## On existence of a unique generalized solution to systems of elliptic PDEs at resonance

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#### Abstract

The Dirichlet boundary value problem for systems of elliptic partial differential equations at resonance is studied. The existence of a unique generalized solution is proved using a new min-max principle and a global inversion theorem.


1. Introduction. In this paper, we consider the following Dirichlet boundary value problem:

$$
\begin{equation*}
L u+g(x, u)=h(x), \quad x \text { in } \Omega, \quad u=0, \quad x \text { on } \partial \Omega, \tag{1.1}
\end{equation*}
$$

where $\Omega \subseteq \mathbb{R}^{n}$ is a bounded domain with smooth boundary, $g: \Omega \times \mathbb{R}^{m} \rightarrow$ $\mathbb{R}^{m}$ is continuous and continuously differentiable with respect to $u$, and $h \in L^{2}(\Omega)$.

By a global inversion theorem and a non-variational version of the minmax principle, the existence of a unique generalized solution to (1.1) has been proved for $L=\Delta$ in QL. The main contribution of this paper is that the existence of a unique solution of (1.1) at resonance can be derived for much more general operators $L$.

Let $\Omega \subseteq \mathbb{R}^{n}$ be a bounded domain with boundary $\partial \Omega$ of class $C^{2}$, and $L$ be a second order differential operator of the form

$$
L u=\left(\begin{array}{c}
L_{1} u_{1} \\
L_{2} u_{2} \\
\vdots \\
L_{m} u_{m}
\end{array}\right), \quad L_{k} \varphi=\sum_{i=1}^{n} \sum_{j=1}^{n} \frac{\partial}{\partial x_{i}}\left(a_{i j}^{(k)}(x) \frac{\partial \varphi}{\partial x_{j}}\right), \quad k=1, \ldots, m
$$

where the operators $L_{k}, k=1, \ldots, m$, are assumed to be strongly elliptic

[^0]and symmetric, i.e.,
$$
a_{i j}^{(k)}=a_{j i}^{(k)}, \quad \exists \mu^{(k)}>0, \quad \sum_{i=1}^{n} \sum_{j=1}^{n} a_{i j}^{(k)} \zeta_{i} \zeta_{j}>\sum_{i=1}^{n} \sum_{j=1}^{n} \mu^{(k)} \zeta_{i} \zeta_{j}
$$
for all $x \in \bar{\Omega}$ and all $\zeta \in \mathbb{R}^{n} \backslash\{0\}$.
In this paper, we suppose that $a_{i j}^{(k)} \in C^{1}(\bar{\Omega})$ and
\[

A_{k}=\left(a_{i j}^{(k)}\right)_{n \times n} \in \mathbb{R}^{n \times n}, \quad A=\left($$
\begin{array}{ccc}
A_{1} & \ldots & 0 \\
\vdots & \ddots & \vdots \\
0 & \ldots & A_{m}
\end{array}
$$\right) \in \mathbb{R}^{n m \times n m}
\]

It is easy to see that $A_{k}, k=1, \ldots, m$, are symmetric matrices. Thus, $A$ is also symmetric. Moreover, $L_{k}, k=1, \ldots, m$, induce a self-adjoint differential operator in $L^{2}(\Omega)$ with domain $H_{0}^{1}(\Omega) \cap H^{2}(\Omega)$. Let $0<\lambda_{1}^{(k)}<\lambda_{2}^{(k)}<\cdots$ be all different eigenvalues of the eigenvalue problem $-L_{k} \varphi=\lambda \varphi$ in $\Omega$ with $\varphi=0$ on $\partial \Omega, k=1, \ldots, m$ (see, e.g., Gilbarg and Trudinger [GT]).

Now, we first introduce the following two lemmas which play an important role in this paper. Their detailed proofs are omitted; the interested readers are referred to [L1, LL, L2, L,

Lemma 1.1 ([L1, $\overline{\mathrm{LL}})$ ). Let $X$ and $Y$ be two closed subspaces of a real Hilbert space $H$, and $H=X \oplus Y$. Suppose that $T: H \rightarrow H$ is a $C^{1}$ mapping. If there exist continuous functions $\alpha, \beta:[0, \infty) \rightarrow(0, \infty)$ such that

$$
\begin{align*}
\left\langle T^{\prime}(u) v, v\right\rangle & \leq-\alpha(\|u\|)\|v\|^{2}  \tag{1.2}\\
\left\langle T^{\prime}(u) w, w\right\rangle & \geq \beta(\|u\|)\|w\|^{2}  \tag{1.3}\\
\left\langle T^{\prime}(u) v, w\right\rangle & =\left\langle v, T^{\prime}(u) w\right\rangle \tag{1.4}
\end{align*}
$$

for all $u \in H, v \in X, w \in Y$, and

$$
\begin{equation*}
\int_{1}^{\infty} \min \{\alpha(s), \beta(s)\} d s=\infty \tag{1.5}
\end{equation*}
$$

then $T$ is a diffeomorphism from $H$ onto $H$.
Lemma 1.2 ([L2, $\mathbb{L}])$. Let $H$ be a vector space such that $H=Z \oplus Y$ for some subspaces $Y$ and $Z$. If $Z$ is finite-dimensional and $X$ is a subspace of $H$ such that $X \cap Y=\{0\}$ and $\operatorname{dim} X=\operatorname{dim} Z$, then $H=X \oplus Y$.
2. Existence of unique solution. To show the existence of a unique solution to equation (1.1), we assume throughout this section that the following condition holds.
(C) $\partial g(x, u) / \partial u$ is symmetric and there exist continuous functions $\alpha, \beta:[0, \infty) \rightarrow(0, \infty)$ and constant symmetric $m \times m$ matrices $B_{1}$ and $B_{2}$ such that

$$
\begin{equation*}
B_{1}+\alpha(\|u\|) I \leq \frac{\partial g(x, u)}{\partial u} \leq B_{2}-\beta(\|u\|) I \tag{2.1}
\end{equation*}
$$

on $\mathbb{R}^{n} \times \mathbb{R}^{m}$ and the eigenvalues of $B_{1}$ and $B_{2}$ are $\lambda_{N_{k}}^{(k)}$ and $\lambda_{N_{k}+1}^{(k)}$, $k=1, \ldots, m$, respectively. Here $\lambda_{N_{k}}^{(k)}$ and $\lambda_{N_{k}+1}^{(k)}$ are two consecutive eigenvalues of $-L_{k} \varphi=\lambda \varphi$ in $\Omega$ with $\varphi=0$ on $\partial \Omega, k=1, \ldots, m$. Moreover

$$
\begin{equation*}
\int_{1}^{\infty} \min \{\alpha(s), \beta(s)\} d s=\infty \tag{2.2}
\end{equation*}
$$

For convenience, we introduce the following notations:
If $f, g \in L^{2}(\Omega)$, let

$$
\begin{aligned}
B[\phi, \psi] & =\int_{\Omega} \sum_{k=1}^{m} \sum_{i=1}^{n} \sum_{j=1}^{n} a_{i j}^{(k)} \frac{\partial \phi_{k}}{\partial x_{i}} \frac{\partial \psi_{k}}{\partial x_{j}} d x=\int_{\Omega} \sum_{k=1}^{m} \nabla \phi_{k}^{T} A_{k} \nabla \psi_{k} d x \\
\langle f, g\rangle_{0} & =\int_{\Omega} f^{T} g d x, \quad\|f\|_{0}^{2}=\int_{\Omega} f^{T} f d x \\
\langle u, v\rangle & =\int_{\Omega}\left[u^{T}(x) v(x)+\sum_{k=1}^{m} \sum_{i=1}^{n} \sum_{j=1}^{n} a_{i j}^{(k)} \frac{\partial u_{k}}{\partial x_{i}} \frac{\partial v_{k}}{\partial x_{j}}\right] d x \\
\|u\|^{2} & =\int_{\Omega}\left[u^{T}(x) u(x)+\sum_{k=1}^{m} \sum_{i=1}^{n} \sum_{j=1}^{n} a_{i j}^{(k)} \frac{\partial u_{k}}{\partial x_{i}} \frac{\partial u_{k}}{\partial x_{j}}\right] d x .
\end{aligned}
$$

By (1.1), we have

$$
\begin{equation*}
B[u, v]-\langle g(x, u), v\rangle_{0}=-\langle h(x), v\rangle_{0}, \quad \forall v \in H_{0}^{1}(\Omega) \tag{2.3}
\end{equation*}
$$

Main Theorem 2.1. If $g(x, u)$ satisfies condition (C) for all $u \in \mathbb{R}^{m}$, $x \in \mathbb{R}^{n}$, then there exists a unique generalized solution to equation (1.1) for every $h \in L^{2}(\Omega)$.

Proof. By the Riesz representation theorem and the above assumptions, we can define a mapping $T(u): H_{0}^{1}(\Omega) \rightarrow H_{0}^{1}(\Omega)$ as follows:

$$
\begin{equation*}
\langle T(u), v\rangle=B[u, v]-\langle g(x, u), v\rangle_{0} . \tag{2.4}
\end{equation*}
$$

Obviously, $T(u)$ is continuously differentiable for all $u \in H_{0}^{1}(\Omega)$ and

$$
\begin{equation*}
\left\langle T^{\prime}(u) w, v\right\rangle=B[w, v]-\int_{\Omega} v^{T}(x) \frac{\partial g(x, u)}{\partial u} w(x) d x \tag{2.5}
\end{equation*}
$$

Invoking the Riesz representation theorem again, there exists a function $d \in H_{0}^{1}(\Omega)$ such that

$$
\begin{equation*}
\langle d, v\rangle=-\int_{\Omega} v^{T}(x) h(x) d x \tag{2.6}
\end{equation*}
$$

By (2.3), we just need to prove that there exists a unique $u$ satisfying

$$
\begin{equation*}
T(u)=d \tag{2.7}
\end{equation*}
$$

We assume that $e_{1}, e_{2}, \ldots, e_{m}$ are eigenvectors corresponding to the eigenvalues $\lambda_{N_{1}}^{(1)}, \lambda_{N_{2}}^{(2)}, \ldots, \lambda_{N_{m}}^{(m)}$ of $B_{1}$, and $\xi_{1}, \xi_{2}, \ldots, \xi_{m}$ are eigenvectors corresponding to the eigenvalues $\lambda_{N_{1}+1}^{(1)}, \lambda_{N_{2}+1}^{(2)}, \ldots, \lambda_{N_{m}+1}^{(m)}$ of $B_{2}$, with $\left\|e_{i}\right\|_{0}=1$ and $\left\|\xi_{i}\right\|_{0}=1$ for all $i=1, \ldots, m$. And let $\tau_{i 1}^{(k)}(x), \tau_{i 2}^{(k)}(x), \ldots, \tau_{i l_{i}}^{(k)}(x)$ be eigenfunctions of the problem $-L_{k} \varphi=\lambda_{i}^{(k)} \varphi, x \in \Omega$, and $\left.\varphi\right|_{\partial \Omega}=0$.

Let

$$
\begin{aligned}
& X=\left\{v \in H_{0}^{1}(\Omega) \mid v=\sum_{k=1}^{m} b_{k}(x) e_{k}, b_{k}(x)=\sum_{i=1}^{N_{k}} \sum_{j=1}^{l_{i}} p_{i j}^{(k)} \tau_{i j}^{(k)}(x)\right\} \\
& Y=\left\{w \in H_{0}^{1}(\Omega) \mid w=\sum_{k=1}^{m} r_{k}(x) \xi_{k}, r_{k}(x)=\sum_{i=N_{k}+1}^{\infty} \sum_{j=1}^{l_{i}} q_{i j}^{(k)} \tau_{i j}^{(k)}(x)\right\}, \\
& Z=\left\{z \in H_{0}^{1}(\Omega) \mid z=\sum_{k=1}^{m} s_{k}(x) \xi_{k}, s_{k}(x)=\sum_{i=1}^{N_{k}} \sum_{j=1}^{l_{i}} q_{i j}^{(k)} \tau_{i j}^{(k)}(x)\right\}
\end{aligned}
$$

where $N_{i}, i=1, \ldots, m$, are as in condition (C) and $p_{i j}^{(k)}, q_{i j}^{(k)}$ are constants. Obviously, $H_{0}^{1}(\Omega)=Z \oplus Y$.

For all $v \in X$ and $u \in H$, we have

$$
\begin{align*}
& \left\langle T^{\prime}(u) v, v\right\rangle=B[v, v]-\int_{\Omega} v^{T}(x) \frac{\partial g(x, u)}{\partial u} v(x) d x  \tag{2.8}\\
& =\int_{\Omega} \sum_{k=1}^{m} \sum_{i=1}^{n} \sum_{j=1}^{n} a_{i j}^{(k)} \frac{\partial v_{k}}{\partial x_{i}} \frac{\partial v_{k}}{\partial x_{j}} d x-\int_{\Omega} v^{T}(x) \frac{\partial g(x, u)}{\partial u} v(x) d x \\
& \leq \int_{\Omega}\langle-L v, v\rangle_{0} d x-\int_{\Omega} v^{T}(x)\left(B_{1}+\alpha(\|u\|) I\right) v(x) d x \\
& =\int_{\Omega}\langle-L v, v\rangle_{0} d x-\int_{\Omega} v^{T}(x) B_{1} v(x) d x-\alpha(\|u\|)\|v\|_{0}^{2} \\
& \leq \sum_{k=1}^{m} \lambda_{N_{k}}^{(k)} \int_{\Omega} b_{k}^{2}(x) d x-\int_{\Omega} \sum_{i=1}^{m} \sum_{j=1}^{m} b_{i}(x) b_{j}(x) e_{i}^{T} B_{1} e_{j} d x-\alpha(\|u\|)\|v\|_{0}^{2} \\
& \leq-\alpha(\|u\|)\|v\|_{0}^{2} .
\end{align*}
$$

Since

$$
\begin{align*}
\|v\|^{2} & =\|v\|_{0}^{2}+B[v, v]=\|v\|_{0}^{2}-\langle L v, v\rangle_{0}  \tag{2.9}\\
& \leq\|v\|_{0}^{2}+\sum_{k=1}^{m} \lambda_{N_{k}}^{(k)} b_{k}^{2}(x)\left\langle e_{k}, e_{k}\right\rangle_{0} \leq(1+M)\|v\|_{0}^{2}
\end{align*}
$$

where $M=\max \left\{\lambda_{N_{k}}^{(k)}: k=1, \ldots, m\right\}$, we have

$$
\begin{equation*}
\left\langle T^{\prime}(u) v, v\right\rangle \leq-\frac{\alpha(\|u\|)}{M+1}\|v\|^{2} \tag{2.10}
\end{equation*}
$$

Similarly, denoting $N=\max \left\{\lambda_{N_{k}+1}^{(k)}: k=1, \ldots, m\right\}$, we have

$$
\begin{equation*}
\left\langle T^{\prime}(u) w, w\right\rangle \geq \frac{\beta(\|u\|)}{N+1}\|w\|^{2}, \quad \forall w \in Y, \forall u \in H \tag{2.11}
\end{equation*}
$$

By the symmetry of $B$, it is easy to see that

$$
\begin{equation*}
\left\langle T^{\prime}(u) v, w\right\rangle=\left\langle v, T^{\prime}(u) w\right\rangle \tag{2.12}
\end{equation*}
$$

Let $\alpha_{1}(s)=\frac{\alpha(s)}{M+1}$ and $\beta_{1}(s)=\frac{\beta(s)}{N+1}$; then

$$
\begin{equation*}
\delta(s)=\min \left\{\alpha_{1}(s), \beta_{1}(s)\right\} \geq \frac{\min \{\alpha(s), \beta(s)\}}{N+1} \tag{2.13}
\end{equation*}
$$

By (2.2), it is easy to see that

$$
\begin{equation*}
\int_{1}^{\infty} \delta(s) d s=\infty \tag{2.14}
\end{equation*}
$$

Obviously, $X \cap Y=\{0\}$ and $\operatorname{dim} X=\operatorname{dim} Z=\sum_{k=1}^{m} \sum_{i=1}^{N_{k}} l_{i}$. From $H=Z \oplus Y$ and Lemma 1.2, we get $H=X \oplus Y$. Thus, $T(u)$ is a diffeomorphism from $H_{0}^{1}(\Omega)$ onto itself by Lemma 1.1. Therefore, there exists a unique $u$ satisfying $T(u)=d$, that is, $u$ is the unique generalized solution to (1.1).
3. An example. Consider the following coupled partial differential equations:

$$
\left\{\begin{align*}
2 \frac{\partial^{2} u_{1}}{\partial x^{2}}+\frac{\partial^{2} u_{1}}{\partial y^{2}}+\frac{5}{2} u_{1}+3 \sin ^{2} \gamma(x, y) u_{1} & -\frac{1}{2} u_{2}  \tag{3.1}\\
+\frac{1}{2} \ln \left(u_{1}+u_{2}+\sqrt{1+\left(u_{1}+u_{2}\right)^{2}}\right) & =h_{1}(x, y) \\
\frac{\partial^{2} u_{2}}{\partial x^{2}}+\frac{\partial^{2} u_{2}}{\partial y^{2}}+\frac{5}{2} u_{2}+2 \cos ^{2} \gamma(x, y) u_{2} & -\frac{1}{2} u_{1} \\
+\frac{1}{2} \ln \left(u_{1}+u_{2}+\sqrt{1+\left(u_{1}+u_{2}\right)^{2}}\right) & =h_{2}(x, y)
\end{align*}\right.
$$

for $(x, y) \in \Omega=[0, \pi] \times[0, \pi] \subseteq \mathbb{R}^{2}$, and $u(x, y)=\left(u_{1}(x, y), u_{2}(x, y)\right) \in \mathbb{R}^{2}$, with boundary conditions $u(0, y)=u(\pi, y)=u(x, 0)=u(x, \pi)=u(0,0)=$ $0, h_{1}, h_{2} \in L^{2}(\Omega)$. Since

$$
\begin{aligned}
& \frac{\partial g(x, u)}{\partial u}= \\
& \left(\begin{array}{cc}
\frac{5}{2}+3 \sin ^{2} \gamma+\frac{1}{2 \sqrt{1+\left(u_{1}+u_{2}\right)^{2}}} & -\frac{1}{2}+\frac{1}{2 \sqrt{1+\left(u_{1}+u_{2}\right)^{2}}} \\
-\frac{1}{2}+\frac{1}{2 \sqrt{1+\left(u_{1}+u_{2}\right)^{2}}} & \frac{5}{2}+2 \cos ^{2} \gamma+\frac{1}{2 \sqrt{1+\left(u_{1}+u_{2}\right)^{2}}}
\end{array}\right)
\end{aligned}
$$

we have $\gamma=\gamma(x, y)$. It is easy to see that

$$
\begin{aligned}
& \left(\begin{array}{cc}
\frac{5}{2} & -\frac{1}{2} \\
-\frac{1}{2} & \frac{5}{2}
\end{array}\right)+\left(\begin{array}{cc}
\frac{1}{2 \sqrt{1+\left(u_{1}+u_{2}\right)^{2}}} & \frac{1}{2} \\
0 & 2 \sqrt{1+\left(u_{1}+u_{2}\right)^{2}}
\end{array}\right) \\
& \leq \frac{\partial g(x, u)}{\partial u} \leq \\
& \left(\begin{array}{cc}
6 & 0 \\
0 & 5
\end{array}\right)-\left(\begin{array}{cc}
\frac{1}{2}-\frac{1}{2 \sqrt{1+\left(u_{1}+u_{2}\right)^{2}}} & \frac{1}{2}-\frac{1}{2 \sqrt{1+\left(u_{1}+u_{2}\right)^{2}}}
\end{array}\right)
\end{aligned}
$$

Here, the set of all eigenvalues of $L_{1}$ under the given boundary conditions is $C_{1}=\left\{2 m^{2}+n^{2}: n\right.$ and $m$ are natural numbers $\}=\{3,6,9,11,12,17, \ldots\}$, and the set of all eigenvalues of $L_{2}$ under the given boundary conditions is $C_{2}=\left\{m^{2}+n^{2}: n\right.$ and $m$ are natural numbers $\}=\{2,5,8,10,13,17, \ldots\}$.
The eigenvalues of $B_{1}$ are 3 and 2 , and the eigenvalues of $B_{2}$ are 6 and 5 . Obviously, the integral is divergent, $\int_{0}^{\infty} d s / \sqrt{1+s^{2}}=\infty$. If $\alpha(s)=\frac{1}{2} \sqrt{1+s^{2}}$ and $\beta(s)=\frac{1}{2}\left(1-1 / \sqrt{1+s^{2}}\right)$, then the conditions of Theorem 2.1 are satisfied. Thus, there exists a unique generalized solution to equation (3.1) for arbitrary $h_{1}(x, y), h_{2}(x, y) \in L^{2}(\Omega)$.

REmark. It should be highlighted that the existence of a unique generalized solution to our system cannot be deduced from the previously known results [AP, BF, HN, INW, IN, LS].

Acknowledgements. This research was partly supported by Fund of Oceanic Telemetry Engineering and Technology Research Center, State Oceanic Administration (grant no. 2012003), and NSFC (61101208, 61002048).

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Received 21.11.2011
and in final form 12.11.2012


[^0]:    2010 Mathematics Subject Classification: Primary 35J47; Secondary 35J61.
    Key words and phrases: PDE, resonance, existence, unique, elliptic system.

