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## LOCAL EXISTENCE OF SOLUTIONS OF THE FREE BOUNDARY PROBLEM FOR THE EQUATIONS OF A MAGNETOHYDRODYNAMIC INCOMPRESSIBLE FLUID

*Abstract.* Local existence of solutions is proved for equations describing the motion of a magnetohydrodynamic incompressible fluid in a domain bounded by a free surface. In the exterior domain we have an electromagnetic field which is generated by some currents located on a fixed boundary. First by the Galerkin method and regularization techniques the existence of solutions of the linearized equations is proved; next by the method of successive approximations the local existence is shown for the nonlinear problem.

**1. Introduction.** In this paper we prove the existence of local solutions to equations describing the motion of a magnetohydrodynamic incompressible fluid in a domain  $\Omega_t \subset \mathbb{R}^3$  bounded by a free surface  $S_t$ . In the domain  $D_t \subset \mathbb{R}^3$  which is exterior to  $\Omega_t$  we have a gas under constant pressure  $p_0$ . Moreover in the domain  $D_t$  we have an electromagnetic field which is generated by some currents which are located on a fixed boundary  $B$  of  $D_t$ .

In the domain  $\Omega_t$  the motion is described by the following problem,

$$\begin{aligned}
 (1.1) \quad & v_t + v \cdot \nabla v - \operatorname{div} \mathbb{T}(v, p) - \mu_1 \overset{1}{H} \cdot \nabla \overset{1}{H} + \mu_1 \nabla \overset{1}{H}^2 = f && \text{in } \tilde{\Omega}^T, \\
 & \operatorname{div} v = 0 && \text{in } \tilde{\Omega}^T, \\
 & \mu_1 \overset{1}{H}_t = -\operatorname{rot} \overset{1}{E} && \text{in } \tilde{\Omega}^T, \\
 & \operatorname{rot} \overset{1}{H} = \sigma_1 (\overset{1}{E} + \mu_1 v \times \overset{1}{H}), && \text{in } \tilde{\Omega}^T, \\
 & \operatorname{div}(\mu_1 \overset{1}{H}) = 0, && \text{in } \tilde{\Omega}^T,
 \end{aligned}$$

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where  $\tilde{\Omega}^T = \bigcup_{0 \leq t \leq T} \Omega_t \times \{t\}$ ,  $v = v(x, t)$  is the velocity of the fluid,  $p = p(x, t)$  is the pressure,  $\overset{1}{H} = \overset{1}{H}(x, t)$  is the magnetic field,  $f = f(x, t)$  is the external force field per unit mass,  $\mu_1$  is the constant magnetic permeability,  $\sigma_1$  is the constant electric conductivity,  $\overset{1}{E} = \overset{1}{E}(x, t)$  is the electric field, and

$$(1.2) \quad \mathbb{T}(v, p) = \{\nu(\partial_{x_i} v_j + \partial_{x_j} v_i) - p\delta_{ij}\}$$

is the stress tensor, where  $\nu$  is the viscosity of the fluid. Moreover, by

$$(1.3) \quad \mathbb{D}(v) = \{\nu(\partial_{x_i} v_j + \partial_{x_j} v_i)\}$$

we denote the dilatation tensor.

In the domain  $D_t$  which is a dielectric (gas) we assume that there is no fluid motion inside ( $v = 0$ ). Therefore we have the electromagnetic field only described by the following system:

$$(1.4) \quad \begin{aligned} \mu_2 \overset{2}{H}_t &= -\text{rot } \overset{2}{E} && \text{in } \tilde{D}^T, \\ \sigma_2 \overset{2}{E} &= \text{rot } \overset{2}{H} && \text{in } \tilde{D}^T, \\ \text{div}(\mu_2 \overset{2}{H}) &= 0 && \text{in } \tilde{D}^T, \end{aligned}$$

where  $\tilde{D}^T = \bigcup_{0 \leq t \leq T} D_t \times \{t\}$ .

On  $S_t = \partial\Omega_t \cap \partial D_t$  we assume the following transmission and boundary conditions:

$$(1.5) \quad \begin{aligned} n \cdot \mathbb{T}(v, p) &= -p_0 n && \text{on } \tilde{S}^T, \\ \frac{1}{\sigma_1} \overset{1}{H} &= \frac{1}{\sigma_2} \overset{2}{H} && \text{on } \tilde{S}^T, \\ \overset{1}{E} \cdot \tau_\alpha &= \overset{2}{E} \cdot \tau_\alpha, \quad \alpha = 1, 2, && \text{on } \tilde{S}^T \\ v \cdot n &= -\frac{\phi_t}{|\nabla\phi_t|} && \text{on } \tilde{S}^T, \end{aligned}$$

where  $\tilde{S}^T = \bigcup_{0 \leq t \leq T} S_t \times \{t\}$ ,  $n$  is the unit outward vector to  $\Omega_t$  and normal to  $S_t$ ,  $\tau_\alpha$ ,  $\alpha = \bar{1}, \bar{2}$ , is the tangent vector to  $S_t$ ,  $\phi(x, t) = 0$  describes  $S_t$  at least locally.

Next we assume the following boundary conditions on  $B$ :

$$(1.6) \quad \begin{aligned} \overset{2}{H} &= H_* && \text{on } B, \\ \overset{2}{E} &= E_* && \text{on } B, \end{aligned}$$

where  $H_*$  and  $E_*$  are connected by

$$\begin{aligned} \sigma_1 E_{*n} &= \frac{1}{A_{\tau_1} A_{\tau_2}} (\partial_{\tau_1} (H_{*\tau_1} A_{\tau_2}) - \partial_{\tau_2} (H_{*\tau_1} A_{\tau_1})), \\ \mu_2 \partial_t H_{*n} &= \frac{1}{A_{\tau_1} A_{\tau_2}} (\partial_{\tau_2} (E_{*\tau_1} A_{\tau_1}) - \partial_{\tau_1} (E_{*\tau_2} A_{\tau_2})), \\ -\partial_t \partial_{\tau_1} (H_{*\tau_1} A_{\tau_2} A_n) - \partial_t \partial_{\tau_2} (H_{*\tau_2} A_{\tau_1} A_n) &= \partial_{\tau_1} \partial_{\tau_2} (E_{*n} A_n) \\ -\mu_2 \partial_{\tau_1} (A_{\tau_2} A_{\tau_3} \partial_t H_{*\tau_1}) - \mu_2 \partial_{\tau_2} (A_{\tau_1} A_{\tau_3} \partial_t H_{*\tau_2}) - \partial_{\tau_2} \partial_{\tau_1} (E_{*n} A_n), \end{aligned}$$

where  $(\tau_1, \tau_2, n)$  are curvilinear coordinates and  $A_{\tau_1}, A_{\tau_2}, A_n$  the Lamé coefficients of the transformation  $(\tau_1, \tau_2, n) \mapsto (x_1, x_2, x_3)$ .

Finally, we assume the initial conditions

$$\begin{aligned} \Omega_t|_{t=0} &= \Omega, & S_t|_{t=0} &= S, & D_t|_{t=0} &= D, \\ (1.7) \quad v|_{t=0} &= v_0, & \overset{1}{H}|_{t=0} &= \overset{1}{H}_0 & \text{in } \Omega, \\ \overset{2}{H}|_{t=0} &= \overset{2}{H}_0 & & & \text{in } D. \end{aligned}$$

To prove the existence of solutions to the above problem we introduce the Lagrangian coordinates  $\xi \in \Omega$ . The Lagrangian coordinates are connected with the velocity  $v$  as the initial data for the Cauchy problem

$$(1.8) \quad \frac{dx}{dt} = v(x, t), \quad x|_{t=0} = \xi \in \Omega.$$

Therefore  $x_v(\xi, t) = \xi + \int_0^t \bar{v}(\xi, \tau) d\tau$ , where

$$\bar{v}(\xi, t) = v(x_v(\xi, t), t).$$

To introduce the Lagrangian coordinates in  $D_t$  we extend  $v$  onto  $D_t$ . Let us denote the extended functions by  $v'$ . Then we define  $\xi \in D$  to be the Cauchy data for the problem

$$(1.9) \quad \frac{dx}{dt} = v'(x, t), \quad x|_{t=0} = \xi \in D.$$

Therefore  $x_{v'}(\xi, t) = \xi + \int_0^t \bar{v}'(\xi, \tau) d\tau$ , where  $\bar{v}'(\xi, t) = v'(x_{v'}(\xi, t), t)$ . Then by (1.1)<sub>5</sub>,

$$\begin{aligned} \Omega_t &= \{x \in \mathbb{R}^3 : x = x_v(\xi, t), \xi \in \Omega\}, \\ S_t &= \{x \in \mathbb{R}^3 : x = x_v(\xi, t), \xi \in S\}. \end{aligned}$$

Since  $S_t$  is determined at least locally by the equation  $\phi(x, t) = 0$ ,  $S$  is described by  $\phi(x_v(\xi, t), t)|_{t=0} = 0$ . Moreover we have

$$\bar{n}_v = n(x_v(\xi, t), t) = \frac{\nabla_x \phi(x, t)}{|\nabla_x \phi(x, t)|} \Big|_{x=x_v(\xi, t)}.$$

We introduce the following notation:

$$\|u\|_{l,Q} = \|u\|_{H^l(Q)}, \quad Q \in \{\Omega, S, D, \Pi, B\}, \quad 0 \leq l \in \mathbb{Z},$$

$$\|u\|_{k,p,q,Q^T} = \|u\|_{L_q(0,T,W_p^k(Q))}, \quad Q \in \{\Omega, S, D, \Pi, B\},$$

$$p, q \in [1, \infty], \quad 0 \leq k \in \mathbb{Z},$$

where  $Q^t = Q \times (0, t)$ , and

$$\|u\|_{p,Q} = \|u\|_{L_p(Q)}, \quad Q \in \{\Omega, S, D, \Pi, B\}, \quad p \in [1, \infty].$$

**2. Weak solutions.** Weak solutions to problem (1.1)–(1.7) are defined in Lagrangian coordinates. From the equality

$$\int_{\Omega_t} \operatorname{rot} H\psi \, dx = \int_{\Omega_t} H \operatorname{rot} \psi \, dx - \int_{S_t} (n \times H)\psi \, dx_{S_t}$$

and from (1.1)–(1.5) we get

$$\begin{aligned} & \frac{1}{\sigma_2} \int_{S_t} (n \times \operatorname{rot} \overset{2}{H})\psi \, dx_{S_t} - \frac{1}{\sigma_1} \int_{S_t} (n \times \operatorname{rot} \overset{1}{H})\psi \, dx_{S_t} \\ & + \mu_1 \int_{S_t} n \times (v \times \overset{1}{H})\psi \, dx_{S_t} = \int_{S_t} (n \times \overset{2}{E})\psi \, dx_{S_t} \\ & - \int_{S_t} n \times (\overset{1}{E} + \mu_1 v \times \overset{1}{H})\psi \, dx_{S_t} + \mu_1 \int_{S_t} n \times (v \times \overset{1}{H})\psi \, dx_{S_t} = 0. \end{aligned}$$

DEFINITION 2.1. By *weak solutions* to problem (1.1)–(1.7) we mean functions  $\bar{v}, \bar{H}$  which satisfy the integral identities

$$\begin{aligned} (2.1) \quad & \int_0^T \int_{\Omega} (-\bar{v}\bar{\varphi}_t + \mathbb{D}_v(\bar{v})\mathbb{D}_v(\bar{\varphi})) \, d\xi \, dt - \int_0^T \int_{\Omega} (\mu_1 \bar{H}\nabla_v \bar{H} \cdot \bar{\varphi} - \mu_1 \nabla_v \bar{H}^2 \bar{\varphi}) \, d\xi \, dt \\ & = \int_0^T \int_{\Omega} \bar{f}\bar{\varphi} \, d\xi \, dt - \int_0^T \int_S p_0 \bar{n}_v \varphi \, d\xi_S \, dt - \int_{\Omega} \bar{v}_0 \bar{\varphi}(0) \, d\xi, \end{aligned}$$

$$\begin{aligned} (2.2) \quad & \int_0^T \int_{\Pi} \left( -\mu \bar{H}\bar{\psi}_t - \mu \bar{v}\nabla_v \bar{H}\bar{\psi} + \frac{1}{\sigma} \operatorname{rot}_v \bar{H} \operatorname{rot}_v \bar{\psi} \right) d\xi \, dt \\ & - \int_0^T \int_{\Omega} \mu_1 (\bar{v} \times \overset{1}{H}) \operatorname{rot}_v \bar{\psi} \, d\xi \, dt = \frac{1}{\sigma_2} \int_0^T \int_B (\bar{n}_v \times \bar{E}_*) \bar{\psi} \, d\xi_B \, dt \\ & - \mu \int_{\Pi} \bar{H}_0 \bar{\psi}(0) \, d\xi, \end{aligned}$$

where  $\varphi, \psi$  are sufficiently regular with  $\varphi(x, T) = \psi(x, T) = 0$ , and  $\bar{n}_v$  is the unit outward vector normal to  $S$  or  $B$ .

In (2.1), (2.2) we use the notation  $\bar{A}(\xi, t) = A(x_v(\xi, t), t)$ ,  $\bar{H}|_\Omega = \bar{H}^{\frac{1}{2}}$ ,  $\bar{H}|_D = \bar{H}^{\frac{2}{2}}$ ,  $\sigma|_\Omega = \sigma_1$ ,  $\sigma|_D = \sigma_2$ ,  $\Pi = \Omega \cup D$ ,  $\mu|_\Omega = \mu_1$ ,  $\mu|_D = \mu_2$ ; in (2.2)  $v$  is extended onto  $\Pi$ ,

$$\begin{aligned} \mathbb{D}_v(\bar{v}) &= \{\nu(\partial_{x_i}\xi_k \nabla_{\xi_k} \bar{v}_j + \partial_{x_j}\xi_k \nabla_{\xi_k} \bar{v}_i)\}, \quad \text{rot}_v \bar{v} = \nabla_v \times \bar{v}, \\ \nabla_v &= \partial_x \xi_i \nabla_{\xi_i}, \quad \text{div}_v \bar{v} = \nabla_v \cdot \bar{v} = \partial_{x_i}\xi_k \nabla_{\xi_k} \bar{v}_i, \quad \partial_{\xi_i} = \nabla_{\xi_i}. \end{aligned}$$

Let  $A$  be the Jacobi matrix of the transformation  $x = x_v(\xi, t)$ . Then  $\det A = \exp(\int_0^t \text{div}_v \bar{v} d\tau) = 1$ .

Moreover  $x_{\xi_j}^i = \delta_{ij} + \int_0^t \partial_{\xi_j} \bar{v}^i(\xi, \tau) d\tau$  and  $\xi_x = x_\xi^{-1}$ . Then we get

$$\begin{aligned} \sup_{\xi \in \Omega} |x_\xi| &\leq 1 + \sup_{\xi \in \Omega} \int_0^t |\bar{v}(\xi, \tau)| d\tau \leq 1 + c \int_0^t \|\bar{v}\|_{3,\Omega} d\tau \\ &\leq 1 + c\sqrt{t} \left( \int_0^t \|\bar{v}\|_{3,\Omega}^2 d\tau \right)^{1/2} \leq 1 + c\sqrt{t} \|\bar{v}\|_{3,2,2,\Omega^t}. \end{aligned}$$

Thus  $\sup_{x \in \Omega_t} |\xi_x| \leq \varphi(a)$ , where  $a = \sqrt{t} \|\bar{v}\|_{3,2,2,\Omega^t}$  and  $\varphi$  is an increasing positive function.

To prove the existence of a solution to the above problem we introduce Lagrangian coordinates connected with a given divergence-free function  $u$ . Moreover we linearize the terms with  $v$  in (1.1) writing them in the form  $\bar{u} \cdot \nabla \bar{v}$  and  $\bar{u} \times \bar{H}$ . Then from (2.1), (2.2) we get

$$\begin{aligned} (2.3) \quad \int_0^T \int_\Omega (-\bar{v} \bar{\varphi}_t + \mathbb{D}_u(\bar{v}) \mathbb{D}_u(\bar{\varphi})) d\xi dt - \int_0^T \int_\Omega (\mu_1 \bar{H}' \nabla_u \bar{H}' \bar{\varphi} - \mu_1 \nabla_u \bar{H}'^2) d\xi dt \\ = \int_0^T \int_\Omega \bar{f} \bar{\varphi} d\xi dt - \int_0^T \int_S p_0 \bar{n}_u \varphi d\xi_S dt - \int_\Omega \bar{v}_0 \bar{\varphi}(0) d\xi, \end{aligned}$$

$$\begin{aligned} (2.4) \quad \int_0^T \int_\Pi \left( -\mu \bar{H} \bar{\psi}_t - \mu \bar{u} \nabla_u \bar{H} \bar{\psi} + \frac{1}{\sigma} \text{rot}_u \bar{H} \text{rot}_u \bar{\psi} \right) d\xi dt \\ - \int_0^T \int_\Omega \mu_1 (\bar{u} \times \bar{H}) \text{rot}_u \bar{\psi} d\xi dt \\ = \frac{1}{\sigma_2} \int_0^T \int_B (\bar{n}_u \times \bar{E}_*) \bar{\psi} d\xi_B dt - \mu \int_\Pi \bar{H}_0 \bar{\psi}(0) d\xi, \end{aligned}$$

where  $\bar{H}'$  is a given function.

We have the following main theorem of this paper, which is proved in Section 5.

MAIN THEOREM. Assume that  $\bar{v}_0 \in H^2(\Omega)$ ;  $\bar{v}_t(0), \bar{v}_{tt}(0) \in L_2(\Omega)$ ;  $\bar{f}_t, \bar{f}_{tt} \in L_2(0, T, L_2(\Omega))$ ;  $\bar{f} \in L_2(0, T, H^2(\Omega))$ ;  $\bar{H}(0) \in H^2(\Pi)$ ;  $\bar{H}_t(0) \in H^1(\Pi)$ ;  $\bar{E}_* \in L_\infty(0, T, H^1(B))$ ;  $E_{*t}, \bar{H}_{*tt} \in L_2(0, T, L_2(B))$ ;  $E_{*t}, \bar{H}_{*tt} \in L_2(0, T, L_2(B))$ ;  $\bar{H}_{*t} \in L_2(0, T, H^2(B))$ ;  $\bar{H}_* \in L_2(0, T, H^3(B))$  and  $S, B \in H^{5/2}$ . Then there exists  $T^{**}$  sufficiently small such that for  $T \leq T^{**}$  there exists a solution to problem (1.1)–(1.7) such that

$$\begin{aligned} &\bar{v} \in L_2(0, T, H^3(\Omega)) \cap L_\infty(0, T, H^1(\Omega)); \\ &\bar{v}_t \in L_\infty(0, T, H^1(\Omega)) \cap L_2(0, T, H^2(\Omega)); \\ &\bar{v}_{tt} \in L_\infty(0, T, L_2(\Omega)) \cap L_2(0, T, H^1(\Omega)); \\ &\bar{p} \in L_2(0, T, H^2(\Omega)), \quad \bar{p}_t \in L_2(0, T, H^1(\Omega)); \\ &\bar{H} \in L_2(0, T, H^3(\Pi)) \cap L_\infty(0, T, H^1(\Pi)); \\ &\bar{H}_t \in L_\infty(0, T, H^1(\Pi)) \cap L_2(0, T, H^2(\Pi)); \\ &\bar{H}_{tt} \in L_\infty(0, T, L_2(\Pi)) \cap L_2(0, T, H^1(\Pi)) \end{aligned}$$

and

$$\begin{aligned} &\|\bar{v}_t\|_{1,2,\infty,\Omega^T}^2 + \|\bar{v}_{tt}\|_{1,2,\infty,\Omega^T}^2 + \|\bar{v}\|_{3,2,2,\Omega^T}^2 + \|\bar{v}_t\|_{2,2,2,\Omega^T}^2 \\ &\quad + \|\bar{v}_{tt}\|_{1,2,2,\Omega^T}^2 + \|\bar{p}\|_{2,2,2,\Omega^T}^2 + \|\bar{p}_t\|_{1,2,2,\Omega^T}^2 + \|\bar{H}_t\|_{1,2,\infty,\Pi^T}^2 + \|\bar{H}\|_{1,2,\infty,\Pi^T}^2 \\ &\quad + \|\bar{H}\|_{3,2,2,\Pi^T}^2 + \|\bar{H}_t\|_{2,2,2,\Pi^T}^2 + \|\bar{H}_{tt}\|_{1,2,2,\Pi^T}^2 \leq A, \end{aligned}$$

where  $A$  is a positive constant.

**3. Existence and regularity of solutions of the linearized problem (1.12).** To prove the existence of a solution to the problems (1.12), (1.13) we use the Galerkin method. Take a basis  $\{\bar{\varphi}_k\}$  in  $L_2(\Omega)$  and  $\{\bar{\psi}_k\}$  in  $L_2(\Pi)$ . Then we are looking for an approximate solution in the form

$$(3.1) \quad \bar{v}_n = \sum_{k=1}^n c_{kn}(t)\bar{\varphi}_k(\xi), \quad \bar{H}_n = \sum_{k=1}^n d_{kn}(t)\bar{\psi}_k(\xi),$$

where the functions  $c_{kn}, d_{kn}, k = 1, \dots, n$ , are solutions of the following system of ordinary differential equations:

$$(3.2) \quad \int_{\Omega} (\bar{v}_{nt}\bar{\varphi}_k + \mathbb{D}_u(\bar{v}_n)\mathbb{D}_u(\bar{\varphi}_k)) d\xi - \mu_1 \int_{\Omega} (\bar{H}'\nabla_u\bar{H}'\bar{\varphi}_k - \nabla_u\bar{H}'^2\bar{\varphi}_k) d\xi \\ = \int_{\Omega} \bar{f}\bar{\varphi}_k d\xi - \int_S p_0\bar{n}_u\bar{\varphi}_k d\xi_S,$$

$$(3.3) \quad \int_{\Pi} (\mu\bar{H}_{nt}\bar{\psi}_k - \mu\bar{u}\nabla_u\bar{H}_n\bar{\psi}_k + \frac{1}{\sigma}\text{rot}_u\bar{H}_n\text{rot}_u\bar{\psi}_k) d\xi \\ - \int_{\Omega} \mu_1(u \times \bar{H})\text{rot}_u\bar{\psi}_k d\xi = \frac{1}{\sigma_2} \int_B (\bar{n}_u \times \bar{E}_*)\bar{\psi}_k d\xi_B,$$

for  $k = 1, \dots, n$ . The equations (3.2), (3.3) can be written in the form

$$(3.2)_1 \quad \frac{d}{dt}c_{kn} + a_{ki}(t)c_{in} = f_k(t),$$

$$(3.3)_1 \quad \frac{d}{dt}d_{nk} + b_{ki}(t)d_{in} = g_k(t),$$

where  $k = 1, \dots, n$  and summation over repeated indices is assumed. Then from (3.2), (3.3) we see that

$$\sum_{k,i} \int_0^T |a_{ki}(t)| dt \leq \varphi(a)(\|\bar{H}'\|_{1,2,2,\Omega^T}^2 + 1),$$

$$\sum_{k,i} \int_0^T |b_{ki}(t)| dt \leq \varphi(a)(\|\bar{u}\|_{1,2,2,\Omega^T}^2 + 1),$$

where  $a = t^{1/2}\|\bar{u}\|_{3,2,2,\Omega^t}$  and  $\varphi$  is an increasing positive function. Assuming

$$(3.4) \quad v_0 = \sum_{k=1}^\infty c_k \bar{\varphi}_k, \quad H_0 = \sum_{k=1}^\infty d_k \bar{\psi}_k$$

we have the following initial conditions for solutions to the problem (3.2), (3.3):

$$(3.5) \quad c_{kn}(0) = c_k, \quad d_{kn}(0) = d_k, \quad n \in \mathbb{N}, \quad k = 1, \dots, n.$$

The existence and uniqueness of solution to the problem (3.2), (3.3), (3.5) follows from the theory of ordinary differential equations.

Next we have to assume that

$$(3.6) \quad \sup_{t \in [0, T]} \sup_{\xi \in \Omega} |I - \xi_x| \leq \delta,$$

where  $\delta$  is sufficiently small and  $I$  is the unit matrix.

Now we obtain estimates for solutions of (3.2), (3.3).

LEMMA 3.1. Assume that  $\bar{H}' \in L_\infty(0, T, L_2(\Omega)) \cap L_2(0, T, H^3(\Omega))$ ;  $\bar{H}'_t \in L_2(0, T, L_2(\Omega))$ ;  $\bar{f} \in L_2(0, T, L_2(\Omega))$ ;  $\bar{v}(0) \in L_2(\Omega)$ ;  $\bar{u} \in L_2(0, T, H^3(\Omega))$ . Then

$$(3.7) \quad \|\bar{v}_n\|_{0,2,\infty,\Omega^t}^2 + \|\bar{v}_n\|_{1,2,2,\Omega^t}^2$$

$$\leq \alpha(t, a) [\|\bar{H}'\|_{0,2,\infty,\Omega^t}^2 (\varepsilon \|\bar{H}'\|_{3,2,2,\Omega^t}^2 + c(\varepsilon)t(\|\bar{H}'(0)\|_{0,\Omega} + \|\bar{H}'_t\|_{0,2,2,\Omega^t}^2))$$

$$+ c\|\bar{f}\|_{0,2,2,\Omega^t}^2 + cp_0^2 + \|\bar{v}_n(0)\|_{0,\Omega}^2],$$

where  $\alpha$  is an increasing positive function,  $a = t^{1/2}\|\bar{u}\|_{3,2,2,\Omega^t}$  and  $0 \leq t \leq T$ .

*Proof.* Multiplying (3.2) by  $c_{kn}$  and summing over  $k$  from 1 to  $n$  we get

$$(3.8) \quad \int_{\Omega} \left( \frac{1}{2} \frac{d}{dt} \bar{v}_n^2 + |\mathbb{D}_u(\bar{v}_n)|^2 \right) d\xi \\ = \mu_1 \int_{\Omega} (\bar{H}' \nabla_u \bar{H}' \bar{v}_n - \nabla_u \bar{H}'^2 \bar{v}_n) d\xi + \int_{\Omega} \bar{f} \bar{v}_n d\xi - \int_S p_0 \bar{n}_u \bar{v}_n d\xi_S.$$

Using in (3.8) the Hölder and Young inequalities together with the Korn inequality

$$\|\bar{v}_n\|_{1,\Omega}^2 \leq c(\delta)(\|\mathbb{D}(\bar{v}_n)\|_{0,\Omega}^2 + \|\bar{v}_n\|_{0,\Omega}^2),$$

we get

$$(3.9) \quad \frac{d}{dt} \|\bar{v}_n\|_{0,\Omega}^2 + \|\bar{v}_n\|_{1,\Omega}^2 \\ \leq \varphi(a) \|\bar{H}'\|_{0,\Omega}^2 \|\bar{H}'\|_{2,\Omega}^2 + c\|p_0 \bar{n}_u\|_{0,S}^2 + c\|\bar{f}\|_{0,\Omega}^2 + c\|\bar{v}_n\|_{0,\Omega}^2.$$

Integrating (3.9) with respect to time and using the Gronwall inequality we get (3.7).

We want to obtain more regular solutions of (3.2); therefore we show

LEMMA 3.2. *Let the assumptions of Lemma 3.1 be satisfied and  $\bar{v}(0) \in H^1(\Omega)$ . Then*

$$(3.10) \quad \|\bar{v}_{nt}\|_{0,2,2,\Omega^t}^2 + \|\bar{v}_n\|_{1,2,\infty,\Omega^t}^2 \leq \alpha(t, a) [\|\bar{H}'\|_{0,2,\infty,\Omega^t}^2 (\varepsilon \|\bar{H}'\|_{3,2,2,\Omega^t}^2 \\ + c(\varepsilon)t(\|\bar{H}'(0)\|_{0,\Omega}^2 + \|\bar{H}'_t\|_{0,2,2,\Omega^t}^2)) + c\|\bar{f}\|_{0,2,2,\Omega^t}^2 + c\|\bar{v}_n(0)\|_{1,\Omega}^2 \\ + c\|\bar{v}_n\|_{0,2,2,\Omega^t}^2 + \varepsilon\|\bar{v}_{nt}\|_{1,2,2,\Omega^t}^2 + cp_0^2],$$

where  $\alpha$  is an increasing positive function.

*Proof.* Multiplying (3.2) by  $\frac{d}{dt} c_{kn}$  and summing over  $k$  from 1 to  $n$  we get

$$(3.11) \quad \int_{\Omega} \left( \frac{1}{2} \bar{v}_{nt}^2 + \mathbb{D}_u(\bar{v}_n) \mathbb{D}_u(\bar{v}_{nt}) \right) d\xi \\ = \mu_1 \int_{\Omega} (\bar{H}' \nabla_u \bar{H}' \bar{v}_{nt} - \nabla_u \bar{H}'^2 \bar{v}_{nt}) d\xi + \int_{\Omega} \bar{f} \bar{v}_{nt} d\xi - \int_S p_0 \bar{n}_u \bar{v}_{nt} d\xi_S.$$

Using in (3.11) the Hölder and Young inequalities we get

$$(3.12) \quad \|\bar{v}_{nt}\|_{0,\Omega}^2 + \frac{d}{dt} \|\mathbb{D}_u(\bar{v}_n)\|_{0,\Omega}^2 \leq \varphi(a) \left( \|\bar{H}'\|_{0,\Omega}^2 \|\bar{H}'\|_{2,\Omega}^2 + \varepsilon \|\bar{v}_{nt}\|_{0,\Omega}^2 \right. \\ \left. + \int_{\Omega} |\bar{u}_\xi| |\mathbb{D}_u(\bar{v}_n)| |\bar{v}_{n\xi}| d\xi \right) + c\|p_0 \bar{n}_u\|_{0,S}^2 + c\|\bar{f}\|_{0,\Omega}^2 + \varepsilon \|\bar{v}_{nt}\|_{1,\Omega}^2.$$



Integrating (3.12) with respect to time and applying the Korn and Gronwall inequalities we get (3.10).

To estimate  $\|\bar{v}_{nt}\|_{1,2,2,\Omega^t}^2$  we need the following result.

LEMMA 3.3. *Let the assumptions of Lemma 3.1 be satisfied and  $\bar{f}_t \in L_2(0, T, L_2(\Omega))$ ;  $\frac{1}{H'}, \frac{1}{H'_t} \in L_\infty(0, T, H^1(\Omega))$ ;  $\bar{v}_t(0) \in L_2(\Omega)$ . Then*

$$(3.13) \quad \|\bar{v}_{nt}\|_{0,2,\infty,\Omega^t}^2 + \|\bar{v}_{nt}\|_{1,2,2,\Omega^t}^2 \leq \alpha(a)[(\varepsilon\|\bar{u}\|_{3,2,2,\Omega^t}^2 + c(\varepsilon)t(\|\bar{u}_t\|_{0,2,2,\Omega^t}^2 + \|\bar{u}(0)\|_{0,\Omega}^2))(\|\bar{v}_n\|_{1,2,\infty,\Omega^t}^2 + cp_0^2 + \|\frac{1}{H'}\|_{1,2,\infty,\Omega^t}^4) + \|\frac{1}{H'_t}\|_{1,2,\infty,\Omega^t}^2(\varepsilon\|\frac{1}{H'}\|_{3,2,2,\Omega^t}^2 + c(\varepsilon)t(\|\frac{1}{H'}(0)\|_{0,\Omega}^2 + \|\frac{1}{H'_t}\|_{0,2,2,\Omega^t}^2))] + c\|\bar{v}_{nt}\|_{0,2,2,\Omega^t}^2 + c\|\bar{f}_t\|_{0,2,2,\Omega^t}^2 + \|\bar{v}_{nt}(0)\|_{0,\Omega}^2,$$

where  $\alpha$  is an increasing positive function.

*Proof.* Differentiating (3.2) with respect to  $t$ , multiplying by  $\frac{d}{dt}c_{kn}$ , summing over  $k$  from 1 to  $n$  and using the Korn, Hölder and Young inequalities we get

$$(3.14) \quad \frac{d}{dt}\|\bar{v}_{nt}\|_{0,\Omega}^2 + \|\bar{v}_{nt}\|_{1,\Omega}^2 \leq \varphi(a)[\|\bar{u}_\xi\|_{\infty,\Omega}(\varepsilon\|\bar{v}_{nt}\|_{1,\Omega}^2 + \|\bar{v}_n\|_{1,\Omega}^2 + \|\frac{1}{H'}\|_{1,\Omega}^4) + \|\frac{1}{H'_t}\|_{1,\Omega}^2\|\frac{1}{H'}\|_{1,\Omega}^2 + \varepsilon\|\bar{v}_{nt}\|_{1,\Omega}^2] + c\|\bar{f}_t\|_{0,\Omega}^2 + c\|(p_0\bar{n}_u)_t\|_{0,S}^2 + \|\bar{v}_{nt}\|_{0,\Omega}^2.$$

Integrating (3.14) with respect to time we get (3.13).

From (3.7), (3.10), (3.13) we have

LEMMA 3.4. *Let the assumptions of Lemmas 3.1–3.3 be satisfied. Then*

$$(3.15) \quad \|\bar{v}_n\|_{1,2,2,\Omega^t}^2 + \|\bar{v}_{nt}\|_{1,2,2,\Omega^t}^2 + \|\bar{v}_n\|_{1,2,\infty,\Omega^t}^2 + \|\bar{v}_{nt}\|_{0,2,\infty,\Omega^t}^2 \leq \alpha(a, t)[(\varepsilon\|\frac{1}{H'}\|_{3,2,2,\Omega^t}^2 + c(\varepsilon)t(\|\frac{1}{H'}(0)\|_{0,\Omega}^2 + \|\frac{1}{H'_t}\|_{0,2,2,\Omega^t}^2)) \cdot (\|\frac{1}{H'}\|_{0,2,\infty,\Omega^t}^2 + \|\frac{1}{H'_t}\|_{1,2,\infty,\Omega^t}^2) + (\varepsilon\|\bar{u}\|_{3,2,2,\Omega^t}^2 + c(\varepsilon)t(\|\bar{u}_t\|_{0,2,2,\Omega^t}^2 + \|\bar{u}(0)\|_{0,\Omega}^2))(cp_0^2 + \|\frac{1}{H'}\|_{1,2,\infty,\Omega^t}^4) + c\|\bar{f}\|_{0,2,2,\Omega^t}^2 + c\|\bar{f}_t\|_{0,2,2,\Omega^t}^2 + \|\bar{v}_n(0)\|_{1,\Omega}^2 + cp_0^2] + \|\bar{v}_{nt}(0)\|_{0,\Omega}^2 \equiv \bar{F},$$

where  $\alpha$  is an increasing positive function.

Now choosing a subsequence and letting  $n \rightarrow \infty$  we get

LEMMA 3.5. Assume that  $\bar{u} \in L_2(0, T, H^3(\Omega))$ ;  $\bar{v}(0) \in H^1(\Omega)$ ;  $\bar{v}_t(0) \in L_2(\Omega)$ ;  $\bar{H}' \in L_\infty(0, T, H^1(\Omega)) \cap L_2(0, T, H^3(\Omega))$ ;  $\bar{H}'_t \in L_\infty(0, T, H^1(\Omega)) \cap L_2(0, T, L_2(\Omega))$  and  $\bar{f}_1, \bar{f}_t \in L_2(0, T, L_2(\Omega))$ . Then there exists a weak solution of problem (3.2) such that

$$\begin{aligned} \bar{v} &\in L_\infty(0, T, H^1(\Omega)) \cap L_2(0, T, H^1(\Omega)), \\ \bar{v}_t &\in L_2(0, T, H^1(\Omega)) \cap L_\infty(0, T, L_2(\Omega)) \end{aligned}$$

and

$$(3.16) \quad \|\bar{v}\|_{1,2,2,\Omega^t}^2 + \|\bar{v}_t\|_{1,2,2,\Omega^t}^2 + \|\bar{v}\|_{1,2,\infty,\Omega^t}^2 + \|\bar{v}_t\|_{0,2,\infty,\Omega^t}^2 \leq \bar{F}.$$

To show that  $\bar{v} \in L_2(0, T, H^3(\Omega))$  we consider the following elliptic problem:

$$(3.17) \quad \begin{aligned} -\operatorname{div}_u \mathbb{T}_u(\bar{v}, \bar{p}) &= -\bar{v}_t + \mu_1 \bar{H}' \nabla_u \bar{H}' + \mu_1 \nabla_u \bar{H}'^2 + \bar{f} && \text{in } \Omega^T, \\ \operatorname{div}_u \bar{v} &= 0 && \text{in } \Omega^T, \\ \bar{n}_u \mathbb{T}_u(\bar{v}, \bar{p}) &= -p_0 \bar{n}_u && \text{on } S^T. \end{aligned}$$

Applying to (3.17) the regularization technique for elliptic problems (see [4]) and Lemma 3.5 we get

$$(3.18) \quad \begin{aligned} &\|\bar{v}\|_{3,\Omega}^2 + \|\bar{p}\|_{2,\Omega}^2 \\ &\leq \varphi(a) (\|\bar{H}'\|_{2,\Omega}^4 + \|\bar{H}'\|_{\infty,\Omega}^2 \|\bar{H}'\|_{2,\Omega}^2 \\ &\quad + \|\bar{H}'\|_{1,\Omega}^2 \|\bar{H}'\|_{3,\Omega}^2) + \|\bar{v}_t\|_{1,\Omega}^2 + \|\bar{f}\|_{1,\Omega}^2 + \|p_0 \bar{n}_u\|_{3/2,S}^2. \end{aligned}$$

Integrating (3.18) with respect to time we get

LEMMA 3.6. Let the assumptions of Lemma 3.5 be satisfied and  $\bar{H}' \in L_\infty(0, T, H^2(\Omega))$ ;  $\bar{f} \in L_2(0, T, H^1(\Omega))$ . Then

$$(3.19) \quad \begin{aligned} \|\bar{v}\|_{3,2,2,\Omega^t}^2 + \|\bar{p}\|_{2,2,2,\Omega^t}^2 &\leq \alpha(a, t) [(\|\bar{H}'\|_{3,2,2,\Omega^t}^2 + \|\bar{H}'\|_{2,2,\infty,\Omega^t}^2) \\ &\quad \cdot (\varepsilon \|\bar{H}'\|_{3,2,2,\Omega^t}^2 + c(\varepsilon)t(\|\bar{H}'(0)\|_{0,\Omega}^2 + \|\bar{H}'_t\|_{0,2,2,\Omega^t}^2)) \\ &\quad + cp_0^2 + \|\bar{v}_t\|_{1,2,2,\Omega^t}^2 + \|\bar{f}\|_{1,2,2,\Omega^t}^2], \end{aligned}$$

where  $\alpha$  is an increasing positive function.

To show that  $\bar{v}_t \in L_2(0, T, H^2(\Omega))$  we differentiate (3.17) with respect to time. Then we get the following elliptic problem:

$$\begin{aligned}
 (3.20) \quad & -\operatorname{div}_u \mathbb{T}_u(\bar{v}_t, \bar{p}_t) = (\operatorname{div}_u)_t \mathbb{T}_u(\bar{v}, \bar{p}) + \operatorname{div}_u(\mathbb{T}_u)_t(\bar{v}, \bar{p}) - \bar{v}_{tt} \\
 & + (\mu_1 \bar{H}' \nabla_u \bar{H}')_t - (\mu_1 \nabla_u \bar{H}'^2)_t + \bar{f}_t \quad \text{in } \Omega^T, \\
 & \operatorname{div}_u \bar{v}_t = -(\operatorname{div}_u)_t \bar{v} \quad \text{in } \Omega^T, \\
 & \bar{n}_u \mathbb{T}_u(\bar{v}_t, \bar{p}_t) = -p_0(\bar{n}_u)_t - (\bar{n}_u \mathbb{T}_u)_t(\bar{v}, \bar{p}), \quad \text{on } S^T.
 \end{aligned}$$

Applying to (3.20) the regularization technique for elliptic problems we get (see [4])

$$\begin{aligned}
 (3.21) \quad & \|\bar{v}_t\|_{2,\Omega}^2 + \|\bar{p}_t\|_{1,\Omega}^2 \\
 & \leq \varphi(a)[\|\bar{u}\|_{2,\Omega}^2(\|\bar{v}\|_{3,\Omega}^2 + \|\bar{p}\|_{2,\Omega}^2 \\
 & \quad + \|\bar{H}'\|_{2,\Omega}^4) + \|\bar{H}'\|_{1,\Omega}^2 \|\bar{H}'\|_{2,\Omega}^2 + \|\bar{H}'_t\|_{2,\Omega}^2 \|\bar{H}'\|_{1,\Omega}^2 \\
 & \quad + cp_0^2 \|\bar{u}\|_{1,\Omega}^2] + \|\bar{v}_{tt}\|_{0,\Omega}^2 + \|\bar{f}_t\|_{0,\Omega}^2.
 \end{aligned}$$

Integrating (3.21) with respect to time and using Lemmas 3.5–3.6 we get

LEMMA 3.7. *Let the assumptions of Lemmas 3.5–3.6 be satisfied and  $\bar{u} \in L_\infty(0, T, H^2(\Omega))$ ;  $\bar{H}'_t \in L_2(0, T, H^2(\Omega))$ ;  $\bar{H}'_{tt}, \bar{f}_t \in L_2(0, T, L_2(\Omega))$ ;  $\bar{H}'(0), \bar{H}'_t(0) \in L_2(\Omega)$ . Then*

$$\begin{aligned}
 (3.22) \quad & \|\bar{v}_t\|_{2,2,2,\Omega^t}^2 + \|\bar{p}_t\|_{1,2,2,\Omega^t}^2 \\
 & \leq \alpha(a, t)[\|\bar{u}\|_{2,2,\infty,\Omega^t}^2(\|\bar{v}\|_{3,2,2,\Omega^t}^2 + \|\bar{p}\|_{1,2,2,\Omega^t}^2) \\
 & \quad + (\|\bar{H}'\|_{2,2,\infty,\Omega^t}^2(1 + \|\bar{u}\|_{2,2,\infty,\Omega^t}^2) + \|\bar{H}'_t\|_{2,2,2,\Omega^t}^2) \\
 & \quad \cdot (\varepsilon(\|\bar{H}'\|_{3,2,2,\Omega^t}^2 + \|\bar{H}'_t\|_{2,2,2,\Omega^t}^2) \\
 & \quad + c(\varepsilon)t(\|\bar{H}'_t\|_{0,2,2,\Omega^t}^2 + \|\bar{H}'_{tt}\|_{0,2,2,\Omega^t}^2 + \|\bar{H}'(0)\|_{0,\Omega}^2 \\
 & \quad + \|\bar{H}'_t(0)\|_{0,\Omega}^2)) + \|\bar{v}_{tt}\|_{0,2,2,\Omega^t}^2 + cp_0^2 + \|\bar{f}_t\|_{0,2,2,\Omega^t}^2],
 \end{aligned}$$

where  $\alpha$  is an increasing positive function.

Next we have to obtain an estimate for  $\|\bar{v}_{tt}\|_{1,2,2,\Omega^t}$ . Let

$$\begin{aligned}
 \phi_n & \in L_2(0, T, H^3(\Omega)) \cap L_\infty(0, T, H^1(\Omega)), \\
 \phi_{nt} & \in L_\infty(0, T, H^1(\Omega)) \cap L_2(0, T, H^2(\Omega))
 \end{aligned}$$

and suppose that

$$\phi_n = \sum_{k=1}^n c_k(t) \varphi_k(\xi) \in A_n,$$

where  $\{\varphi_n\}$  is an orthonormal basis in  $H^3(\Omega)$ . Then  $\bigcup_{n=1}^\infty A_n$  is dense in  $L_2(0, T, H^3(\Omega)) \cap L_\infty(0, T, H^1(\Omega))$  and  $\bigcup_{n=1}^\infty A_{nt}$  is dense in  $L_\infty(0, T, H^1(\Omega)) \cap L_2(0, T, H^2(\Omega))$ . Let  $\bar{v}_n \in \bigcup_{k=1}^n A_k$ . Then from the weak solution of the linearized problem (1.12), where  $\bar{n}_u \mathbb{T}(\bar{v}_n, \bar{q}) = 0$  on  $S^T$  and  $\bar{q} = \bar{p} - p_0$  in  $\Omega^T$ , if we let  $n \rightarrow \infty$  we get

LEMMA 3.8. *Let the assumptions of Lemmas 3.5–3.7 be satisfied and  $\bar{u}_t \in L_2(0, T, H^2(\Omega)) \cap L_\infty(0, T, H^1(\Omega))$ ;  $\frac{1}{2} \bar{H}'_{tt} \in L_2(0, T, H^1(\Omega))$ ;  $\bar{v}_{tt}(0) \in L_2(\Omega)$  and  $\bar{u}_{tt}, \bar{f}_{tt} \in L_2(0, T, L_2(\Omega))$ . Then*

$$\begin{aligned}
 (3.23) \quad & \|\bar{v}_{tt}\|_{0,2,\infty,\Omega^t}^2 + \|\bar{v}_{tt}\|_{1,2,2,\Omega^t}^2 \\
 & \leq \alpha(a, t)[(\varepsilon\|\bar{v}\|_{3,2,2,\Omega^t}^2 + c(\varepsilon)t(\|\bar{v}_t\|_{0,2,2,\Omega^t}^2 + \|\bar{v}(0)\|_{0,\Omega}^2))\|\bar{u}_t\|_{1,2,\infty,\Omega^t}^2 \\
 & \quad + \varepsilon\|\bar{H}'_{tt}\|_{1,2,2,\Omega^t}^4 + ct^2\|\bar{H}'_t\|_{1,2,2,\Omega^t}^4 + c\|\bar{H}'(0)\|_{2,\Omega}^4 \\
 & \quad + (\varepsilon\|\bar{H}'_t\|_{2,2,2,\Omega^t}^2 + c(\varepsilon)t(\|\bar{H}'_{tt}\|_{0,2,2,\Omega^t}^2 + \|\bar{H}'_t(0)\|_{0,\Omega}^2)) \\
 & \quad \cdot (\|\bar{u}\|_{2,2,\infty,\Omega^t}^2\|\bar{H}'\|_{2,2,\infty,\Omega^t}^2 + \|\bar{H}'_t\|_{1,2,2,\Omega^t}^2) \\
 & \quad + \|\bar{H}'\|_{2,2,\infty,\Omega^t}^2(\varepsilon\|\bar{u}_t\|_{2,2,2,\Omega^t}^2 + c(\varepsilon)t(\|\bar{u}_{tt}\|_{0,2,2,\Omega^t}^2 + \|\bar{u}_t(0)\|_{0,\Omega}^2)) \\
 & \quad + \|\bar{f}_{tt}\|_{0,2,2,\Omega^t}^2 + \|\bar{v}_{tt}(0)\|_{0,\Omega}^2],
 \end{aligned}$$

where  $\alpha$  is an increasing positive function.

Now adding inequalities (3.22) and (3.23) for  $p = q + p_0$  we get

LEMMA 3.9. *Let the assumptions of Lemmas 3.5–3.8 be satisfied. Then*

$$\begin{aligned}
 (3.24) \quad & \|\bar{v}_{tt}\|_{0,2,\infty,\Omega^t}^2 + \|\bar{v}_{tt}\|_{1,2,2,\Omega^t}^2 + \|\bar{v}_t\|_{2,2,2,\Omega^t}^2 + \|\bar{p}_t\|_{1,2,2,\Omega^t}^2 \\
 & \leq \alpha(a, t)[(t\|\bar{u}_t\|_{2,2,2,\Omega^t}^2 + \|\bar{u}(0)\|_{2,\Omega}^2)(\|\bar{v}\|_{3,2,2,\Omega^t}^2 + \|\bar{p}\|_{2,2,2,\Omega^t}^2) \\
 & \quad + (\varepsilon\|\bar{v}\|_{3,2,2,\Omega^t}^2 + c(\varepsilon)t(\|\bar{v}_t\|_{0,2,2,\Omega^t}^2 + \|\bar{v}(0)\|_{0,\Omega}^2))\|\bar{u}_t\|_{1,2,\infty,\Omega^t}^2 \\
 & \quad + (\varepsilon(\|\bar{H}'_t\|_{2,2,2,\Omega^t}^2 + \|\bar{H}'\|_{3,2,2,\Omega^t}^2) + c(\varepsilon)t(\|\bar{H}'\|_{0,2,2,\Omega^t}^2 \\
 & \quad + \|\bar{H}'_{tt}\|_{0,2,2,\Omega^t}^2 + \|\bar{H}'(0)\|_{0,\Omega}^2 + \|\bar{H}'_t(0)\|_{0,\Omega}^2)) \\
 & \quad \cdot (\|\bar{H}'\|_{2,2,\infty,\Omega^t}^2(1 + \|\bar{u}\|_{2,2,\infty,\Omega^t}^2) + \|\bar{H}'_t\|_{2,2,2,\Omega^t}^2) \\
 & \quad + \|\bar{H}'\|_{2,2,\infty,\Omega^t}^2(\varepsilon\|\bar{u}_t\|_{2,2,2,\Omega^t}^2 + c(\varepsilon)t(\|\bar{u}_{tt}\|_{0,2,2,\Omega^t}^2 + \|\bar{u}_t(0)\|_{0,\Omega}^2)) \\
 & \quad + \|\bar{f}_{tt}\|_{0,2,2,\Omega^t}^2 + \|\bar{v}_{tt}(0)\|_{0,\Omega}^2 + cp_0^2 + \|\bar{f}_t\|_{0,2,2,\Omega^t}^2],
 \end{aligned}$$

where  $\alpha$  is an increasing positive function.

From Lemmas 3.5–3.9 we get  $\bar{v} \in L_2(0, T, H^3(\Omega)) \cap L_\infty(0, T, H^1(\Omega))$ ;  $\bar{v}_t \in L_\infty(0, T, H^1(\Omega)) \cap L_2(0, T, H^2(\Omega))$ ;  $\bar{v}_{tt} \in L_\infty(0, T, L_2(\Omega)) \cap L_2(0, T, H^1(\Omega))$  and  $\bar{p} \in L_2(0, T, H^2(\Omega))$ ,  $\bar{p}_t \in L_2(0, T, H^1(\Omega))$ .

**4. Existence and regularity of solutions of the linearized problem (1.13)**

LEMMA 4.1. *Assume that  $\bar{E}_* \in L_2(0, T, L_2(B))$ ;  $\bar{H}(0) \in L_2(\Pi)$ ;  $H_* \in L_2(0, T, H^1(B))$ . Then for solutions of (3.3) the following inequality holds:*

$$(4.1) \quad \|\bar{H}_n\|_{0,2,\infty,\Pi^t}^2 + \|\bar{H}_n\|_{1,2,2,\Pi^t}^2 \leq \alpha(a, t)[\|\bar{E}_*\|_{0,2,2,B^t}^2 + \|\bar{H}_n(0)\|_{0,\Pi}^2 + \|\bar{H}_*\|_{1,2,2,B^t}^2],$$

where  $\alpha$  is an increasing positive function.

*Proof.* Multiplying (3.3) by  $d_{kn}$  and summing over  $k$  from 1 to  $n$  we get

$$(4.2) \quad \int_{\Pi} \left( \frac{1}{2} \mu \frac{d}{dt} \bar{H}_n^2 + \frac{1}{\sigma} |\text{rot}_u \bar{H}_n|^2 \right) d\xi = \int_{\Omega} \mu(\bar{u} \times \bar{H}_n) \text{rot}_u \bar{H}_n d\xi + \int_{\Pi} \mu \bar{u} \nabla_u \bar{H}_n \bar{H}_n d\xi + \frac{1}{\sigma_2} \int_B (\bar{n}_u \times \bar{E}_*) \bar{H}_n d\xi_B.$$

Applying to (4.2) the Hölder and Young inequalities and the inequality (see Appendix)

$$(4.3) \quad \|\bar{H}_n\|_{1,\Pi}^2 \leq c(\delta)(\|\text{rot}_u \bar{H}_n\|_{0,\Pi}^2 + \|\bar{H}_*\|_{1,B}^2),$$

we get

$$(4.4) \quad \frac{d}{dt} \|\bar{H}_n\|_{0,\Pi}^2 + c \|\bar{H}_n\|_{1,\Pi}^2 \leq \varphi(a)(\|\bar{u}\|_{\infty,\Pi} \|\bar{H}_n\|_{0,\Pi}^2 + \|\bar{E}_*\|_{0,B}^2 + \|\bar{H}_*\|_{1,B}^2).$$

Integrating (4.4) with respect to time and using the Gronwall inequality we get (4.1).

LEMMA 4.2. *Let the assumptions of Lemma 4.1 be satisfied and  $\bar{H}(0) \in H^1(\Pi)$ ;  $\bar{H}_* \in L_\infty(0, T, H^1(B))$ . Then*

$$(4.5) \quad \|\bar{H}_{nt}\|_{0,2,2,\Pi^t}^2 + \|\bar{H}_n\|_{1,2,\infty,\Pi^t}^2 \leq \alpha(a, t)[\varepsilon \|\bar{H}_{nt}\|_{1,2,2,\Pi^t}^2 + \|\bar{E}_*\|_{0,2,2,B^t}^2 + \|\bar{H}_n(0)\|_{1,\Pi}^2 + \|\bar{H}_*\|_{1,2,\infty,B^t}^2],$$

where  $\alpha$  is an increasing positive function.

*Proof.* Multiplying (3.3) by  $\frac{d}{dt} d_{kn}$  and summing over  $k$  from 1 to  $n$  we get,

$$(4.6) \quad \int_{\Pi} \left( \frac{1}{2} \mu \bar{H}_{nt}^2 + \frac{1}{\sigma} \text{rot}_u \bar{H}_n \text{rot}_u \bar{H}_{nt} \right) d\xi = \int_{\Omega} \mu(\bar{u} \times \bar{H}_n) \text{rot}_u \bar{H}_{nt} d\xi + \int_{\Pi} \mu \bar{u} \nabla_u \bar{H}_n \bar{H}_{nt} d\xi + \frac{1}{\sigma_2} \int_B (\bar{n}_u \times \bar{E}_*) \bar{H}_{nt} d\xi_B.$$

Applying to (4.6) the Hölder and Young inequalities we get

$$(4.7) \quad \|\bar{H}_{nt}\|_{0,\Pi}^2 + \frac{d}{dt} \left( \frac{1}{\sigma} \|\text{rot}_u \bar{H}_n\|_{0,\Pi}^2 \right) \leq \varphi(a) (\|\bar{u}\|_{\infty,\Pi} \|\bar{H}_n\|_{1,\Pi}^2 + \varepsilon \|\bar{H}_{nt}\|_{1,\Pi}^2 + \|\bar{E}_*\|_{0,B}^2).$$

Integrating (4.7) with respect to time using (4.3) and the Gronwall inequality we get (4.5).

To estimate  $\|\bar{H}_{nt}\|_{1,2,2,\Pi^t}^2$  we need the following result:

LEMMA 4.3. *Let the assumptions of Lemma 4.1 be satisfied and  $\bar{u}_t \in L_2(0, T, H^2(\Pi))$ ;  $\bar{u}_{tt} \in L_2(0, T, L_2(\Pi))$ ;  $\bar{u} \in L_2(0, T, H^3(\Pi))$ ;  $\bar{E}_{*t} \in L_2(0, T, L_2(B))$ ;  $\bar{E}_* \in L_\infty(0, T, H^1(B))$ ;  $\bar{H}_t(0) \in L_2(\Pi)$ ; and  $\bar{H}_{*t} \in L_2(0, T, H^1(B))$ . Then*

$$(4.8) \quad \begin{aligned} \|\bar{H}_{nt}\|_{0,2,\infty,\Pi^t}^2 + \|\bar{H}_{nt}\|_{1,2,2,\Pi^t}^2 &\leq \alpha(a) [(\varepsilon \|\bar{u}_t\|_{2,2,2,\Pi^t}^2 \\ &+ c(\varepsilon)t(\|\bar{u}_{tt}\|_{0,2,2,\Pi^t}^2 + \|\bar{u}_t(0)\|_{0,\Pi}^2) \|\bar{H}_n\|_{1,2,\infty,\Pi^t}^2 \\ &+ (\varepsilon \|\bar{u}\|_{3,2,2,\Pi^t}^2 + c(\varepsilon)t(\|\bar{u}_t\|_{0,2,2,\Pi^t}^2 + \|\bar{u}_t(0)\|_{0,\Pi}^2)) \\ &\cdot (\|\bar{H}_n\|_{1,2,\infty,\Pi^t}^2 + \|\bar{H}_{nt}\|_{0,2,\infty,\Pi^t}^2 + \|\bar{E}_*\|_{1,2,\infty,B^t}^2)] \\ &+ c\|\bar{E}_{*t}\|_{0,2,2,B^t}^2 + \|\bar{H}_{nt}(0)\|_{0,\Pi}^2 + \|\bar{H}_{*t}\|_{1,2,2,B^t}^2, \end{aligned}$$

where  $\alpha$  is an increasing positive function.

*Proof.* Differentiating (3.3) with respect to time, multiplying by  $\frac{d}{dt}d_{kn}$ , summing over  $k$  from 1 to  $n$  and using the Korn, Hölder and Young inequalities we get

$$(4.9) \quad \begin{aligned} \frac{1}{2} \mu \frac{d}{dt} \|\bar{H}_{nt}\|_{0,\Pi}^2 + \frac{1}{\sigma} \int_{\Pi} |\text{rot}_u(\bar{H}_{nt})|^2 d\xi \\ \leq \varphi(a) [ \|\bar{H}_n\|_{1,\Pi}^2 (\|\bar{u}_t\|_{1,\Pi}^2 + \|\bar{u}\|_{2,\Pi}^2) \\ + \varepsilon \|\bar{H}_{nt}\|_{1,\Pi}^2 + \|\bar{u}_\xi\|_{\infty,\Pi}^2 \|\bar{H}_n\|_{1,\Pi}^2 + \|\bar{H}_{nt}\|_{0,\Pi}^2 \|\bar{u}_\xi\|_{\infty,\Pi}^2 \\ + c\|\bar{E}_{*t}\|_{0,B}^2 + c\|\bar{E}_*\|_{\infty,B}^2 \|\bar{u}\|_{2,\Pi}^2 ]. \end{aligned}$$

Integrating (4.9) with respect to time using (4.3) we get (4.8). From (4.1), (4.5), (4.8) we have

LEMMA 4.4. *Let the assumptions of Lemmas 4.1–4.3 be satisfied. Then*

$$(4.10) \quad \begin{aligned} \|\bar{H}_n\|_{1,2,2,\Pi^t}^2 + \|\bar{H}_{nt}\|_{1,2,2,\Pi^t}^2 + \|\bar{H}_{nt}\|_{0,2,\infty,\Pi^t}^2 + \|\bar{H}_n\|_{1,2,\infty,\Pi^t}^2 \\ \leq \alpha(a, t) [(\varepsilon \|\bar{u}\|_{3,2,2,\Pi^t}^2 + c(\varepsilon)t(\|\bar{u}_t\|_{0,2,2,\Pi^t}^2 + \|\bar{u}(0)\|_{0,\Pi}^2)) \|\bar{E}_*\|_{1,2,\infty,B^t}^2 \\ + \|\bar{H}_n(0)\|_{1,\Pi}^2 + \|\bar{H}_{nt}(0)\|_{0,\Pi}^2 + \|\bar{E}_{*t}\|_{0,2,2,B^t}^2 + \|\bar{H}_*\|_{1,2,2,B^t}^2 \\ + \|\bar{H}_*\|_{1,2,\infty,B^t}^2 + \|\bar{H}_{*t}\|_{1,2,2,B^t}^2] \equiv \bar{G}, \end{aligned}$$

where  $\alpha$  is an increasing positive function.

Now choosing a subsequence and letting  $n \rightarrow \infty$  we get

LEMMA 4.5. *Assume that  $\bar{u} \in L_2(0, T, H^3(\Pi))$ ;  $\bar{u}_t \in L_2(0, T, H^2(\Pi))$ ;  $\bar{u}_{tt} \in L_2(0, T, L_2(\Pi))$ ;  $\bar{H}_*, \bar{H}_{*t} \in L_2(0, T, H^1(B))$ ;  $\bar{H}_*, \bar{E}_* \in L_\infty(0, T, H^1(B))$ ;  $\bar{E}_{*t} \in L_2(0, T, L_2(B))$ ;  $\bar{H}(0) \in H^1(\Pi)$ ;  $\bar{H}_t(0) \in L_2(\Pi)$ . Then there exists a weak solution of problem (3.3) such that*

$$\begin{aligned} \bar{H} &\in L_\infty(0, T, H^1(\Pi)) \cap L_2(0, T, H^1(\Pi)), \\ \bar{H}_t &\in L_2(0, T, H^1(\Pi)) \cap L_\infty(0, T, L_2(\Pi)) \end{aligned}$$

and

$$(4.11) \quad \|\bar{H}\|_{1,2,2,\Pi^t}^2 + \|\bar{H}_t\|_{1,2,2,\Pi^t}^2 + \|\bar{H}_t\|_{0,2,\infty,\Pi^t}^2 + \|\bar{H}\|_{1,2,\infty,\Pi^t}^2 \leq \bar{G}.$$

Next we introduce a partition of unity  $(\tilde{\Omega}_i, \{\xi_i\})$  with  $\Pi = \Omega \cup D = \bigcup_i \tilde{\Omega}_i$ . Let  $\tilde{\Omega}$  be one of the  $\tilde{\Omega}_i$  and  $\zeta(\xi) = \zeta_i(\xi)$  the corresponding function. If  $\tilde{\Omega} \subset \Omega$  or  $\tilde{\Omega} \subset D$ , let  $\tilde{\omega}$  be such that  $\tilde{\omega} \subset \tilde{\Omega}$  and  $\zeta(\xi) = 1$  for  $\xi \in \tilde{\omega}$ .

In a boundary subdomain  $\tilde{\Omega} \cap B \neq \emptyset$ ,  $\tilde{\omega} \cap B \neq \emptyset$ ,  $\tilde{\omega} \subset \tilde{\Omega}$  we introduce local coordinates  $\{y\}$  connected with  $\{\xi\}$  by

$$(4.12) \quad y_k = \alpha_{kl}(\xi_l - \beta_l), \quad \alpha_{3k} = n_k(\beta), \quad k = 1, 2, 3,$$

where  $\beta \in \tilde{\omega} \cap B \subset \tilde{\omega} \cap B$ ,  $\tilde{B} = \tilde{\Omega} \cap B$  and  $\{\alpha_{kl}\}$  is a constant orthogonal matrix such that  $\tilde{B}$  is determined by  $y_3 = F(y_1, y_2)$ ,  $F \in H^{5/2}$  and

$$\tilde{\Omega} = \{y : |y_i| < d, i = 1, 2, F(y') < y_3 < F(y') + d, y' = (y_1, y_2)\}.$$

Next we introduce a function  $v'$  by

$$(4.13) \quad v'_i(y) = \alpha_{ij}v_j(\xi)|_{\xi=\xi(y)},$$

where  $\xi = \xi(y)$  is the inverse transformation to (4.12).

Further, we introduce new variables by

$$(4.14) \quad z_i = y_i, \quad i = 1, 2, \quad z_3 = y_3 - \tilde{F}(y), \quad y \in \tilde{\Omega},$$

which will be denoted by  $z = \phi(y)$ , where  $\tilde{F}$  is an extension of  $F$  to  $\tilde{\Omega}$  with  $\tilde{F} \in H^{5/2}(\tilde{\Omega})$ . Let  $\hat{\Omega} = \phi(\tilde{\Omega}) = \{z : |z_i| < d, i = 1, 2, 0 < z_3 < d\}$  and  $\hat{B} = \phi(\tilde{B})$ . Define

$$(4.15) \quad \hat{v}(z) = v'(y)|_{y=\phi^{-1}(z)},$$

Introduce  $\hat{\nabla}_k = \xi_{lx_k}(\xi)z_{i\xi_l} \nabla_{z_i}|_{\xi=\chi^{-1}(z)}$ , where  $\chi(\xi) = \phi(\psi(\xi))$  and  $y = \psi(\xi)$  is described by (4.12).

If  $\tilde{\Omega} \cap S \neq \emptyset$  we similarly introduce local coordinates, but now  $\tilde{\Omega} = \{y : |y_i| < d, i = 1, 2, |F(y') - y_3| < d, y' = (y_1, y_2)\}$  and  $\beta \in S \cap \tilde{\Omega} = \tilde{S}$ .

We also introduce the following notation:

$$(4.16) \quad \tilde{v}(\xi) = \bar{v}(\xi)\zeta(\xi) \quad \text{for } \tilde{\Omega} \cap S = \emptyset \text{ and } \tilde{\Omega} \cap B = \emptyset,$$

$$(4.17) \quad \tilde{v}(z) = \hat{v}(z)\hat{\zeta}(z) \quad \text{for } z \in \hat{\Omega} = \phi(\tilde{\Omega}), \quad \tilde{\Omega} \cap B \neq \emptyset \text{ or } \tilde{\Omega} \cap S \neq \emptyset,$$

where  $\hat{\zeta}(z) = \zeta(\xi)_{\xi=\chi^{-1}(z)}$  and  $z' = \tau$ ,  $z_3 = n$ .

From (1.1), (1.4) and (1.5) we get

$$(4.18) \quad \left( \frac{1}{\sigma_1} \operatorname{rot} \overset{1}{H} - \mu_1 u \times \overset{1}{H} \right) \cdot \tau_\alpha = \left( \frac{1}{\sigma_2} \operatorname{rot} \overset{2}{H} \right) \cdot \tau_\alpha, \quad \alpha = 1, 2, \text{ on } \widetilde{S}^T.$$

In local  $z$  coordinates, (4.18) has the form

$$\left( \frac{1}{\sigma_1} \widehat{\operatorname{rot}} \overset{1}{H} - \mu_1 \widehat{u} \times \overset{1}{H} \right) \cdot \widehat{\tau}_\alpha = \left( \frac{1}{\sigma_2} \widehat{\operatorname{rot}} \overset{2}{H} \right) \cdot \widehat{\tau}_\alpha, \quad \alpha = 1, 2, \text{ on } \widetilde{S}^T.$$

Thus we get

$$(4.19) \quad \begin{aligned} \frac{1}{A_2 A_3} \left( \frac{1}{\sigma_2} \partial_3 (\overset{2}{H}_2 A_2) - \frac{1}{\sigma_1} \partial_3 (\overset{1}{H}_2 A_2) \right) &= \mu_1 (\widehat{u} \times \overset{1}{H}) \cdot \frac{\widehat{\tau}_1}{|\widehat{\tau}_1|} \quad \text{on } \widetilde{S}^T, \\ \frac{1}{A_1 A_2} \left( \frac{1}{\sigma_1} \partial_3 (\overset{1}{H}_1 A_1) - \frac{1}{\sigma_2} \partial_3 (\overset{2}{H}_1 A_1) \right) &= \mu_1 (\widehat{u} \times \overset{1}{H}) \cdot \frac{\widehat{\tau}_2}{|\widehat{\tau}_2|} \quad \text{on } \widetilde{S}^T, \end{aligned}$$

where  $A_i, i = 1, 2, 3$ , are the Lamé coefficients of the transformation  $z \mapsto x$  and from (1.1)<sub>5</sub>, (1.4)<sub>3</sub>, (1.5)<sub>2</sub> we have

$$(4.19)_1 \quad \frac{1}{\sigma_1} \partial_3 \overset{1}{H}_3 - \frac{1}{\sigma_2} \partial_3 \overset{2}{H}_3 = 0.$$

Under the above notation problem (1.1), (1.4) has the following form in a boundary subdomain  $\widetilde{\Omega} \cap S \neq \emptyset$ :

$$(4.20) \quad \begin{aligned} \mu_1 \overset{1}{H}_1 - \frac{1}{\sigma_1} \widehat{\nabla}_u^2 \overset{1}{H} &= \mu_1 ((\overset{1}{H} \cdot \widehat{\nabla}_u \widehat{u} - (\widehat{u} \cdot \widehat{\nabla}_u) \overset{1}{H}) + \mu_1 \widehat{\nabla}_u \overset{1}{H} \widehat{u} \\ &\quad + \frac{1}{\sigma_1} ((\widehat{\nabla}_u \zeta) \cdot \widehat{\nabla}_u \overset{1}{H} + \widehat{\nabla}_u \cdot (\overset{1}{H} \widehat{\nabla}_u \zeta)) \quad \text{in } \Omega \cap \widehat{\Omega}, \end{aligned}$$

$$(4.21) \quad \begin{aligned} \mu_2 \overset{2}{H}_t - \frac{1}{\sigma_2} \widehat{\nabla}_u^2 \overset{2}{H} &= \frac{1}{\sigma_2} ((\widehat{\nabla}_u \cdot \zeta) \widehat{\nabla}_u \overset{2}{H} + \widehat{\nabla}_u \cdot (\overset{2}{H} \widehat{\nabla}_u \zeta)) \\ &\quad + \mu_2 \widehat{\nabla}_u \overset{2}{H} \cdot \widehat{u} \quad \text{in } D \cap \widehat{\Omega}. \end{aligned}$$

Multiplying (4.20) by  $\overset{1}{H}$  and (4.21) by  $\overset{2}{H}$ , integrating respectively over  $\Omega \cap \widehat{\Omega}$  and  $D \cap \widehat{\Omega}$ , adding the resulting equalities, using the transmission conditions (4.19), assuming that  $|\widehat{\Omega}|$  is sufficiently small, integrating with respect to time and using the Gronwall inequality we get

$$(4.22) \quad \begin{aligned} \|\overset{1}{H}\|_{0,2,\infty,\widehat{\Omega}^t}^2 + \|\overset{1}{H}\|_{1,2,2,\widehat{\Omega}^t}^2 &\leq \alpha(\widehat{a}, t) [(\varepsilon \|\widehat{u}\|_{3,2,2,\widehat{\Omega}^t}^2 \\ &\quad + c(\varepsilon)t(\|\widehat{u}_t\|_{0,2,2,\widehat{\Omega}^t}^2 + \|\widehat{u}(0)\|_{0,\widehat{\Omega}}^2))(1 + \|\widehat{u}\|_{0,2,\infty,\widehat{\Omega}^t}^2) + \|\overset{1}{H}(0)\|_{0,\widehat{\Omega}}^2]. \end{aligned}$$

Here and below,  $\alpha$  is an increasing positive function.

Now differentiating (4.20), (4.21) with respect to  $\tau$ , multiplying respectively by  $\overset{1}{H}_\tau, \overset{2}{H}_\tau$ , integrating over  $\Omega \cap \widehat{\Omega}, D \cap \widehat{\Omega}$ , adding the resulting



equalities and using the transmission conditions (4.19) we get

$$(4.23) \quad \frac{d}{dt} \|\tilde{H}_\tau\|_{0,\hat{\Omega}}^2 + \|\tilde{H}_\tau\|_{1,\hat{\Omega}}^2 \leq \varphi(\hat{a})(\varepsilon \|\tilde{H}_\tau\|_{1,\hat{\Omega}}^2 + c \|\hat{u}\|_{2,\hat{\Omega}}^2 + c \|\hat{u}\|_{2,\hat{\Omega}}^2 \|\tilde{H}\|_{1,\hat{\Omega}}^2 \\ + \|\hat{H}\|_{1,\hat{\Omega}}^2 \|\tilde{u}_\tau\|_{1,\hat{\Omega}}^2 + \varepsilon \|\hat{H}_\tau\|_{1,\hat{\Omega}}^2 + c \|\hat{H}_\tau\|_{0,\hat{\Omega}}^2 + c\hat{a}^2 \|\tilde{H}\|_{3,\hat{\Omega}}^2 \\ + c\hat{a}(\|\tilde{H}\|_{2,\hat{\Omega}}^2 + \|\tilde{H}_\tau\|_{1,\hat{\Omega}}^2) + \|\hat{u}\|_{2,\hat{\Omega}}^2 \|\tilde{H}_\tau\|_{0,\hat{\Omega}}^2).$$

Integrating (4.23) with respect to time, using (4.22) and the Gronwall inequality we get

$$(4.24) \quad \|\tilde{H}_\tau\|_{0,2,\infty,\hat{\Omega}^t}^2 + \|\tilde{H}_\tau\|_{1,2,2,\hat{\Omega}^t}^2 \leq \alpha(\hat{a}, t)[(\varepsilon \|\hat{u}\|_{3,2,2,\hat{\Omega}^t}^2 \\ + c(\varepsilon)t(\|\hat{u}_t\|_{0,2,2,\hat{\Omega}^t}^2 + \|\hat{u}(0)\|_{0,\hat{\Omega}}^2)(\|\hat{u}\|_{0,2,\infty,\hat{\Omega}^t}^2 + 1) \\ + \|\tilde{H}(0)\|_{0,\hat{\Omega}}^2(\|\hat{u}\|_{0,2,\infty,\hat{\Omega}^t}^2 + \|\hat{u}\|_{2,2,\infty,\hat{\Omega}^t}^2) + c\hat{a}\|\tilde{H}\|_{2,2,2,\hat{\Omega}^t}^2 \\ + c\hat{a}^2\|\tilde{H}\|_{3,2,2,\hat{\Omega}^t}^2 + \|\tilde{H}_\tau(0)\|_{0,\hat{\Omega}}^2].$$

Now we estimate  $\|\tilde{H}_{nn}\|_{0,2,2,\hat{\Omega}^t}^2$ . From (4.20), (4.21) we get

$$(4.25) \quad \|\tilde{H}_{nn}\|_{0,\hat{\Omega}}^2 \leq \varphi(\hat{a})(\hat{a}\|\tilde{H}\|_{2,\hat{\Omega}}^2 + \|\tilde{H}_\tau\|_{1,\hat{\Omega}}^2 + \|\tilde{H}\|_{1,\hat{\Omega}}^2 \|\hat{u}\|_{2,\hat{\Omega}}^2 \\ + \|\tilde{u}\|_{2,\hat{\Omega}}^2 \|\hat{H}\|_{1,\hat{\Omega}}^2 + \|\hat{u}\|_{1,\hat{\Omega}}^2 + \|\hat{H}\|_{1,\hat{\Omega}}^2) + \|\tilde{H}_t\|_{0,\hat{\Omega}}^2.$$

Integrating (4.25) with respect to time we get

$$(4.26) \quad \|\tilde{H}_{nn}\|_{0,2,2,\hat{\Omega}^t}^2 \leq \alpha(\hat{a})(c\hat{a}\|\tilde{H}\|_{2,2,2,\hat{\Omega}^t}^2 + \|\tilde{H}_\tau\|_{1,2,2,\hat{\Omega}^t}^2 \\ + \|\tilde{H}\|_{1,2,2,\hat{\Omega}^t}^2 \|\hat{u}\|_{2,2,\infty,\hat{\Omega}^t}^2 + \|\tilde{u}\|_{2,2,\infty,\hat{\Omega}^t}^2 \|\hat{H}\|_{1,2,2,\hat{\Omega}^t}^2 \\ + \|\hat{u}\|_{1,2,2,\hat{\Omega}^t}^2 + \|\hat{H}\|_{1,2,2,\hat{\Omega}^t}^2) + \|\tilde{H}_t\|_{0,2,2,\hat{\Omega}^t}^2.$$

Now we estimate  $\|\tilde{H}_t\|_{0,2,2,\hat{\Omega}^t}^2$ . Multiplying (4.20), (4.21) respectively by  $\overset{1}{\tilde{H}_t}$ ,  $\overset{2}{\tilde{H}_t}$ , integrating over  $\Omega \cap \hat{\Omega}$ ,  $D \cap \hat{\Omega}$  using the transmission conditions (4.19) and integrating with respect to time we get

$$(4.27) \quad \|\tilde{H}_t\|_{0,2,2,\hat{\Omega}^t}^2 + \|\tilde{H}\|_{1,2,\infty,\hat{\Omega}^t}^2 \leq \varphi(\hat{a}, t)[(1 + \|\hat{u}\|_{2,2,\infty,\hat{\Omega}^t}^2)(\|\hat{H}\|_{1,2,2,\hat{\Omega}^t}^2 \\ + \varepsilon \|\hat{u}\|_{3,2,2,\hat{\Omega}^t}^2 + c(\varepsilon)t(\|\hat{u}_t\|_{0,2,2,\hat{\Omega}^t}^2 + \|\hat{u}(0)\|_{0,\hat{\Omega}}^2)) + \|\tilde{H}(0)\|_{1,\hat{\Omega}}^2].$$

Then from (4.22), (4.26), (4.27) we get

$$(4.28) \quad \|\tilde{H}_{nn}\|_{0,2,2,\hat{\Omega}^t}^2 \leq \alpha(\hat{a}, t)[(\varepsilon \|\hat{u}\|_{3,2,2,\hat{\Omega}^t}^2 + c(\varepsilon)t(\|\hat{u}_t\|_{0,2,2,\hat{\Omega}^t}^2 + \|\hat{u}(0)\|_{0,\hat{\Omega}}^2)) \\ \cdot (\|\hat{u}\|_{2,2,\infty,\hat{\Omega}^t}^2 + 1) + \|\tilde{H}(0)\|_{0,\hat{\Omega}}^2(t\|\hat{u}_t\|_{2,2,2,\hat{\Omega}^t}^2 + \|\hat{u}(0)\|_{2,\hat{\Omega}}^2 + 1) \\ + c\hat{a}\|\tilde{H}\|_{2,2,2,\hat{\Omega}^t}^2 + c\hat{a}^2\|\tilde{H}\|_{3,2,2,\hat{\Omega}^t}^2 + \|\tilde{H}(0)\|_{1,\hat{\Omega}}^2].$$

Now differentiating (4.20), (4.21) with respect to  $\tau$ , multiplying respectively by  $\overset{1}{\tilde{H}}_{\tau\tau}$ ,  $\overset{2}{\tilde{H}}_{\tau\tau}$ , integrating over  $\widehat{\Omega} \cap \Omega$ ,  $\widehat{\Omega} \cap D$  and integrating with respect to time using the transmission conditions (4.19) we get

$$\begin{aligned}
 (4.29) \quad & \|\tilde{H}_{\tau\tau}\|_{0,2,\infty,\widehat{\Omega}^t}^2 + \|\tilde{H}_{\tau\tau}\|_{1,2,2,\widehat{\Omega}^t}^2 \\
 & \leq \alpha(\widehat{a}, t)[((1 + \|\widehat{u}\|_{3,2,2,\widehat{\Omega}^t}^2)(\varepsilon\|\widehat{u}\|_{3,2,2,\widehat{\Omega}^t}^2 + c(\varepsilon)t(\|\widehat{u}_t\|_{0,2,2,\widehat{\Omega}^t}^2 \\
 & \quad + \|\widehat{u}(0)\|_{0,\widehat{\Omega}}^2)(\|\widehat{u}\|_{2,2,\infty,\widehat{\Omega}^t}^2 + 1) + \|\tilde{H}(0)\|_{0,\widehat{\Omega}}^2)(t\|\widehat{u}_t\|_{2,2,2,\widehat{\Omega}^t}^2 \\
 & \quad + \|\widehat{u}(0)\|_{2,2,\Omega}^2 + 1)^2 + c\widehat{a}^2\|\tilde{H}\|_{3,2,2,\widehat{\Omega}^t}^2 + c\widehat{a}\|\tilde{H}\|_{2,2,2,\widehat{\Omega}^t}^2 \\
 & \quad + c\|\widehat{u}\|_{2,2,\infty,\widehat{\Omega}^t}^2 + \|\tilde{H}(0)\|_{1,\widehat{\Omega}}^2].
 \end{aligned}$$

Now we estimate  $\tilde{H}_{nn\tau}$ . Differentiating (4.20), (4.21) with respect to  $\tau$ , multiplying respectively by  $\overset{1}{\tilde{H}}_{nn\tau}$ ,  $\overset{2}{\tilde{H}}_{nn\tau}$ , integrating over  $\widehat{\Omega} \cap \Omega$ ,  $\widehat{\Omega} \cap D$  and with respect to time we get

$$\begin{aligned}
 (4.30) \quad & \|\tilde{H}_{nn\tau}\|_{0,2,2,\widehat{\Omega}^t}^2 \leq \alpha(\widehat{a}, t)[((1 + \|\widehat{u}\|_{3,2,2,\widehat{\Omega}^t}^2)(\varepsilon\|\widehat{u}\|_{3,2,2,\widehat{\Omega}^t}^2 \\
 & \quad + c(\varepsilon)t(\|\widehat{u}_t\|_{0,2,2,\widehat{\Omega}^t}^2 + \|\widehat{u}(0)\|_{0,\widehat{\Omega}}^2)(\|\widehat{u}\|_{2,2,\infty,\widehat{\Omega}^t}^2 + 1) \\
 & \quad + \|\tilde{H}(0)\|_{0,\widehat{\Omega}}^2)(t\|\widehat{u}_t\|_{2,2,2,\widehat{\Omega}^t}^2 + \|\widehat{u}(0)\|_{2,\widehat{\Omega}}^2 + 1)^2 \\
 & \quad + c(\widehat{a} + \widehat{a}^2)\|\tilde{H}\|_{3,2,2,\widehat{\Omega}^t}^2 + c\|\widehat{u}\|_{2,2,\infty,\widehat{\Omega}^t}^2 + \|\tilde{H}(0)\|_{1,\widehat{\Omega}}^2 \\
 & \quad + \|\tilde{H}_t\|_{1,2,2,\widehat{\Omega}^t}^2].
 \end{aligned}$$

Next we estimate  $\|\tilde{H}_t\|_{1,2,2,\widehat{\Omega}^t}^2$ . Differentiating (4.20), (4.21) with respect to time, multiplying respectively by  $\overset{1}{\tilde{H}}_t$ ,  $\overset{2}{\tilde{H}}_t$ , integrating over  $\widehat{\Omega} \cap \Omega$ ,  $\widehat{\Omega} \cap D$  and with respect to time, and using the transmission conditions (4.19) we get

$$\begin{aligned}
 (4.31) \quad & \|\tilde{H}_t\|_{0,2,\infty,\widehat{\Omega}^t}^2 + \|\tilde{H}_t\|_{1,2,2,\widehat{\Omega}^t}^2 \leq \alpha(\widehat{a}, t)[((1 + \|\widehat{u}\|_{2,2,\infty,\widehat{\Omega}^t}^2 \\
 & \quad + \|\widehat{u}\|_{3,2,2,\widehat{\Omega}^t}^2 + \|\widehat{u}\|_{2,2,\infty,\widehat{\Omega}^t}^4 + \|\widehat{u}_t\|_{2,2,2,\widehat{\Omega}^t}^2 + \|\widehat{u}_{tt}\|_{0,2,2,\widehat{\Omega}^t}^2) \\
 & \quad \cdot (\varepsilon\|\widehat{u}\|_{3,2,2,\widehat{\Omega}^t}^2 + c(\varepsilon)t(\|\widehat{u}_t\|_{0,2,2,\widehat{\Omega}^t}^2 + \|\widehat{u}(0)\|_{0,\widehat{\Omega}}^2)(\|\widehat{u}\|_{2,2,\infty,\widehat{\Omega}^t}^2 + 1) \\
 & \quad + \|\tilde{H}(0)\|_{0,\widehat{\Omega}}^2)(t\|\widehat{u}_t\|_{2,2,2,\widehat{\Omega}^t}^2 + \|\widehat{u}(0)\|_{2,\widehat{\Omega}}^2 + 1)^2 \\
 & \quad + c(\widehat{a} + \widehat{a}^2)\|\tilde{H}\|_{3,2,2,\widehat{\Omega}^t}^2 + c\|\widehat{u}\|_{2,2,\infty,\widehat{\Omega}^t}^2 + \varepsilon\|\widehat{u}_t\|_{2,2,2,\widehat{\Omega}^t}^2 \\
 & \quad + c(\varepsilon)t(\|\widehat{u}_{tt}\|_{0,2,2,\widehat{\Omega}^t}^2 + \|\widehat{u}(0)\|_{0,\widehat{\Omega}}^2) + \|\tilde{H}(0)\|_{1,\widehat{\Omega}}^2].
 \end{aligned}$$

Now we estimate  $\|\tilde{H}_{nnn}\|_{0,2,2,\hat{\Omega}^t}^2$ . Differentiating (4.20), (4.21) with respect to  $n$  and multiplying respectively by  $\overset{1}{\tilde{H}}_{nnn}$ ,  $\overset{2}{\tilde{H}}_{nnn}$ , from the above inequalities we get

$$\begin{aligned}
 (4.32) \quad \|\tilde{H}_{nnn}\|_{0,2,2,\hat{\Omega}^t}^2 &\leq \alpha(\hat{a}, t)[(1 + \|\hat{u}\|_{2,2,\infty,\hat{\Omega}^t}^2 + \|\hat{u}\|_{3,2,2,\hat{\Omega}^t}^2 \\
 &\quad + \|\hat{u}\|_{2,2,\infty,\hat{\Omega}^t}^4 + \|\hat{u}_t\|_{2,2,2,\hat{\Omega}^t}^2)(\varepsilon\|\hat{u}\|_{3,2,2,\hat{\Omega}^t}^2 + c(\varepsilon)t(\|\hat{u}_t\|_{0,2,2,\hat{\Omega}^t}^2 \\
 &\quad + \|\hat{u}(0)\|_{0,\hat{\Omega}}^2))(\|\hat{u}\|_{2,2,\infty,\hat{\Omega}^t}^2 + 1) + |\tilde{H}(0)|_{0,\hat{\Omega}}^2(t\|\hat{u}_t\|_{2,2,2,\hat{\Omega}^t}^2 \\
 &\quad + \|\hat{u}(0)\|_{2,\hat{\Omega}}^2 + 1)^2 + c(\hat{a} + \hat{a}^2)\|\tilde{H}\|_{3,2,2,\hat{\Omega}^t}^2 + c\|\hat{u}\|_{2,2,\infty,\hat{\Omega}^t}^2 \\
 &\quad + \varepsilon\|\hat{u}_t\|_{2,2,2,\hat{\Omega}^t}^2 + c(\varepsilon)t(\|\hat{u}_{tt}\|_{0,2,2,\hat{\Omega}^t}^2 + \|\hat{u}_t(0)\|_{0,\hat{\Omega}}^2) + \|\tilde{H}(0)\|_{1,\hat{\Omega}}^2].
 \end{aligned}$$

Adding the above inequalities we get the estimates

$$\begin{aligned}
 (4.33) \quad \|\tilde{H}\|_{3,2,2,\hat{\Omega}^t}^2 &\leq \alpha(\hat{a}, t)[(1 + \|\hat{u}\|_{2,2,\infty,\hat{\Omega}^t}^2 + \|\hat{u}\|_{3,2,2,\hat{\Omega}^t}^2 \\
 &\quad + \|\hat{u}\|_{2,2,\infty,\hat{\Omega}^t}^4 + \|\hat{u}_t\|_{2,2,2,\hat{\Omega}^t}^2)(\varepsilon\|\hat{u}\|_{3,2,2,\hat{\Omega}^t}^2 + c(\varepsilon)t(\|\hat{u}_t\|_{0,2,2,\hat{\Omega}^t}^2 \\
 &\quad + \|\hat{u}(0)\|_{0,\hat{\Omega}}^2))(\|\hat{u}\|_{2,2,\infty,\hat{\Omega}^t}^2 + 1) + \|\tilde{H}(0)\|_{0,\hat{\Omega}}^2(t\|\hat{u}_t\|_{2,2,2,\hat{\Omega}^t}^2 \\
 &\quad + \|\hat{u}(0)\|_{2,\hat{\Omega}}^2 + 1)^2 + c\|\hat{u}\|_{2,2,\infty,\hat{\Omega}^t}^2 + \varepsilon\|\hat{u}_t\|_{2,2,2,\hat{\Omega}^t}^2 \\
 &\quad + c(\varepsilon)t(\|\hat{u}_{tt}\|_{0,2,2,\hat{\Omega}^t}^2 + \|\hat{u}_t(0)\|_{0,\hat{\Omega}}^2) + \|\tilde{H}(0)\|_{1,\hat{\Omega}}^2].
 \end{aligned}$$

Now we consider interior subdomains. Let  $\tilde{\Omega} \subset \Omega$ . Then as above we prove

$$\begin{aligned}
 (4.34) \quad \|\overset{1}{\tilde{H}}\|_{3,2,2,\tilde{\Omega}^t}^2 &\leq \alpha(a, t)[(1 + \|\bar{u}\|_{2,2,\infty,\tilde{\Omega}^t}^2 + \|\bar{u}\|_{3,2,2,\tilde{\Omega}^t}^2 \\
 &\quad + \|\bar{u}\|_{2,2,\infty,\tilde{\Omega}^t}^4)(\varepsilon\|\bar{u}\|_{3,2,2,\tilde{\Omega}^t}^2 + c(\varepsilon)t(\|\bar{u}\|_{0,2,2,\tilde{\Omega}^t}^2 \\
 &\quad + \|\bar{u}(0)\|_{0,\tilde{\Omega}}^2))(\|\bar{u}\|_{2,2,\infty,\tilde{\Omega}^t}^2 + 1) + \|\overset{1}{\tilde{H}}(0)\|_{0,\tilde{\Omega}}^2(t\|\bar{u}_t\|_{2,2,2,\tilde{\Omega}^t}^2 \\
 &\quad + \|\bar{u}(0)\|_{2,\tilde{\Omega}}^2 + 1)^2 + c\|\bar{u}\|_{2,2,\infty,\tilde{\Omega}^t}^2 + \varepsilon\|\bar{u}_t\|_{2,2,2,\tilde{\Omega}^t}^2 \\
 &\quad + c(\varepsilon)t(\|\bar{u}_{tt}\|_{0,2,2,\tilde{\Omega}^t}^2 + \|\bar{u}_t(0)\|_{0,\tilde{\Omega}}^2) + \|\overset{1}{\tilde{H}}(0)\|_{1,\tilde{\Omega}}^2].
 \end{aligned}$$

Now let  $\tilde{\Omega} \subset D$ . Then similarly to (4.32) we prove

$$(4.35) \quad \|\overset{2}{\tilde{H}}\|_{3,2,2,\tilde{\Omega}^t}^2 \leq \alpha(a, t)[(1 + \|\bar{u}\|_{2,2,\infty,\tilde{\Omega}^t}^2)\|\overset{2}{\tilde{H}}(0)\|_{0,\tilde{\Omega}}^2 + \|\overset{2}{\tilde{H}}(0)\|_{1,\tilde{\Omega}}^2].$$

Now we consider the following problem in boundary subdomains  $\widetilde{\Omega} \cap B \neq \emptyset$ :

$$(4.36) \quad \begin{cases} \mu_2(\widetilde{H}_* + \widetilde{H}^*)_t - \frac{1}{\sigma_2} \widehat{\nabla}_u^2(\widetilde{H}_* + \widetilde{H}^*) \\ = \frac{1}{\sigma_2} (\widehat{\nabla}_u \cdot \zeta) \widehat{\nabla}(\widehat{H}_* + \widehat{H}^*) + \widehat{\nabla}_u \cdot ((\widehat{H}_* + \widehat{H}^*) \widehat{\nabla}_u \zeta), \\ \widetilde{H}_* = \widetilde{H} - \widetilde{H}^*|_B = 0, \end{cases}$$

where  $\widetilde{H}^*$  is a given function. Then we get the system

$$(4.37) \quad \begin{cases} \mu_2 \widetilde{H}_{*t} - \frac{1}{\sigma_2} \widehat{\nabla}_u^2 \widetilde{H}_* = \frac{1}{\sigma_2} ((\widehat{\nabla}_u \cdot \zeta) \widehat{\nabla}_u \widehat{H}_* + \widehat{\nabla}_u \cdot (\widehat{H}_* \widehat{\nabla}_u \zeta) \\ + \frac{1}{\sigma_2} (\widehat{\nabla}_u \cdot \zeta) \widehat{\nabla}_u \widehat{H}^* + \widehat{\nabla}_u \cdot (\widehat{H}^* \widehat{\nabla}_u \zeta)) - \mu_2 \widetilde{H}_t^* + \frac{1}{\sigma_2} \widehat{\nabla}_u^2 \widetilde{H}^*, \\ \widetilde{H}_*|_B = 0. \end{cases}$$

Then as above we get

$$(4.38) \quad \begin{aligned} \|\widetilde{H}_*\|_{3,2,2,\widehat{\Omega}^t}^2 &\leq \alpha(\widehat{a}, t) [(1 + \|\widehat{u}\|_{2,2,\infty,\widehat{\Omega}^t}^2) (\|\widetilde{H}_*(0)\|_{0,\widehat{\Omega}}^2 + \|\widetilde{H}^*\|_{3,2,2,\widehat{\Omega}^t}^2) \\ &\quad + \|\widetilde{H}_t^*\|_{2,2,2,\widehat{\Omega}^t}^2 + \|\widetilde{H}_{tt}^*\|_{0,2,2,\widehat{\Omega}^t}^2 \\ &\quad + \|\widetilde{H}_{*\tau\tau}\|_{0,2,2,\widehat{B}^t}^2 + \|\widetilde{H}^*\|_{3,2,2,\widehat{B}^t}^2]. \end{aligned}$$

Then from the trace theorem we get

$$(4.39) \quad \begin{aligned} \|\widetilde{H}\|_{3,2,2,\widehat{\Omega}^t}^2 &\leq \alpha(\widehat{a}, t) [(1 + \|\widehat{u}\|_{2,2,\infty,\widehat{\Omega}^t}^2) (\|\widetilde{H}(0)\|_{0,\widehat{\Omega}}^2 + \|\widetilde{H}\|_{3,2,2,\widehat{B}^t}^2 \\ &\quad + \|\widetilde{H}(0)\|_{0,\widehat{B}}^2) + \|\widetilde{H}_t\|_{2,2,2,\widehat{B}^t}^2 + \|\widetilde{H}_{tt}\|_{0,2,2,\widehat{B}^t}^2]. \end{aligned}$$

Summing the inequalities (4.33), (4.34), (4.35), (4.39) and going back to the variables  $\xi$  we get

LEMMA 4.6. *Let the assumptions of Lemma 4.5 be satisfied and  $\overline{H}_* \in L_2(0, T, H^3(B))$ ;  $\overline{H}_{*t} \in L_2(0, T, H^2(B))$ ;  $\overline{H}_{*tt} \in L_2(0, T, L_2(B))$ . Then*

$$(4.40) \quad \begin{aligned} \|\overline{H}\|_{3,2,2,\Pi^t}^2 &\leq \alpha(a, t) [(1 + \|\overline{u}\|_{2,2,\infty,\Pi^t}^2 + \|\overline{u}\|_{3,2,2,\Pi^t}^2 \\ &\quad + \|\overline{u}\|_{2,2,\infty,\Pi^t}^4 + \|\overline{u}_t\|_{2,2,2,\Pi^t}^2) (\varepsilon \|\overline{u}\|_{3,2,2,\Pi^t}^2 + c(\varepsilon)t (\|\overline{u}_t\|_{0,2,2,\Pi^t}^2 \\ &\quad + \|\overline{u}(0)\|_{0,\Pi}^2)) (\|\overline{u}\|_{2,2,\infty,\Pi^t}^2 + 1) + \|\overline{H}(0)\|_{0,\Pi}^2 + \|\overline{H}\|_{3,2,2,B^t}^2 \\ &\quad + \|\overline{u}(0)\|_{2,\Pi}^2 + 1)^2 + ct \|\overline{u}_t\|_{2,2,2,\Pi^t}^2 + \|\overline{H}(0)\|_{0,B}^2) (t \|\overline{u}_t\|_{2,2,2,\Pi^t}^2 \\ &\quad + \|\overline{u}(0)\|_{2,\Pi}^2 + \varepsilon \|\overline{u}_t\|_{2,2,2,\Pi^t}^2 + c(\varepsilon)t (\|\overline{u}_{tt}\|_{0,2,2,\Pi^t}^2 + \|\overline{u}_t(0)\|_{0,\Pi}^2)) \\ &\quad + \|\overline{H}_t\|_{2,2,2,B^t}^2 + \|\overline{H}_{tt}\|_{0,2,2,B^t}^2 + \|\overline{H}(0)\|_{1,\Pi}^2]. \end{aligned}$$

where  $\alpha$  is an increasing positive function.

Next we have to obtain estimate for  $\|\bar{H}_{tt}\|_{1,2,2,\Pi^t}^2$ . Let  $\phi_n \in L_2(0, T, H^3(\Pi)) \cap L_\infty(0, T, H^1(\Pi))$ ,  $\phi_{nt} \in L_\infty(0, T, H^1(\Pi) \cap L_2(0, T, H^2(\Pi)))$  and  $\phi_n = \sum_{k=1}^n C_k(t)\varphi_k(\xi) \in A_n$  where  $\{\varphi_n\}$  is an orthonormal basis in  $H^3(\Pi)$ . Then  $\bigcup_{n=1}^\infty A_n$  is dense in  $L_2(0, T, H^3(\Pi)) \cap L_\infty(0, T, H^1(\Pi))$  and  $\bigcup_{n=1}^\infty A_{nt}$  is dense in  $L_2(0, T, H^2(\Pi)) \cap L_\infty(0, T, H^1(\Pi))$ .

Let  $\bar{H}_n \in \bigcup_{k=1}^n A_k$ . Then from the weak solution to the linearized problem (1.13) after letting  $n \rightarrow \infty$  we get

LEMMA 4.7. *Let the assumptions of Lemmas 4.5–4.6 be satisfied and  $\bar{E}_*, \bar{E}_{*tt} \in L_2(0, T, L_2(B))$ ;  $\bar{u}_{tt} \in L_2(0, T, H^1(\Pi))$ ;  $\bar{u}_t \in L_\infty(0, T, H^1(\Pi))$ ;  $\bar{u} \in L_\infty(0, T, H^2(\Pi))$  and  $E_{*tt} \in L_2(0, T, L_2(B))$ . Then*

$$(4.41) \quad \begin{aligned} & \|\bar{H}_{tt}\|_{0,2,\infty,\Pi^t}^2 + \|\bar{H}_{tt}\|_{1,2,2,\Pi^t}^2 \leq \alpha(a, t)[\varepsilon(\|\bar{u}_{tt}\|_{1,2,2,\Pi^t}^4 \\ & \quad + \|\bar{u}_t\|_{1,2,\infty,\Pi^t}^4 \|\bar{u}\|_{2,2,\infty,\Pi^t}^4 + \|\bar{u}_t\|_{1,2,\infty,\Pi^t}^4 + \|\bar{u}\|_{2,2,\infty,\Pi^t}^4 \\ & \quad + \|\bar{u}_{tt}\|_{1,2,2,\Pi^t}^4)(t^2\|\bar{H}_t\|_{1,2,2,\Pi^t}^4 + \|\bar{H}(0)\|_{1,\Pi}^4) + (\varepsilon\|\bar{u}\|_{3,2,2,\Pi^t}^2 \\ & \quad + c(\varepsilon)t(\|\bar{u}_t\|_{0,2,2,\Pi^t}^2 + \|\bar{u}(0)\|_{0,\Pi}^2)(\|\bar{u}_t\|_{1,2,\infty,\Pi^t}^2 + 1)\|\bar{H}_{tt}(0)\|_{0,\Pi}^2 \\ & \quad + \|\bar{u}_t\|_{1,2,\infty,\Pi^t}^2(\varepsilon\|\bar{H}\|_{3,2,2,\Pi^t}^2 + c(\varepsilon)t(\|\bar{H}_t\|_{0,2,2,\Pi^t}^2 + \|\bar{H}(0)\|_{0,\Pi}^2)) \\ & \quad + \|E_*\|_{0,2,2,B^t}^2 + \|E_{*tt}\|_{0,2,2,B^t}^2]. \end{aligned}$$

where  $\alpha$  is an increasing positive function.

Now we estimate  $\|\bar{H}_t\|_{2,2,2,\Pi^t}^2$ . As in (4.40) we prove

LEMMA 4.8. *Let the assumptions of Lemmas 4.5–4.7 be satisfied. Then*

$$(4.42) \quad \begin{aligned} & \|\bar{H}_t\|_{2,2,2,\Pi^t}^2 \leq \alpha(a, t)[(\varepsilon(\|\bar{H}\|_{3,2,2,\Pi^t}^2 + \|\bar{H}\|_{3,2,2,B^t}^2 + \|\bar{u}_t\|_{2,2,2,\Pi^t}^2) \\ & \quad + c(\varepsilon)t(\|\bar{H}_t\|_{0,2,2,\Pi^t}^2 + \|\bar{H}_{tt}\|_{0,2,2,\Pi^t}^2 + \|\bar{H}_t\|_{0,2,2,B^t}^2 + \|\bar{H}_{tt}\|_{0,2,2,B^t}^2 \\ & \quad + \|\bar{u}_t\|_{0,2,2,\Pi^t}^2 + \|\bar{u}_{tt}\|_{0,2,2,\Pi^t}^2 + \|\bar{u}(0)\|_{0,\Pi}^2 + \|\bar{u}_t(0)\|_{0,\Pi}^2 + \|\bar{H}(0)\|_{0,\Pi}^2 \\ & \quad + \|\bar{H}_t(0)\|_{0,\Pi}^2)(\|\bar{u}_t\|_{2,2,2,\Pi^t}^2 + \|\bar{u}_t\|_{1,2,\infty,\Pi^t}^2 + \|\bar{u}(0)\|_{2,\Pi}^2 \\ & \quad + \|\bar{H}(0)\|_{2,\Pi}^2) + \varepsilon\|\bar{u}_t\|_{2,2,2,\Pi^t}^4 + \|\bar{H}_t\|_{1,2,2,\Pi^t}^4 + \|\bar{H}_t\|_{1,2,\infty,B^t}^2 \\ & \quad + \|\bar{H}_t\|_{1,2,2,B^t}^4 + \|\bar{H}_t\|_{1,2,2,\Pi^t}^2 + \|\bar{H}_t\|_{1,2,2,B^t}^2 + (t\|\bar{u}_t\|_{2,2,2,\Pi^t}^2 \\ & \quad + \|\bar{u}(0)\|_{2,\Pi}^2)\|\bar{H}\|_{3,2,2,\Pi^t}^2]. \end{aligned}$$

where  $\alpha$  is an increasing positive function.

Then from Lemmas 4.5–4.8 we deduce that

$$\begin{aligned} & \bar{H} \in L_2(0, T, H^3(\Pi)) \cap L_\infty(0, T, H^1(\Pi)); \\ & \bar{H}_t \in L_\infty(0, T, H^1(\Pi)) \cap L_2(0, T, H^2(\Pi)); \\ & \bar{H}_{tt} \in L_\infty(0, T, L_2(\Pi)) \cap L_2(0, T, H^1(\Pi)). \end{aligned}$$

**5. Existence of solution of (1.1)–(1.7).** We prove the existence of a solution of (1.1)–(1.7) by the method of successive approximations determined by the problem

$$\begin{aligned}
 & \bar{v}_{m+1t} - \operatorname{div}_{\bar{v}_m} \mathbb{T}_{\bar{v}_m}(\bar{v}_{m+1}, \bar{p}_{m+1}) \\
 & \quad = \mu_1(\overset{1}{\bar{H}}_m \nabla_{\bar{v}_m} \overset{1}{\bar{H}}_m - \nabla_{\bar{v}_m} \overset{1}{\bar{H}}_m^2) + \bar{f} \quad \text{in } \Omega^T, \\
 & \operatorname{div}_{\bar{v}_m} \bar{v}_{m+1} = 0 \quad \text{in } \Omega^T, \\
 (5.1) \quad & \mu_1 \overset{1}{\bar{H}}_{m+1,t} = -\operatorname{rot}_{\bar{v}_m} \overset{1}{\bar{E}}_m + \mu_1 \nabla_{\bar{v}_m} \overset{1}{\bar{H}}_{m+1} \bar{v}_m \quad \text{in } \Omega^T, \\
 & \operatorname{rot}_{\bar{v}_m} \overset{1}{\bar{H}}_{m+1} = \sigma_1(\overset{1}{\bar{E}}_m + \mu_1 \bar{v}_m \times \overset{1}{\bar{H}}_{m+1}) \quad \text{in } \Omega^T, \\
 & \operatorname{div}_{\bar{v}_m}(\mu_1 \overset{1}{\bar{H}}_{m+1}) = 0 \quad \text{in } \Omega^T,
 \end{aligned}$$

$$\begin{aligned}
 & \bar{n}_{\bar{v}_n} \mathbb{T}_{\bar{v}_n}(\bar{v}_{m+1}, \bar{p}_{m+1}) = -p_0 \bar{n}_{\bar{v}_m} \quad \text{on } S^T, \\
 & \frac{1}{\sigma_1} \overset{1}{\bar{H}}_{m+1} = \frac{1}{\sigma_2} \overset{2}{\bar{H}}_{m+1} \quad \text{on } S^T, \\
 (5.2) \quad & \overset{1}{\bar{E}}_m \cdot \bar{\tau}_{\alpha \bar{v}_m} = \overset{2}{\bar{E}}_m \cdot \bar{\tau}_{\alpha \bar{v}_m}, \quad \alpha = 1, 2, \quad \text{on } S^T, \\
 & \bar{v}_{m+1}|_{t=0} = \bar{v}_0, \quad \overset{1}{\bar{H}}_{m+1}|_{t=0} = \overset{1}{\bar{H}}_0 \quad \text{in } \Omega,
 \end{aligned}$$

$$\begin{aligned}
 & \mu_2 \overset{2}{\bar{H}}_{m+1,t} = -\operatorname{rot}_{\bar{v}_m} \overset{2}{\bar{E}}_m + \mu_2 \nabla_{\bar{v}_m} \overset{2}{\bar{H}}_{m+1} \bar{v}_m \quad \text{in } D^T, \\
 (5.3) \quad & \sigma_2 \overset{2}{\bar{E}}_m = \operatorname{rot}_{\bar{v}_m} \overset{2}{\bar{H}}_{m+1} \quad \text{in } D^T, \\
 & \operatorname{div}_{\bar{v}_m}(\mu_2 \overset{2}{\bar{H}}_{m+1}) = 0 \quad \text{in } D^T,
 \end{aligned}$$

where in (5.3)  $\bar{v}_m$  is an extension of the function  $\bar{v}_m$  from (5.1) onto  $D^T$ , such that  $\bar{v}_m \rightarrow 0$  as  $\xi \rightarrow B$ ,

$$\begin{aligned}
 (5.4) \quad & \overset{2}{\bar{H}}_{m+1} = \bar{H}_*, \quad \overset{2}{\bar{E}}_m = \bar{E}_* \quad \text{on } B^T, \\
 & \overset{2}{\bar{H}}_{m+1}|_{t=0} = \overset{2}{\bar{H}}_0 \quad \text{in } D.
 \end{aligned}$$

First we show the boundedness of the sequence described by (5.1)–(5.4) in the norm

$$\begin{aligned}
 (5.5) \quad \beta_m(t) = & \|\bar{v}_{mt}\|_{1,\Omega}^2 + \|\bar{v}_m\|_{1,\Omega}^2 + \|\bar{v}_m\|_{3,2,2,\Omega^t}^2 + \|\bar{v}_{mt}\|_{2,2,2,\Omega^t}^2 \\
 & + \|\bar{v}_{mtt}\|_{1,2,2,\Omega^t}^2 + \|\bar{p}_m\|_{2,2,2,\Omega^t}^2 + \|\bar{p}_{mt}\|_{1,2,2,\Omega^t}^2 + \|\bar{H}_{mt}\|_{1,\Pi}^2 \\
 & + \|\bar{H}_m\|_{1,\Pi}^2 + \|\bar{H}\|_{3,2,2,\Pi^t}^2 + \|\bar{H}_{mt}\|_{2,2,2,\Pi^t}^2 + \|\bar{H}_{mtt}\|_{1,2,2,\Pi^t}^2.
 \end{aligned}$$

LEMMA 5.1. *There exists a positive function  $P$  and  $A > 0$  such that  $P(0, 0, F_0) < A$ ,  $\beta_m(0) < A$ , and there exists  $T_*$  such that  $\beta_m(t) \leq A$  for  $t \leq T_*$ ,  $m = 1, 2, \dots$ , where  $F_0$  is a function depending on  $p_0, \bar{H}_*, \bar{E}_*, \bar{H}(0), \bar{v}(0), \bar{u}(0), \bar{f}$  in the norms from Lemmas 3.1–4.8.*

*Proof.* From Lemmas 3.1–4.8 assuming that  $\varepsilon = \sqrt{t}$  and  $\bar{u} = \bar{v}_m, \bar{H}' = \bar{H}_m$ , using the fact that

$$\begin{aligned} \|\bar{u}\|_{2,2,\infty,\Omega^t}^2 &\leq ct(\|\bar{u}_t\|_{2,2,2,\Omega^t}^2 + \|\bar{u}(0)\|_{2,\Omega}^2), \\ \|\bar{H}'\|_{2,2,\infty,\Pi^t}^2 &\leq ct(\|\bar{H}'_t\|_{2,2,2,\Pi^t}^2 + \|\bar{H}(0)\|_{2,\Pi}^2), \end{aligned}$$

we get

$$(5.6) \quad \beta_{m+1}(t) \leq P(t, t^\gamma \beta_m(t), F_0),$$

where  $\gamma > 0$ . Let  $A$  be such that  $P(0, 0, F_0) < A$ . Since  $P$  is a continuous function there exists  $T_* > 0$  such that for  $t < T_*$  we get

$$(5.7) \quad P(t, t^\gamma A, F_0) \leq A.$$

From (5.7) we see that if  $\beta_m(t) \leq A$  then  $\beta_{m+1}(t) \leq A$ .

Now we prove the convergence of the sequence  $\{\bar{v}_m, \bar{H}_m\}$ . Set

$$\bar{V}_m = \bar{v}_m - \bar{v}_{m-1}, \quad \bar{\mathcal{H}}_m = \bar{H}_m - \bar{H}_{m-1}, \quad \bar{\mathcal{P}}_m = \bar{p}_m - \bar{p}_{m-1},$$

From (5.1)–(5.4), we obtain the following system of problems:

$$\begin{aligned} (5.8) \quad \bar{V}_{m+1,t} - \operatorname{div}_{\bar{v}_m} \mathbb{D}_{\bar{v}_m}(\bar{V}_{m+1}) - \nabla_{\bar{v}_m} \bar{\mathcal{P}}_{m+1} \\ = \operatorname{div}_{\bar{v}_m} [\mathbb{D}_{\bar{v}_m} - \mathbb{D}_{\bar{v}_{m-1}}](\bar{v}_m) + [\operatorname{div}_{\bar{v}_m} - \operatorname{div}_{\bar{v}_{m-1}}] \mathbb{D}_{\bar{v}_m}(\bar{v}_m) \\ + [\nabla_{\bar{v}_m} - \nabla_{\bar{v}_{m-1}}] \bar{\mathcal{P}}_m - \mu_1 \bar{H}_m [\nabla_{\bar{v}_m} - \nabla_{\bar{v}_{m-1}}] \bar{H}_m \\ - \mu_1 \bar{H}_{m-1} \nabla_{\bar{v}_{m-1}} \bar{\mathcal{H}}_m - \mu_1 \bar{\mathcal{H}}_m \nabla_{\bar{v}_{m-1}} \bar{H}_m \\ + \mu_1 \nabla_{\bar{v}_m} \bar{H}_m \bar{\mathcal{H}}_m + \mu_1 [\nabla_{\bar{v}_m} - \nabla_{\bar{v}_{m-1}}] \bar{H}_m \bar{H}_{m-1} \\ + \mu_1 \nabla_{\bar{v}_{m-1}} \bar{\mathcal{H}}_{m-1} \bar{H}_{m-1} \equiv F^*, \end{aligned}$$

$$\begin{aligned} (5.9) \quad \mu \bar{\mathcal{H}}_{m+1,t} - \frac{1}{\sigma} \operatorname{rot}_{\bar{v}_m} \operatorname{rot}_{\bar{v}_m} \bar{\mathcal{H}}_{m+1} = [\operatorname{rot}_{\bar{v}_m} - \operatorname{rot}_{\bar{v}_{m-1}}] \operatorname{rot}_{\bar{v}_m} \bar{H}_m \\ + \operatorname{rot}_{\bar{v}_{m-1}} [\operatorname{rot}_{\bar{v}_m} - \operatorname{rot}_{\bar{v}_{m-1}}] \bar{H}_m + \operatorname{rot}_{\bar{v}_m} (\bar{v}_m \times \bar{\mathcal{H}}_{m+1}) \\ + \operatorname{rot}_{\bar{v}_m} (\bar{V}_m \times \bar{H}_m) + [\operatorname{rot}_{\bar{v}_m} - \operatorname{rot}_{\bar{v}_{m-1}}] (\bar{v}_{m-1} \times \bar{H}_m) \\ + [\operatorname{rot}_{\bar{v}_m} - \operatorname{rot}_{\bar{v}_{m-1}}] \operatorname{rot}_{\bar{v}_m} \bar{H}_m + \operatorname{rot}_{\bar{v}_{m-1}} [\operatorname{rot}_{\bar{v}_m} - \operatorname{rot}_{\bar{v}_{m-1}}] \bar{H}_m \\ + \nabla_{\bar{v}_m} \bar{\mathcal{H}}_{m+1} \bar{v}_m + (\nabla_{\bar{v}_m} - \nabla_{\bar{v}_{m-1}}) \bar{H}_m \bar{V}_m \equiv K^* \end{aligned}$$

and

$$\begin{aligned}
 (5.10) \quad \bar{n}_{\bar{v}_m} \mathbb{T}_{\bar{v}_m} (\mathcal{V}_{m+1}, \bar{\mathcal{P}}_{m+1}) &= [\bar{n}_{\bar{v}_m} - \bar{n}_{\bar{v}_{m-1}}] \mathbb{D}_{\bar{v}_m} (\bar{v}_m) \\
 &\quad + \bar{n}_{\bar{v}_{m-1}} [\mathbb{D}_{\bar{v}_m} - \mathbb{D}_{\bar{v}_{m-1}}] \bar{v}_m + [\bar{n}_{\bar{v}_m} - \bar{n}_{\bar{v}_{m-1}}] \bar{\mathcal{P}}_m \\
 &\quad + p_0 [\bar{n}_{\bar{v}_m} - \bar{n}_{\bar{v}_{m-1}}] \equiv G^*.
 \end{aligned}$$

LEMMA 5.2. *Let the assumptions of Lemma 5.1 be satisfied. Then there exists  $T_{**}$  sufficiently small such that*

$$\begin{aligned}
 (5.11) \quad &\|\bar{\mathcal{V}}_{m+1}\|_{2,2,2,\Omega^t}^2 + \|\bar{\mathcal{V}}_{m+1,t}\|_{0,2,2,\Omega^t}^2 + \|\bar{\mathcal{V}}_{m+1}\|_{1,\Omega}^2 + \|\mathcal{P}_{m+1}\|_{1,2,2,\Omega^t}^2 \\
 &\quad + \|\bar{\mathcal{H}}_{m+1}\|_{0,\Pi}^2 + \|\bar{\mathcal{H}}_{m+1}\|_{1,2,2,\Pi^t}^2 \\
 &\leq \alpha(A,t)t(t+1)(\|\bar{\mathcal{V}}_m\|_{2,2,2,\Pi^t}^2 + \|\bar{\mathcal{H}}_m\|_{1,2,2,\Pi^t}^2),
 \end{aligned}$$

where  $\alpha(A,t)t(t+1) < 1$  for  $t < T_{**}$ , and  $\alpha$  is an increasing positive function.

*Proof.* Multiplying (5.8) by  $\bar{\mathcal{V}}_{m+1}$  and integrating over  $\Omega$  we get

$$\begin{aligned}
 (5.12) \quad \frac{1}{2} \int_{\Omega} \frac{d}{dt} \bar{\mathcal{V}}_{m+1}^2 d\xi + \int_{\Omega} |\mathbb{D}_{\bar{v}_m} (\bar{\mathcal{V}}_{m+1})|^2 d\xi \\
 = \int_{\Omega} F^* \bar{\mathcal{V}}_{m+1} d\xi + \int_S G^* \bar{\mathcal{V}}_{m+1} d\xi_S.
 \end{aligned}$$

We have proved that

$$\begin{aligned}
 (5.13) \quad \int_{\Omega^t} |F^* \bar{\mathcal{V}}_{m+1}| d\xi dt \leq \alpha(A)t(t+1)(\|\bar{\mathcal{V}}_m\|_{2,2,2,\Omega^t}^2 + \|\bar{\mathcal{H}}_m\|_{1,2,2,\Omega^t}^2) \\
 + \varepsilon \|\bar{\mathcal{V}}_{m+1}\|_{1,2,2,\Omega^t}^2,
 \end{aligned}$$

$$(5.14) \quad \int_{S^t} |G^* \bar{\mathcal{V}}_{m+1}| d\xi_S dt \leq \alpha(A)t(t+1)\|\bar{\mathcal{V}}_m\|_{2,2,2,\Omega^t}^2 + \varepsilon \|\bar{\mathcal{V}}_{m+1}\|_{1,2,2,\Omega^t}^2.$$

Applying to (5.12) the Korn inequality, integrating with respect to time and using (5.13), (5.14) we get

$$\begin{aligned}
 (5.15) \quad \|\bar{\mathcal{V}}_{m+1}\|_{0,\Omega}^2 + \|\bar{\mathcal{V}}_{m+1}\|_{1,2,2,\Omega^t}^2 \\
 \leq \alpha(A)t(t+1)(\|\bar{\mathcal{V}}_m\|_{2,2,2,\Omega^t}^2 + \|\bar{\mathcal{H}}_m\|_{1,2,2,\Omega^t}^2).
 \end{aligned}$$

Now using in (5.9) local  $z$  coordinates, and (4.19), as in (4.22) we get

$$(5.16) \quad \|\bar{\mathcal{H}}_{m+1}\|_{0,2,\infty,\Pi^t}^2 + \|\bar{\mathcal{H}}_{m+1}\|_{1,2,2,\Pi^t}^2 \leq \alpha(A,t) + (t+1)\|\bar{\mathcal{V}}_m\|_{2,2,2,\Pi^t}^2.$$

Multiplying (5.8) by  $\bar{\mathcal{V}}_{m+1,t}$ , integrating over  $\Omega$  and using Young and Hölder inequalities we get

$$\begin{aligned}
 (5.17) \quad &\|\bar{\mathcal{V}}_{m+1,t}\|_{0,\Omega}^2 + \frac{d}{dt} \int_{\Omega} |\mathbb{D}_{\bar{v}_m} (\bar{\mathcal{V}}_{m+1})|^2 d\xi \\
 &\leq c\|\bar{v}_m\|_{3,\Omega}^2 \int_{\Omega} |\mathbb{D}_{\bar{v}_m} (\bar{\mathcal{V}}_{m+1})|^2 d\xi + \frac{d}{dt} \int_S G^* \bar{\mathcal{V}}_{m+1} d\xi_S + \|G_t^*\|_{0,S}^2 \\
 &\quad + \|\bar{\mathcal{V}}_{m+1}\|_{1,\Omega}^2 + \|F^*\|_{0,\Omega}^2 + \|\bar{v}_m\|_{3,\Omega}^2 \|G^*\|_{0,S}^2.
 \end{aligned}$$



Integrating (5.17) with respect to time and using the Gronwall and Korn inequalities we get

$$(5.18) \quad \|\bar{\mathcal{V}}_{m+1t}\|_{0,2,2,\Omega^t}^2 + \|\bar{\mathcal{V}}_{m+1}\|_{1,\Omega}^2 \leq \alpha(A, t)(\|G_t^*\|_{0,2,2,S^t}^2 + \|\bar{\mathcal{V}}_{m+1}\|_{1,2,2,\Omega^t}^2 + \|F^*\|_{0,2,2,\Omega^t}^2 + c\|G^*\|_{0,2,\infty,S^t}^2 + \varepsilon\|\bar{\mathcal{V}}_{m+1}\|_{1,2,\infty,\Omega^t}^2).$$

We have proved that

$$(5.19) \quad \|F^*\|_{0,2,2,\Omega^t}^2 \leq \alpha(A)t(t+1)(\|\bar{\mathcal{V}}_m\|_{2,2,2,\Omega^t}^2 + \|\bar{\mathcal{H}}_m\|_{1,2,2,\Omega^t}^2),$$

$$(5.20) \quad \|G^*\|_{0,S}^2 \leq \alpha(A)t(t+1)\|\bar{\mathcal{V}}_m\|_{2,2,2,\Omega^t}^2,$$

$$(5.21) \quad \|G_t^*\|_{0,2,2,S^t}^2 \leq \alpha(A)t(t+1)\|\bar{\mathcal{V}}_m\|_{2,2,2,\Omega^t}^2.$$

Using (5.19)–(5.21) in (5.18) we get

$$(5.22) \quad \|\bar{\mathcal{V}}_{m+1,t}\|_{0,2,2,\Omega^t}^2 + \|\bar{\mathcal{V}}_{m+1}\|_{1,\Omega}^2 \leq \alpha(A, t)[t(t+1)(\|\bar{\mathcal{V}}_m\|_{2,2,2,\Omega^t}^2 + \|\bar{\mathcal{H}}_m\|_{1,2,2,\Omega^t}^2) + \|\bar{\mathcal{V}}_{m+1}\|_{1,2,2,\Omega^t}^2 + \varepsilon\|\bar{\mathcal{V}}_{m+1}\|_{1,2,\infty,\Omega^t}^2],$$

where  $\alpha$  is an increasing positive function.

Now applying the regularization technique to the elliptic problem (cf. [4])

$$\begin{aligned} -\operatorname{div}_{\bar{v}_m} \mathbb{T}(\bar{\mathcal{V}}_{m+1}, \bar{\mathcal{P}}_{m+1}) &= F^* + \bar{\mathcal{V}}_{m+1,t} && \text{in } \Omega^T, \\ \bar{n}_{\bar{v}_m} \mathbb{T}_{\bar{v}_m}(\bar{\mathcal{V}}_{m+1}, \bar{\mathcal{P}}_{m+1}) &= G^* && \text{on } S^T, \\ \bar{\mathcal{V}}_{m+1}|_{t=0} &= 0 && \text{in } \Omega, \end{aligned}$$

we get

$$(5.23) \quad \|\bar{\mathcal{V}}_{m+1}\|_{2,2,2,\Omega^t}^2 + \|\bar{\mathcal{P}}_{m+1}\|_{1,2,2,\Omega^t}^2 \leq c(\|F^*\|_{0,2,2,\Omega^t}^2 + \|G^*\|_{1/2,2,2,S^t}^2 + \|\bar{\mathcal{V}}_{m+1,t}\|_{0,2,2,\Omega^t}^2).$$

We have proved that

$$(5.24) \quad \|G^*\|_{1/2,2,2,S^t}^2 \leq \alpha(A)t(t+1)\|\bar{\mathcal{V}}_m\|_{2,2,2,\Omega^t}^2.$$

Now using in (5.23) inequalities (5.19), (5.24) we get

$$(5.25) \quad \|\bar{\mathcal{V}}_{m+1}\|_{2,2,2,\Omega^t}^2 + \|\bar{\mathcal{P}}_{m+1}\|_{1,2,2,\Omega^t}^2 \leq \alpha(A)t(t+1)(\|\bar{\mathcal{V}}_m\|_{2,2,2,\Omega^t}^2 + \|\bar{\mathcal{H}}_m\|_{1,2,2,\Omega^t}^2) + c\|\bar{\mathcal{V}}_{m+1t}\|_{0,2,2,\Omega^t}^2.$$

Now summing up inequalities (5.15), (5.22), (5.25) we get

$$(5.26) \quad \|\bar{\mathcal{V}}_{m+1,t}\|_{0,2,2,\Omega^t}^2 + \|\bar{\mathcal{V}}_{m+1}\|_{1,\Omega}^2 + \|\bar{\mathcal{V}}_{m+1}\|_{2,2,2,\Omega^t}^2 + \|\bar{\mathcal{P}}_{m+1}\|_{1,2,2,\Omega^t}^2 \leq \alpha(A, t)t(t+1)[\|\bar{\mathcal{V}}_m\|_{2,2,2,\Omega^t}^2 + \|\bar{\mathcal{H}}_m\|_{1,2,2,\Omega^t}^2].$$

Combining (5.16)–(5.26) we obtain (5.11).

From Lemmas 5.1 and 5.2 we have

**THEOREM 5.1.** *Let the assumptions of Lemmas 5.1 and 5.2 be satisfied. Then there exists  $T^{**}$  sufficiently small such that for  $T \leq T^{**}$  there exists*

a solution to problem (1.1)–(1.7) such that

$$\begin{aligned} \bar{v} &\in L_2(0, T, H^3(\Omega)) \cap L_\infty(0, T, H^1(\Omega)); \\ \bar{v}_t &\in L_\infty(0, T, H^1(\Omega)) \cap L_2(0, T, H^2(\Omega)); \\ \bar{v}_{tt} &\in L_2(T, 0, H^1(\Omega)); \quad \bar{p} \in L_2(0, T, H^2(\Omega)); \\ p_t &\in L_2(0, T, H^1(\Omega)); \quad \bar{H} \in L_2(0, T, H^3(\Pi)) \cap L_\infty(0, T, H^1(\Pi)); \\ \bar{H}_t &\in L_\infty(0, T, H^1(\Pi)) \cap L_2(0, T, H^2(\Pi)); \quad \bar{H}_{tt} \in L_2(0, T, H^1(\Pi)) \end{aligned}$$

and

$$\begin{aligned} &\|\bar{v}_t\|_{1,2,\infty,\Omega^T}^2 + \|\bar{v}\|_{1,2,\infty,\Omega^T}^2 + \|\bar{v}\|_{3,2,2,\Omega^T}^2 + \|\bar{v}_t\|_{2,2,2,\Omega^T}^2 + \|\bar{v}_{tt}\|_{1,2,2,\Omega^T}^2 \\ &\quad + \|\bar{p}\|_{2,2,2,\Omega^T}^2 + \|\bar{p}_t\|_{1,2,2,\Omega^T}^2 + \|\bar{H}_t\|_{1,2,\infty,\Pi^T}^2 + \|\bar{H}\|_{1,2,\infty,\Pi^T}^2 \\ &\quad + \|\bar{H}\|_{3,2,2,\Pi^T}^2 + \|\bar{H}_t\|_{2,2,2,\Pi^T}^2 + \|\bar{H}_{tt}\|_{1,2,2,\Pi^T}^2 \leq A, \end{aligned}$$

where  $A$  is defined in Lemma 5.1.

### Appendix

LEMMA. Let  $\bar{H}^1(t) \in H^1(\Omega)$  and  $\bar{H}^2(t) \in H^1(D)$  for  $t \in [0, T]$  and suppose assumptions (1.1)–(1.7) are satisfied. Then

$$(1) \quad \|\bar{H}\|_{1,\Pi} \leq c(\|\text{rot } \bar{H}\|_{0,\Pi} + \|\bar{H}_*\|_{1,B}).$$

*Proof.* Using the transmission conditions we get the problems

$$\begin{aligned} &\text{rot } \bar{H}^1 = \text{rot } \bar{H}^2 && \text{in } \Omega, \\ &\text{rot } \bar{H}^1 = \text{rot } \bar{H}^2 && \text{in } D, \\ (2) \quad &\text{div } \bar{H}^1 = 0 && \text{in } \Omega, \\ &\text{div } \bar{H}^2 = 0 && \text{in } D, \\ &\frac{1}{\sigma_1} \bar{n} \bar{H}^1|_S = \frac{1}{\sigma_2} \bar{n} \bar{H}^2|_S && \text{on } S. \end{aligned}$$

From [8, pp. 20–21] solutions of (2) will be sought in the form

$$(3) \quad \bar{H}^1 = \nabla \bar{\varphi}^1 + \bar{u}, \quad \bar{H}^2 = \nabla \bar{\varphi}^2 + \bar{u},$$

where  $\bar{\varphi}^1, \bar{\varphi}^2$  are solutions of the Neumann problems

$$\begin{aligned} &\Delta \bar{\varphi}^1 = 0, \quad \Delta \bar{\varphi}^2 = 0, \\ (4) \quad &\frac{\partial \bar{\varphi}^1}{\partial \bar{n}} \Big|_S = \frac{\sigma_1}{\sigma_2} \bar{n} \bar{H}^2, \quad \frac{\partial \bar{\varphi}^2}{\partial \bar{n}} \Big|_S = \bar{n} \bar{H}^1, \quad \frac{\partial \bar{\varphi}^2}{\partial \bar{n}} \Big|_B = \bar{n} \bar{H}^* \end{aligned}$$

and  $\overset{1}{u}, \overset{2}{u}$  are solutions of the problems

$$(5) \quad \begin{aligned} \operatorname{rot} \overset{1}{u} &= \operatorname{rot} \overset{1}{H}, & \operatorname{rot} \overset{2}{u} &= \operatorname{rot} \overset{2}{H}, \\ \operatorname{div} \overset{1}{u} &= 0, & \operatorname{div} \overset{2}{u} &= 0, \\ \bar{n}\overset{1}{u}|_{\partial\Omega} &= 0, & \bar{n}\overset{2}{u}|_{\partial D} &= 0. \end{aligned}$$

There are vectors  $\overset{1}{e}, \overset{2}{e}$  such that

$$(6) \quad \overset{1}{u} = \operatorname{rot} \overset{1}{e}, \quad \operatorname{div} \overset{1}{e} = 0, \quad \bar{\tau} \cdot \overset{1}{e}|_S = 0,$$

where  $\bar{\tau}$  is any tangent vector to  $S$ , and

$$(7) \quad \overset{2}{u} = \operatorname{rot} \overset{2}{e}, \quad \operatorname{div} \overset{2}{e} = 0, \quad \bar{\tau} \cdot \overset{2}{e}|_{\partial D} = 0,$$

where  $\bar{\tau}$  is any tangent vector to  $\partial D$ . Thus problems (6), (7) can be replaced by problems

$$(8) \quad \begin{aligned} -\Delta \overset{1}{e} &= \operatorname{rot} \overset{1}{H}, & -\Delta \overset{2}{e} &= \operatorname{rot} \overset{2}{H}, \\ \bar{\tau} \cdot \overset{1}{e}|_S &= 0, & \bar{\tau} \cdot \overset{2}{e}|_{\partial D} &= 0, \\ \operatorname{div} \overset{1}{e}|_S &= 0, & \operatorname{div} \overset{2}{e}|_{\partial D} &= 0. \end{aligned}$$

From (2), (4) we get

$$(9) \quad \varphi = \overset{1}{\varphi} + \overset{2}{\varphi} = \int_B \frac{\partial}{\partial n} (\bar{n}\overset{1}{H}_*) \frac{1}{|\xi - x|} dx_B.$$

Using the partition of unity and local coordinates we write (8) in the form

$$(10) \quad -\widehat{\Delta}_u \overset{k}{e} = \widehat{\zeta} \widehat{\operatorname{rot}}_u \overset{k}{H} + 2\widehat{\nabla}_u \overset{k}{e} \widehat{\nabla}_u \widehat{\zeta} + \overset{k}{e} \widehat{\Delta}_u \widehat{\zeta},$$

$$(11) \quad \overset{k}{e} \cdot \widehat{\tau}|_{n=0} = 0, \quad \widehat{\operatorname{div}}_u \overset{k}{e}|_{n=0} = \widehat{\nabla}_u \widehat{\zeta} \overset{k}{e}, \quad k = 1, 2,$$

where  $\overset{k}{e}(z) = \widehat{e}(z)\widehat{\zeta}(z)$ ,  $k = 1, 2$ .

Now applying to (10), (11) Theorem 10.2 of [1], summing the inequalities using Lemma 10.5 of [8] and next going back to variables  $\xi$  we get

$$(12) \quad \|\overset{1}{e}\|_{2,\Omega} \leq c \|\operatorname{rot} \overset{1}{H}\|_{0,\Omega}, \quad \|\overset{2}{e}\|_{2,D} \leq c \|\operatorname{rot} \overset{2}{H}\|_{0,D}.$$

Now from (6), (7), we get

$$(13) \quad \begin{aligned} \|\overset{1}{u}\|_{1,\Omega} &\leq c \|\overset{1}{e}\|_{2,\Omega} \leq c \|\operatorname{rot} \overset{1}{H}\|_{0,\Omega}, \\ \|\overset{2}{u}\|_{1,D} &\leq c \|\overset{2}{e}\|_{2,D} \leq c \|\operatorname{rot} \overset{2}{H}\|_{0,D}. \end{aligned}$$

Then finally from (9), (13) we have

$$\begin{aligned}\|\bar{H}\|_{1,\Pi} &= \|\nabla\varphi + u\|_{1,\Pi} \leq c(\|\nabla\varphi\|_{1,\Pi} + \|u\|_{1,\Pi}) \\ &\leq c(\|\nabla\varphi\|_{1,\Pi} + \|\operatorname{rot} \bar{H}\|_{0,\Pi}) \leq c(\|\bar{H}_*\|_{1,B} + \|\operatorname{rot} \bar{H}\|_{0,\Pi}).\end{aligned}$$

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