# THE VON NEUMANN ALGEBRA ASSOCIATED WITH AN INFINITE NUMBER OF $t$-FREE NONCOMMUTATIVE GAUSSIAN RANDOM VARIABLES 

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#### Abstract

We show that the von Neumann algebras generated by an infinite number of $t$ deformed free gaussian operators are factors of type $I_{\infty}$.


1. Introduction. In [6] we constructed, for each positive real number $t$, families of non-commutative random variables associated with the central limit measures for $t$ transformed classical and free convolutions. In this paper we shall study the families related to $t$-transformed free convolution, in the von Neumann algebras' framework. Let us briefly recall the constructions.

For $t \geq 0$ and a given separable Hilbert space $\mathcal{H}$, (being the complexification of a real Hilbert space $\mathcal{H}_{\mathbb{R}}$ ), with the scalar product $\langle\cdot \mid \cdot\rangle$, we consider the Fock space

$$
\mathcal{F}_{t}(\mathcal{H})=\mathbb{C} \Omega \oplus \bigoplus_{n \geq 1} \mathcal{H}^{\otimes n}
$$

completed with respect to the following scalar product $\langle\cdot \mid \cdot\rangle_{t}$ :

$$
\left\langle x_{1} \otimes x_{2} \otimes \ldots \otimes x_{n} \mid y_{1} \otimes y_{2} \otimes \ldots \otimes y_{k}\right\rangle_{t}=\delta_{n, k} \cdot t^{n-1} \cdot \prod_{j=1}^{n}\left\langle x_{j} \mid y_{j}\right\rangle, \quad\langle\Omega \mid \Omega\rangle_{t}=1
$$

Now, given a vector $f \in \mathcal{H}_{\mathbb{R}}$, we define a creation operator $B_{t}(f)$ and annihilation operator $A_{t}(f)$ on $\mathcal{F}_{t}(\mathcal{H})$. For arbitrary $x_{1}, x_{2}, \ldots, x_{n} \in \mathcal{H}$ we put

$$
B_{t}(f) x_{1} \otimes x_{2} \otimes \ldots \otimes x_{n}=f \otimes x_{1} \otimes x_{2} \otimes \ldots \otimes x_{n}, \quad B_{t}(f) \Omega=f
$$

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where $n \geq 1$, and

$$
\begin{gathered}
A_{t}(f) \Omega=0, \quad A_{t}(f) x_{1}=\left\langle x_{1} \mid f\right\rangle \Omega \\
A_{t}(f) x_{1} \otimes x_{2} \otimes \ldots \otimes x_{n}=t \cdot\left\langle x_{1} \mid f\right\rangle x_{2} \otimes \ldots \otimes x_{n}
\end{gathered}
$$

where $n \geq 2$. Then for every $f \in \mathcal{H}_{\mathbb{R}}$ the operators $A_{t}(f)$ and $B_{t}(f)$ are bounded by $\max \{1, \sqrt{t}\} \cdot\|f\|$, and are adjoint to each other and we shall consider the self-adjoint operators $G_{t}(f)=A_{t}(f)+B_{t}(f)$, which are thus bounded by $2 \cdot \max \{1, \sqrt{t}\}\|f\|$.

Definition 1.1. By $\mathcal{M}_{t}$ we shall denote the von Neumann algebra generated by the set $\left\{G_{t}(f): f \in \mathcal{H}\right\}$, that is its double commutant in $\mathcal{B}\left(\mathcal{F}_{t}(\mathcal{H})\right)$, the $\mathrm{C}^{*}$-algebra of all bounded operators on $\mathcal{H}$ :

$$
\mathcal{M}_{t}:=\left\{G_{t}: f \in \mathcal{H}\right\}^{\prime \prime} \subset \mathcal{B}\left(\mathcal{F}_{t}(\mathcal{H})\right)
$$

For $t=1$, which is the case of free convolution, it was shown by Voiculescu that $\mathcal{M}_{1}=V N\left(\mathbb{F}_{k}\right)$ is the von Neumann type $\mathrm{II}_{1}$ factor of the free group $\mathbb{F}_{k}$ on $k$ free generators. We shall show that for $t \neq 1$ the situation is quite different, if the number of operators $G_{t}(f)$ is infinite.
2. Main result. Our main result is the following.

Theorem 2.1. For $0<t \neq 1$ the von Neumann algebra $\mathcal{M}_{t}$ is the type $I_{\infty}$ factor $\mathcal{M}_{t}=$ $\mathcal{B}\left(\mathcal{F}_{t}(\mathcal{H})\right)$.

In proving the theorem the crucial role is played by the fact that the orthogonal projection $P$ onto the vacuum $\Omega$ is in $\mathcal{M}_{t}$. This follows from the following:

Lemma 2.2. Let $S_{t}:=t \cdot I+(1-t) \cdot P$, where $P$ is the orthogonal projection onto the vacuum $\Omega$ and $I$ is the identity operator in $\mathcal{B}\left(\mathcal{F}_{t}(\mathcal{H})\right)$. Let $G_{i}:=G_{t}\left(x_{i}\right)$, where $\left\{x_{i}: i \geq 1\right\}$ is an orthonormal basis in $\mathcal{H}$, be a sequence of operators in $\mathcal{M}_{t}$. Then the sequence

$$
K_{n}:=\frac{1}{n} \sum_{i=0}^{n}\left(G_{i}\right)^{2}
$$

converges to $S_{t}$ in the strong operator topology, when $n \rightarrow \infty$.
Proof. Since the operators $G_{i}$ are all uniformly bounded, it is sufficient to show that $K_{n}(y) \rightarrow S_{t}(y)$ for any simple tensor of the form $y=x_{j_{1}} \otimes \ldots \otimes x_{j_{m}}$, with $m \geq 1$, or $y=\Omega$. It follows directly from the definition of the creation and annihilation operators $B_{t}\left(x_{i}\right)$ and $A_{t}\left(x_{i}\right)$ that:

$$
K_{n} \Omega=\Omega+\frac{1}{n} \sum_{i=0}^{n} x_{i} \otimes x_{i} \rightarrow \Omega
$$

the convergence being in the norm of $\mathcal{F}_{t}(\mathcal{H})$. On the other hand, in computing the limit of $\frac{1}{n} \sum_{i=0}^{n}\left(G_{i}\right)^{2}(y)$ we consider only the simple tensors $y \in \mathcal{F}_{t}(\mathcal{H})$ of the form $y=x_{j_{1}} \otimes$ $\ldots \otimes x_{j_{m}}, m \geq 1$, in which case $\frac{1}{n} \sum_{i=0}^{j_{1}}\left(G_{i}\right)^{2}(y) \rightarrow 0$ in norm, and $\frac{1}{n} \sum_{i=j_{1}+1}^{n}\left(G_{i}\right)^{2}(y)=$ $\frac{1}{n} \sum_{i=j_{1}+1}^{n}\left(t y+x_{i} \otimes x_{i} \otimes y\right) \rightarrow t y$ in the norm of $\mathcal{F}_{t}(\mathcal{H})$. Hence the lemma follows.

It follows from the lemma that $S_{t} \in \mathcal{M}_{t}$, hence also $\left(S_{t}\right)^{2}-t S_{t}=(1-t) P \in \mathcal{M}_{t}$. Now we shall show that the vacuum vector $\Omega$ is cyclic for $\mathcal{M}_{t}$. This will yield that the commutant $\mathcal{M}_{t}^{\prime}$ is trivial.
Lemma 2.3. The vacuum vector $\Omega$ is cyclic for the von Neumann algebra $\mathcal{M}_{t}$, which means that the linear span of the vectors $\left\{G_{t}(f) \Omega: f \in \mathcal{H}_{\mathbb{R}}\right\}$ is dense in $\mathcal{F}_{t}(\mathcal{H})$.
Proof. This follows the well known scheme used in the free case, since for any finite sequence of indices $i_{1}, i_{2}, \ldots i_{m}$ we have the formula

$$
y=G_{i_{1}} G_{i_{2}} \ldots G_{i_{m}} \Omega=x_{i_{1}} \otimes x_{i_{2}} \otimes \ldots \otimes x_{i_{m}}+\sum_{j=0}^{m-1} y_{j}
$$

where $y_{j}$ is the orthogonal projection of $y$ onto the subspace $\mathcal{H}^{\otimes j}$ of $\mathcal{F}_{t}(\mathcal{H})$, spanned by tensors of length $j$ (i.e. tensors of the form $x_{i_{1}} \otimes x_{i_{2}} \otimes \ldots \otimes x_{i_{j}}$ ). It follows by induction on $m$ that each tensor $x_{i_{1}} \otimes x_{i_{2}} \otimes \ldots \otimes x_{i_{m}}$ can be expressed as a linear combination of vectors of the form $G_{r_{1}} G_{r_{2}} \ldots G_{r_{s}} \Omega$. This proves that $\Omega$ is cyclic.

Now a standard argument shows that if the orthogonal projection onto a cyclic vector for a von Neumann algebra belongs to the algebra, then its commutant is trivial.

Proposition 2.4. The commutant $\mathcal{M}_{t}^{\prime}$ of $\mathcal{M}_{t}$ in $\mathcal{B}\left(\mathcal{F}_{t}(\mathcal{H})\right)$ consists only of multiples of identity.
Proof. For a given $K \in \mathcal{M}_{t}^{\prime}$ we have $K \Omega=K P \Omega=P K \Omega$, so $K \Omega$ is invariant for the orthogonal projection $P$ onto $\Omega$. Thus $K \Omega=c \Omega$ for some constant $c$.

Now, for a vector $f \in \mathcal{F}_{t}(\mathcal{H})$ there exists a sequence $G_{n} \in \mathcal{M}_{t}$ such that $f=\lim _{n} G_{n} \Omega$. Then

$$
K(f)=\lim _{n} K G_{n} \Omega=\lim _{n} G_{n}(K \Omega)=\lim _{n} G_{n}(c \Omega)=c \lim _{n} G_{n} \Omega=c \cdot f
$$

which proves that $K=c \cdot I$. Since $K$ was chosen arbitrary, it follows that $\mathcal{M}_{t}^{\prime}=\{c \cdot I\}$ is trivial.

Proof of Theorem 2.1. Since in $\mathcal{M}_{t}$ there is the orthogonal projection $P$ onto the vector $\Omega$ cyclic for $\mathcal{M}_{t}$, it follows from the above Lemmas and Proposition that $\mathcal{M}_{t}=\left(\left(\mathcal{M}_{t}\right)^{\prime}\right)^{\prime}=$ $\mathcal{B}\left(\mathcal{F}_{t}(\mathcal{H})\right)$. This proves the theorem.

## 3. Final remarks

Remark 3.1. The natural vacuum state $\varphi_{t}$ on $\mathcal{M}_{t}$, defined as $\varphi_{t}(K)=(K \Omega \mid \Omega)_{t}$ for $K \in \mathcal{M}_{t}$, is not tracial, since, for example, $\varphi_{t}\left(\left(G_{i}\right)^{2}\left(G_{j}\right)^{2}\right)=1$ while $\varphi_{t}\left(G_{i}\left(G_{j}\right)^{2} G_{i}\right)=t$. Of course, in general there is no trace on $\mathcal{M}_{t}$ if $\mathcal{M}_{t}=\mathcal{B}\left(\mathcal{F}_{t}(\mathcal{H})\right)$. Moreover, this state is not faithful on $\mathcal{M}_{t}$, since for $Y=1-P=Y^{*}$ we have $\varphi_{t}\left(Y^{*} Y\right)=\varphi_{t}(1-P)=0$.
Remark 3.2. Quanhua Xu [8] showed a general fact, that in the interacting Fock space determined by a sequence $\left(\lambda_{n}\right)_{n=0}^{\infty}$ the vacuum state is tracial if and only if the sequence is constant.

Remark 3.3. Recently Eric Ricard [7] has given the description of the von Neumann algebras generated by finite number of the $t$-gaussian operators.

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