that

$$\int_{Q_k} |f(x)| dx > \lambda |Q_k|, \quad \text{for each } Q_k, \text{ and}$$

$$K \subset \bigcup Q_k^*$$

Proof. For each $y \in K$, there is a cube Q having x in its interior and $\int_{Q} |f(x)| dx > \lambda |Q|$; by compactness we can choose finitely many such cubes covering K. We may then order this finite collection so that their measures form a nonincreasing sequence.

Let Q_1 be the first cube in this sequence, and let Q_2 be the next cube with interior disjoint from that of Q_1 . Having selected Q_1, \ldots, Q_k with pairwise disjoint interiors, choose Q_{k+1} to be the next cube in the cover with this disjointness property.

If Q is not selected, then $|Q \cap Q_i| > 0$ for some selected Q_i with $|Q_i| \ge |Q|$. Consequently, $Q \subset Q_i^*$ and the proof is complete.

We shall call a function Φ on $[0, \infty)$ a Young function if Φ is continuous, convex, nondecreasing, and satisfies

$$\Phi(0) = 0$$
, $\Phi(t)/t \to \infty$ as $t \to \infty$.

The conjugate Young function Φ^{-} may be defined by

$$\Phi^{\sim}(t) = \sup\{st - \Phi(s): s > 0\};$$

it is the minimal Young function which satisfies $st \leq \Phi(s) + \Phi^{\sim}(t)$.

For Φ a Young function, the *Orlicz space* $L^{\Phi}(d\mu)$ consists of all μ -measurable functions f such that $\Phi(\varepsilon|f|)$ is μ -integrable for some $\varepsilon > 0$. It can be normed by

$$||f||_{\Phi} = \inf\{\lambda > 0 \colon \int \Phi(|f|/\lambda) \, d\mu \leqslant 1\}.$$

The Orlicz space version of Hölder's inequality is then

$$\int |fg| d\mu \leq 2 \|f\|_{\mathbf{\Phi}} \|g\|_{\mathbf{\Phi}^{\sim}}.$$

Basic facts about Orlicz spaces may be found in Zygmund [7]; an extensive treatment is given by Krasnosel'skii and Rutickii [5].

Note that monotone convergence implies that

$$\int \Phi(|f|/\lambda) d\mu \le 1 \quad \text{for } \lambda = ||f||_{\Phi};$$

inequality is possible unless $\Phi(|f|/\lambda)$ is μ -integrable for some $\lambda < \|f\|_{\Phi}$. For functions on E^n which are locally in $L^{\Phi}(wdx)$, the norms

$$||f||_{\Phi,w,Q} = \inf\{\lambda > 0: \int_{Q} \Phi(|f|/\lambda) d\mu \le wQ\}$$

are quite useful. For $0 < wQ < \infty$, this is the usual norm in $L^{\Phi}(d\mu)$ with $d\mu$ the restriction to Q of (w/wQ)dx.

We conclude this section with a technical lemma relating conjugate Young functions.

(2.2) LEMMA. For Φ and Φ^{\sim} conjugate Young functions, there is a continuous, monotone function g on $[0, \infty)$ such that

$$\Phi(g(t)) \leqslant tg(t) \leqslant \Phi(2g(t)),$$

$$2\Phi^{\sim}(t/2) \leqslant tg(t) \leqslant \Phi^{\sim}(2t).$$

Proof. If we define

$$h(t) = \sup\{s \colon \Phi(s) \leqslant st\},\,$$

then h is monotone and satisfies $\Phi(h(t)) = th(t)$. Moreover,

$$\Phi^{\sim}(t) = \sup\{st - \Phi(s): \Phi(s) \leqslant st\} \leqslant th(t),$$

while $th(t) = 2th(t) - \Phi(h(t)) \le \Phi^{\sim}(2t)$. However, h has a discontinuity if $\Phi(t)/t$ is a nonzero constant near t = 0.

Let us define

$$\Phi_{0}(t) = \begin{cases} \int_{0}^{t} (1+s)\Phi'(s) ds, & t \leq 1, \\ \Phi_{0}(1) + 2 \int_{1}^{t} \Phi'(s) ds, & t > 1. \end{cases}$$

Then Φ_0 is a Young function with $\Phi(t) \leqslant \Phi_0(t) \leqslant 2\Phi(t)$. Consequently, $2\Phi^{\sim}(t/2) \leqslant \Phi_0^{\sim}(t) \leqslant \Phi^{\sim}(t)$. We then take

$$g(t) = \sup\{s: \ \Phi_0(s) \leqslant st\}.$$

3. The weight class B_{Φ} . Throughout this section, we assume that w is a nonnegative measurable function on E^n and Φ is a Young function.

We call w a nontrivial weight if there is at least one cube Q with $0 < wQ < \infty$. We then define B_{Φ} to consist of the nontrivial weights w for which there is a constant c with

(3.1)
$$w\{x \in E^n: Mf(x) > \lambda\} \leq \int \Phi(c|f(x)|/\lambda)w(x) dx$$

for all $\lambda > 0$ and all $f \in L^{\Phi}(wdx)$. If c_0 is any value of c for which (3.1) is valid, we say $w \in B_{\Phi}$ with constant c_0 .

Note that if f is the characteristic function of an arbitrary cube Q, then $Mf \geqslant 2^{-n}$ on Q^* . Consequently, whenever $w \in B_{\Phi}$ we must have $wQ^* \leqslant \Phi(2^nc)wQ$. Thus w must be a doubling measure which is finite and positive on all cubes.

(3.2) LEMMA. If $w \in B_{\phi}$ with constant c_0 , then for every $\varepsilon > 0$ we have $(w+\varepsilon) \in B_{\phi}$ with constant $2 \cdot 3^n c_0$.

Proof. First note that $\Phi(c_0) \ge 1$ since the maximal function of the characteristic function of any cube Q is 1 on Q.

Next note that

$$|\{x\colon Mf(x)>\lambda\}| \leq 2\cdot 3^n \int\limits_{(x\colon |f(x)|>\lambda/2)} |f/\lambda| \, dx.$$

This is quite standard; it can be proved by using (2.1) and the observation

$$\{x\colon Mf(x)>\lambda\}\subset \{x\colon Mg(x)>\lambda/2\},\$$

where g is the restriction of f to $\{x: |f(x)| > \lambda/2\}$. On this set, we have

$$2 \cdot 3^n |f(x)|/\lambda \le (2 \cdot 3^n |f(x)|/\lambda) \Phi(c_0) \le \Phi(2 \cdot 3^n c_0 |f(x)|/\lambda).$$

Consequently,

$$(w+\varepsilon)\{x: Mf(x) > \lambda\} \leq \int \Phi(c_0|f|/\lambda)w \, dx + \varepsilon \int \Phi(2\cdot 3^n c_0|f(x)|/\lambda) \, dx$$
$$\leq \int \Phi(2\cdot 3^n c_0|f|/\lambda)(w+\varepsilon) \, dx.$$

Now we are ready to characterize the weight class B_{Φ} by means of a condition analogous to the A_p condition in the form used by Chung, Hunt, and Kurtz [2].

(3.3) THEOREM. For w a nontrivial weight, $w \in B_{\Phi}$ if and only if w defines a doubling measure and there is a constant c such-that

$$||1/w||_{\Phi^-, w, Q} \leq c|Q|/wQ$$

for all cubes Q.

Proof. If Q is any cube such that $\int_{Q} |f(x)| dx > \lambda |Q|$, then

$$\frac{|Q|}{wQ} < \int_{Q} \frac{1}{w(x)} \cdot \frac{|f(x)|}{\lambda} \cdot \frac{w(x) dx}{wQ} \le 2 \|1/w\|_{\Phi^-, w, Q} \|f/\lambda\|_{\Phi, w, Q}.$$

If also $||1/w||_{\Phi^*, w, Q} \le c_1 |Q|/wQ$, then $||f/\lambda||_{\Phi, w, Q} > 1/2c_1$ so that

$$\int_{Q} \Phi(2c_1 |f|/\lambda) w \, dx > wQ.$$

Consequently, for any compact $K \subset \{x: Mf(x) > \lambda\}$, we can apply (2.1) to choose a sequence of cubes Q_k with pairwise disjoint interiors such that both $wK \leq \sum wQ_k^*$ and

$$wQ_k \leqslant \int_{Q_k} \Phi(2c_1|\lambda)w dx.$$

If also w defines a doubling measure, then

$$wK \leqslant \sum c_2 wQ_k \leqslant \sum c_2 \int\limits_{Q_k} \Phi(2c_1 |f|/\lambda) w \, dx \leqslant \int \Phi(2c_1 c_2 |f|/\lambda) w \, dx.$$

Taking the supremum over all compact subsets of $\{x: Mf(x) > \lambda\}$ then gives $w \in B_{\Phi}$ with constant $2c_1c_2$.

Now suppose that $w \in B_{\phi}$ with constant c_0 . With the additional assumption that 1/w is bounded, we prove that

$$\int_{Q} \Phi^{\sim}(1/\lambda w) w \, dx \leqslant wQ \quad \text{for } \lambda = cc_0 |Q|/wQ.$$

We then use (3.2) to remove the extra assumption.

For g the function given by (2.2), continuity and monotonicity allow us to choose s > 0 (depending on Q) such that

$$\int_{Q} g(1/sw) \, dx = swQ.$$

Since $tg(t) \ge 2\Phi^{\sim}(t/2)$, we then have

$$2\int_{Q} \Phi^{*}(1/2sw)sw \, dx \leqslant \int_{Q} g(1/sw) \, dx \leqslant swQ, \quad \text{or}$$
$$\int_{Q} \Phi^{*}(1/2sw)w \, dx \leqslant wQ/2.$$

For f(x) = g(1/sw(x)) on Q and 0 elsewhere, we have $Mf(x) \ge swQ/|Q|$ on Q. Thus for $w \in B_{\Phi}$ with constant c_0 ,

$$wQ \leq \int_{Q} \Phi(c_0|Q|g(1/sw)/swQ)w dx.$$

But $\int_{Q} \Phi(g(1/sw))w \, dx \le (1/s) \int_{Q} g(1/sw) \, dx = wQ$, so we must have $1 \le c_0 |Q|/swQ$. Thus for $\lambda = 2c_0 |Q|/wQ$ we have

$$\int_{Q} \Phi^{\sim}(1/\lambda w) w \, dx \leqslant \int_{Q} \Phi^{\sim}(1/2sw) w \, dx \leqslant wQ.$$

For general $w \in B_{\Phi}$ with constant c_0 , (3.2) guarantees that $(w + \varepsilon) \in B_{\Phi}$ with constant $2 \cdot 3^n c_0$ for each $\varepsilon > 0$, and of course $1/(w + \varepsilon)$ is bounded. Consequently, for

$$\lambda = 2^2 \cdot 3^n c_0 |Q|/wQ > 2^2 \cdot 3^n c_0 |Q|/(wQ + \varepsilon |Q|)$$

we have

$$\int_{Q} \Phi^{\sim} (1/\lambda (w+\varepsilon)) (w+\varepsilon) dx \leq wQ + \varepsilon |Q|.$$

Letting $\varepsilon \to 0$ gives $\int_{Q} \Phi^{\sim}(1/\lambda w) w \, dx \leq wQ$ by monotone convergence.

Now we consider inclusion relations among B_{Φ} classes. As we noted in the proof of (3.2), $w\{x \colon Mf(x) > \lambda\}$ is controlled by the restriction of f to $\{x \colon |f(x)| > \lambda/2\}$. Consequently, $w \in B_{\Phi}$ means Φ has sufficiently rapid growth at infinity; we always have $B_{\Phi} \subset B_{\Psi}$ if $\Phi(t) \leq \Psi(ct)$ for all sufficiently large t. On the other hand, Muckenhoupt [6] proved that for $w \in A_p$, we must have $w \in A_q$ for some q < p, which reverses the natural inclusion. In the remainder of this section, we develop a generalization of Muckenhoupt's result by adapting the proof given by Chung, Hunt, and Kurtz [2] to a restricted class of Orlicz spaces.

(3.4) LEMMA. If $w \in B_{\phi}$ with constant c_0 , then for each cube Q we have $wQ \leq \Phi(2c_0)w\{x \in O: w(x) < 2wQ/|Q|\}$.

Proof. Let $E = \{x \in Q: w(x) < 2wQ/|Q|\}$ and F = Q - E. Then $2|F|wQ/|Q| \le \int_F w \, dx \le wQ$, so that $|F| \le |Q|/2$ and $|E| \ge |Q|/2$. Take $f = \chi_E$; then $Mf \ge 1/2$ on Q. Hence

$$wQ \le w\{x: Mf(x) \ge 1/2\} \le \int \Phi(2c_0|f|)w \, dx = \Phi(2c_0)wE.$$

(3.5) Lemma. Suppose $\Phi^{\sim}(st) \leq A\Phi^{\sim}(s)\Phi^{\sim}(t)$ for all s, t > 0, and $w \in B_{\Phi}$ with constant c_0 . Set $E(\lambda, t, Q) = \{x \in Q: \Phi^{\sim}(1/\lambda w) > t\}$. Then there are constants c_1 and ε (depending only on c_0 and A) such that

$$\int_{E(\lambda,t,Q)} \Phi^{\sim}(1/\lambda w) w \, dx \leqslant c_1 t w E(\lambda, \, \varepsilon t, \, Q)$$

whenever

$$\int_{Q} \Phi^{\sim}(1/\lambda w) w \, dx \leqslant t w Q.$$

Proof. Since w is a doubling measure, the Calderón-Zygmund decomposition shows that we can select nonoverlapping subcubes Q_i of Q such that

$$twQ_i < \int_{Q_i} \Phi^{\sim}(1/\lambda w) w \, dx \leq \Phi(2^n c_0) twQ_i$$

for each i and $\Phi^{\sim}(1/\lambda w) \leq t$ a.e. in $Q - \bigcup Q_i$. Hence

$$\int\limits_{E(t,\lambda,Q)}\Phi^{\sim}(1/\lambda w)w\,dx\leqslant \sum \Phi(2^nc_0)twQ_i$$

and

$$wQ_i \leq \Phi(2c_0)w\{x \in Q_i: w < 2wQ_i/|Q_i|\}$$

by (3.4). We complete the proof by finding $\varepsilon > 0$ such that

$$\{x \in Q_i: w < 2wQ_i/|Q_i|\} \subset E(\lambda, \varepsilon t, Q_i),$$

or equivalently, $\Phi^{\sim}(|Q_i|/2\lambda wQ_i) \ge \varepsilon t$.

By (3.3), we can choose a constant c with

$$\int_{Q_i} \Phi^{\sim}(wQ_i/c|Q_i|w)w\,dx \leqslant wQ_i.$$

Since

$$twQ_{i} \leqslant \int_{Q_{i}} \Phi^{\sim}(1/\lambda w)w \, dx \leqslant A \int_{Q_{i}} \Phi^{\sim}(c|Q_{i}|/\lambda wQ_{i}) \Phi^{\sim}(wQ_{i}/c|Q_{i}|w)w \, dx$$
$$\leqslant A\Phi^{\sim}(c|Q_{i}|/\lambda wQ_{i})wQ_{i} \leqslant A^{2}\Phi^{\sim}(2c)\Phi^{\sim}(|Q_{i}|/2\lambda wQ_{i})wQ_{i},$$

we may take $1/\varepsilon = A^2 \Phi^{\sim}(2c)$.

(3.6) THEOREM. Suppose $w \in B_{\Phi}$ with constant c_0 . If $\Phi^{\sim}(st) \leq A\Phi^{\sim}(s)\Phi^{\sim}(t)$ for all s, t > 0, then there is a $\delta > 0$ (depending only on c_0 and A), such that $w \in B_{\Psi}$ for $\Psi^{\sim}(t) = \Phi^{\sim}(t)^{1+\delta}$.

Proof. By (3.3) it suffices to find a constant c for which

$$\int_{Q} \Psi^{\sim}(1/\lambda w) w \, dx \leqslant wQ \quad \text{when } \lambda = c|Q|/wQ;$$

as in the proof of (3.3) it suffices to prove this under the additional assumption that 1/w is bounded.

Let us define a measure μ by $d\mu = \Phi^{\sim}(1/\lambda w)w dx$; then

$$\int_{O} \Psi^{\sim}(1/\lambda w) w \, dx = \int_{O} \Phi^{\sim}(1/\lambda w)^{\delta} \, d\mu = \delta \int_{O}^{\infty} t^{\delta-1} \, \mu E(\lambda, t, Q) \, dt$$

where $E(\lambda, t, Q)$ is the set defined in (3.5). For $t \le 1$ we may bound

$$\mu E(\lambda, t, Q) \leqslant \mu Q = \int_{Q} \Phi^{*}(1/\lambda w) w \, dx,$$

while for t > 1 and $\lambda \ge \|1/w\|_{\Phi^*, w, Q}$ we have $\int_Q \Phi^*(1/\lambda w) w \, dx < twQ$, and hence $\mu E(\lambda, t, Q) \le c_1 tw(\lambda, \varepsilon t, Q)$ by (3.5). Thus

$$\begin{split} \int_{Q} \Psi^{\sim}(1/\lambda w) w \, dx & \leqslant \mu Q + \delta \int_{1}^{\infty} t^{\delta-1} \mu E(\lambda, \, t, \, Q) \, dt \leqslant \mu Q + c_{1} \delta \int_{1}^{\infty} t^{\delta} w E(\lambda, \, \varepsilon t, \, Q) \, dt \\ & \leqslant \mu Q + c_{1} \delta \varepsilon^{-1-\delta} \int_{0}^{\infty} t^{\delta} w E(\lambda, \, t, \, Q) \, dt \\ & = \mu Q + c_{1} \delta \varepsilon^{-1-\delta} (1+\delta)^{-1} \int_{Q} \Phi^{\sim}(1/\lambda w)^{1+\delta} w \, dx \, . \end{split}$$

Choosing δ small enough that $c_1 \delta \varepsilon^{-1-\delta} (1+\delta)^{-1} \le 1/2$ gives

$$\int\limits_{Q}\Psi^{\sim}(1/\lambda w)w\,dx \leqslant 2\mu Q = 2\int\limits_{Q}\Phi^{\sim}(1/\lambda w)w\,dx \leqslant wQ \quad \text{ for } \lambda = 2\,\|1/w\|_{\Phi^{\sim},w,Q}.$$

The condition $\Phi^{\sim}(st) \leq A\Phi^{\sim}(s)\Phi^{\sim}(t)$ excludes quite a few of the standard Young functions. If $\Phi(t) \approx t^p(\log t)^{\alpha p}$ for large t, then $\Phi^{\sim}(t) \approx t^q(\log t)^{-\alpha q}$ for large t, where p and q are Hölder conjugates. Consequently, (3.6) applies when $\alpha \leq 0$

but not for $\alpha > 0$. In the next section we show that the conclusion of (3.6) fails when $\Phi(t) = t^2 (1 + \log_+ t)^2$, so that the condition on Φ^- cannot be eliminated from the hypotheses of (3.6).

In calculations using a given weight, the best possible weighted estimates for Mf would be obtained by finding an optimal Young function Φ for which $w \in B_{\Phi}$. In such cases, (3.6) says that its conjugate Φ^{\sim} should not satisfy the condition $\Phi^{\sim}(st) \leq A\Phi^{\sim}(s)\Phi^{\sim}(t)$.

- 4. Applications and examples. By using (5.3) of [1], we can obtain a variety of weighted Orlicz space bounds for the Hardy-Littlewood maximal function when the weight is in B_{Φ} . We summarize them in (4.1) below.
- (4.1) THEOREM. For $w \in B_{\Phi}$, the Hardy-Littlewood maximal function is bounded from $L^{\Theta}(wdx)$ to $L^{\Psi}(wdx)$ provided

$$\int_{0}^{t} \Psi'(s) \Phi(t/s) ds \leqslant \Theta(ct)$$

for some constant c.

While there are cases (including $\Phi(t) = t^p$, 1) where (3.6) and (4.1) can be used to prove that <math>Mf is bounded in $L^{\Phi}(wdx)$ for all $w \in B_{\Phi}$, this conclusion is generally false. We give two simple examples where B_{Φ} contains weights not in A_{Φ} .

(4.2) Example. Take

$$\Phi(t) = \begin{cases} t^2 & \text{for } t \leq 1, \\ t^3 & \text{for } t > 1. \end{cases}$$

Then $L^{\phi} = L^2 \cap L^3$, so that $L^{\phi^2} = L^2 + L^3$. It follows easily that $B_{\phi} = A_3$ while $A_{\phi} = A_2$.

(4.3) Example. Let

$$\Phi^{\sim}(t) = \begin{cases} \left(\frac{t}{2+t}\right)^2, & 0 \le t \le 1, \\ \left(\frac{t}{3+\log t}\right)^2, & t > 1. \end{cases}$$

One may verify that $\Phi(t) \approx (t \log t)^2$ for large t, so that $w \in B_{\Phi}$ means

$$\int_{\{Mf(x)>\lambda\}} w(x) dx \leqslant c^2 \int \left(\frac{|f(x)|}{\lambda}\right)^2 \left(1 + \log_+ \frac{c|f(x)|}{\lambda}\right)^2 w(x) dx.$$

Let us take n = 1 and w(x) = |x|; then to prove $w \in B_{\Phi}$ it is enough to check that

$$\int_{0}^{a} \Phi^{\sim}(1/\lambda x) x \, dx \leqslant a^{2}/2 \quad \text{for } \lambda = c/a, \ a > 0.$$

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Choosing c = 1 gives

$$\int_{0}^{a} \Phi^{\sim}(a/x)x \, dx = \int_{0}^{a} \frac{a^{2} \, dx}{x(3 + \log a/x)^{2}} = a^{2}/3.$$

On the other hand, for any δ , $\varepsilon > 0$ we have

$$\int_{0}^{a} \Phi^{\sim}(\varepsilon/x)^{1+\delta} x \, dx = \infty.$$

The characterization of A_{ϕ} given by Kerman and Torchinsky [4] shows that $A_{\phi} = A_2$; in this case $w \in A_p$ for all p > 2 but w is not in A_2 .

Let us now apply (4.1) with this choice of Φ . Take $\Psi(t) = t^{2} - 1$ for t > 1 and 0 otherwise. Then L^{Ψ} consists of the functions whose restriction to every set of finite measure is in L^{2} . Since

$$\int_{1}^{t} (2s)(t/s)^{2} (1 + \log t/s)^{2} ds = (2/3)t^{2} (1 + \log t)^{3},$$

we see that the Hardy-Littlewood maximal function is bounded from $L^2(\log L)^3(wdx)$ to $L^2_{loc}(wdx)$ for all $w \in B_{\phi}$.

For the case w(x) = |x| on E^1 , slightly better weighted bounds for Mf can be obtained by using the fact that w is in the weight class A(2,1) of Chung, Hunt, and Kurtz [2]. However, if we modify w by redefining $w(x) = |x|/(\log 2/x)^{\epsilon}$ for |x| < 1, then $w \notin A(2,1)$ for $\varepsilon > 0$ but $w \in B_{\Phi}$ for $\varepsilon > 1$.

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H^p spaces over open subsets of R^n

by

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Abstract. Part of the theory of H^p spaces over R^n , originated by C. Fesserman and E. M. Stein [4], is generalized to the case of arbitrary open subsets of R^n . The following subjects are treated: (1) Desinition of $H^p(\Omega)$, where Ω is an open subset of R^n , by means of maximal functions; (2) Atomic decomposition for $H^p(\Omega)$; (3) Identification of the duals of $H^p(\Omega)$ with certain function spaces over Ω ; (4) The complex method of interpolation for $H^p(\Omega)$ and $L^p(\Omega)$; (5) Extension of a distribution in $H^p(\Omega)$ to a distribution in $H^p(R^n)$. All the results are given in the situation that R^n has a parabolic metric.

1. Introduction. In this paper, we introduce H^p spaces over arbitrary open subsets of R^n by means of certain maximal functions and show that they have some properties similar to the H^p spaces over R^n (for the H^p spaces over R^n , see Calderón-Torchinsky [1], [2] or Torchinsky's book [10; Chapt. XIV]).

We briefly review our results.

Let φ be a function in $C_0^{\infty}(\mathbf{R}^n)$ such that $\operatorname{supp} \varphi \subset \{x \in \mathbf{R}^n \mid |x| < 1\}$ (if $x \in \mathbf{R}^n$, then |x| denotes the usual Euclidean norm of x) and $\int \varphi(x) dx = 1$. For t > 0, we define $(\varphi)_t$ by $(\varphi)_t(x) = t^{-n} \varphi(t^{-1}x)$ (we shall modify this definition afterwards; see the next to the last paragraph in this section). Let Ω be an open subset of \mathbf{R}^n . For $f \in \mathcal{D}'(\Omega)$, we define the radial maximal function $f_{\varphi,\Omega}^+(x)$, $x \in \Omega$, by

$$f_{\varphi,\Omega}^+(x) = \sup\{|\langle f,(\varphi)_t(x-\cdot)\rangle| | 0 < t < \operatorname{dis}(x,\Omega^c)\},\,$$

where Ω° denotes the complement of Ω (throughout this paper, $\mathscr{D}'(\Omega)$ denotes the set of distributions on Ω and $\langle f, \psi \rangle$, where $f \in \mathscr{D}'(\Omega)$ and $\psi \in C_0^{\infty}(\Omega)$, means $f(\psi)$; we use the same notation $\langle f, \psi \rangle$ if f is a distribution with compact support and ψ is a smooth function on R^n). For p with $0 , we define <math>H^p(\Omega)$ as the set of those $f \in \mathscr{D}'(\Omega)$ for which $f_{\psi,\Omega}^+$ belongs to $L^p(\Omega)$. We consider $H^p(\Omega)$ a quasinormed linear space by defining the quasinorm of $f \in H^p(\Omega)$ to be equal to the $L^p(\Omega)$ -norm of $f_{\psi,\Omega}^+$. (By a quasinorm we mean a function σ on a linear space X which has the following properties: (i) $\sigma(x) > 0$ if $x \neq 0$ and $\sigma(0) = 0$; (ii) $\sigma(\lambda x) = |\lambda| \sigma(x)$ for all scalars λ and all $x \in X$; (iii) there exists a positive constant k such that $\sigma(x+y) \le k(\sigma(x)+\sigma(y))$ for all $x, y \in X$.) Then the maximal inequality given by the author [8] shows that the above definition

BMO, complex method of interpolation, parabolic metric.

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