

*NOTE ON BLOW-UP OF SOLUTIONS FOR A POROUS MEDIUM  
EQUATION WITH CONVECTION AND BOUNDARY FLUX*

BY

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**Abstract.** De Pablo et al. [Proc. Roy. Soc. Edinburgh Sect. A 138 (2008), 513–530] considered a nonlinear boundary value problem for a porous medium equation with a convection term, and they classified exponents of nonlinearities which lead either to the global-in-time existence of solutions or to a blow-up of solutions. In their analysis they left open the case of a certain critical range of exponents. The purpose of this note is to fill this gap.

**1. Introduction.** De Pablo et al. [3] studied the initial-boundary value problem

$$(1.1) \quad \begin{cases} u_t = (u^m)_{xx} + \lambda(u^q)_x, & x \in (0, \infty), t \in (0, T), \\ -(u^m)_x(0, t) = u^p(0, t), & t \in (0, T), \\ u(x, 0) = u_0(x), & x > 0, \end{cases}$$

where  $u_0$  is a continuous, non-negative, non-increasing and integrable function, and  $T$  is the maximal existence time which may be either finite or infinite. They established an almost complete characterization, in terms of the parameters  $m \geq 1$ ,  $p, q > 0$  and  $\lambda > 0$ , whether all solutions are global-in-time or there exist solutions which blow up in finite time. In particular, it is shown in [3] that for  $p \neq q$  there exist both non-trivial global solutions and blowing up solutions in the range  $p > \max\{q, \frac{1}{2}(m+1)\}$  while all solutions are global in the complementary range of  $p$ . If  $p = q$ , the parameter  $\lambda$  plays a role in this analysis and it has a critical value  $\lambda_c = 1$ : in the case  $\lambda > 1$  all solutions are global while for  $\lambda < 1$  all non-trivial solutions blow up if  $\frac{1}{2}(m+1) < p \leq m+1$ , and there are non-trivial global solutions and blowing up solutions if  $p > m+1$ . Finally, it is proved in [3] that in the case  $\lambda = 1$ , all solutions are global in the range  $p = q \leq m+1$ , and there are small global solutions when  $p = q > m+1$ .

As stated in [3, Remark 3.4], the question of blow-up of solutions in the range  $p = q > m+1$  and  $\lambda = 1$  was left open. Moreover, as pointed out in [3], there might not exist blowing up subsolutions of self-similar type.

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2010 *Mathematics Subject Classification*: Primary 35K55; Secondary 35B33.

*Key words and phrases*: diffusion and convection, finite time blow-up, critical exponents.

The aim of this note is to solve this problem, which together with the result from [3] provides a complete characterization of exponents  $\lambda > 0$ ,  $m \geq 1$  and  $p, q > 0$  for which problem (1.1) has either global-in-time solutions or blowing up solutions. The following theorem contains the main result of this note.

**THEOREM.** *For problem (1.1) in the case  $\lambda = 1$ ,  $m \geq 1$  and  $p = q > m + 1$ , there exists a smooth, positive, non-increasing and integrable function  $u_0$  such that the corresponding solution  $u$  of problem (1.1) blows up in finite time.*

To prove this theorem, we will employ ideas of Alikakos et al. [1] and of Fisher and Grant [4]. In our proof, the key step is to find a function whose mass is “large enough” in a neighborhood of  $x = 0$  and such that the mass of the corresponding solution remains sufficiently large on some time interval in the same region. Moreover, we shall use the sub- and super-solution approach from [3], which is based on an important observation that the comparison principle for solutions to problem (1.1) is valid as long as the initial data are strictly ordered at  $x = 0$  (see [3]).

**2. Proof of Theorem.** In the following we always assume that  $\lambda = 1$ ,  $m \geq 1$  and  $p = q > m + 1$ . First, we prove a lemma which provides a rough control of mass in the interval  $[0, 1]$  of a certain large solution.

**LEMMA 2.1.** *For fixed  $M > 0$  and  $T > 0$ , there exists a positive smooth function  $\underline{u} \in C(0, T; L^1(\mathbb{R}^+))$  such that  $\underline{u}(x, t)$  is a subsolution of problem (1.1) satisfying*

$$(2.1) \quad \int_0^1 \underline{u}(x, t) \, dx \geq M \quad \text{for } t \in (0, T).$$

*Proof.* Consider the function

$$(2.2) \quad \underline{u}(x, t) = \frac{c}{t + 1} e^{-\varepsilon x}, \quad (x, t) \in [0, \infty) \times [0, T),$$

where  $c, \varepsilon > 0$  are parameters to be determined. Assuming that  $\varepsilon \leq 1$  and choosing  $c = (T + 1)Me$ , we have

$$\int_0^1 \underline{u}(x, t) \, dx \geq \frac{c}{T + 1} \int_0^1 e^{-1} \, dx = \frac{ce^{-1}}{T + 1} \geq M.$$

Recall that a function  $\underline{u}$  is a subsolution of problem (1.1) if it satisfies

$$\begin{cases} \underline{u}_t \leq (\underline{u}^m)_{xx} + (\underline{u}^p)_x, & (x, t) \in (0, \infty) \times (0, T), \\ (\underline{u}^m)_x(0, t) + \underline{u}^p(0, t) \geq 0, & t \in (0, T). \end{cases}$$

Inserting the function  $\underline{u}$  from (2.2) into those inequalities, we obtain

$$(2.3) \quad -\frac{c}{(t+1)^2}e^{-\varepsilon x} \leq \frac{c^m}{(t+1)^m}(m\varepsilon)^2e^{-m\varepsilon x} - \frac{c^p}{(t+1)^p}p\varepsilon e^{-p\varepsilon x},$$

$$(2.4) \quad -\frac{m\varepsilon c^m}{(t+1)^m} + \frac{c^p}{(t+1)^p} \geq 0.$$

Notice that  $p > m + 1 \geq 2$ . By choosing  $\varepsilon \leq \frac{1}{pc^{p-1}}$ , we have

$$(2.5) \quad -\frac{1}{(t+1)^2}ce^{-\varepsilon x} \leq -\frac{1}{(t+1)^2}(c^p p\varepsilon)e^{-\varepsilon x} \leq -\frac{c^p}{(t+1)^p}p\varepsilon e^{-p\varepsilon x}.$$

Since  $\frac{c^m}{(t+1)^m}(m\varepsilon)^2e^{-m\varepsilon x} > 0$ , inequality (2.3) results immediately from (2.5).

On the other hand, we can also choose  $\varepsilon \leq \frac{c^{p-m}}{m(T+1)^{p-m}}$  such that

$$\frac{c^m m\varepsilon}{c^p} \leq \frac{1}{(T+1)^{p-m}} \leq \frac{1}{(t+1)^{p-m}},$$

which leads immediately to inequality (2.4).

To summarize the above arguments, if we take  $c = (T + 1)Me$  and

$$\varepsilon = \min\left\{1, \frac{1}{pc^{p-1}}, \frac{c^{p-m}}{m(T+1)^{p-m}}\right\},$$

then the function  $\underline{u}$  of (2.2) is the desired subsolution of problem (1.1). ■

With this lemma, we are now in a position to prove the main theorem.

*Proof of Theorem.* Consider the solution of the equation

$$(u^m)_{xx} + (u^p)_x = 0 \quad \text{for } 0 < x < \infty$$

given by the explicit formula  $\varphi(x) = \left(\frac{p-m}{m}x\right)^{\frac{1}{m-p}}$ . Notice that  $\varphi(x)$  blows up at  $x = 0$ .

Denote

$$M_c = \int_0^1 \varphi(x) dx = \left(\frac{p-m}{m}\right)^{\frac{1}{m-p}} \frac{p-m}{p-m-1}.$$

For  $M = 2M_c$  and for fixed  $\gamma \in (1/2, 1)$ , we choose  $T$  so large that

$$(2.6) \quad \frac{1}{T} \leq \left(\frac{1}{\gamma^{p-1}} - \frac{1}{\gamma^{m-1}}\right)\varphi^{p-1}(1).$$

For such  $M$  and  $T$ , according to Lemma 2.1, we have a positive smooth subsolution  $\underline{u}$  of problem (1.1) on  $(0, \infty) \times (0, T)$  which satisfies

$$\int_0^1 \underline{u}(x, t) dx \geq M \quad \text{for } t \in (0, T).$$

We claim that a solution  $u$  corresponding to a smooth and integrable initial datum  $u_0(x) > \underline{u}(x, 0)$  with  $(u_0^m)_x(0) + (u_0^p)(0) = 0$  blows up in finite

time. Assume on the contrary that  $u(x, t)$  is global (for a local existence result, cf. [2]). By the comparison principle for problem (1.1), we have

$$u(x, t) \geq \underline{u}(x, t) = \frac{c}{t + 1} e^{-\varepsilon x} \quad \text{for } (x, t) \in (0, \infty) \times (0, T),$$

where  $c$  and  $\varepsilon$  were determined in the proof of Lemma 2.1. Since the solution  $u$  has a positive lower bound, it is a classical solution of problem (1.1) by [5].

Now let

$$\Delta = \{(x, t) \in (0, 1) \times (0, T) \mid x - (1 - t/T) > 0\}.$$

Define the functions  $z, v : \bar{\Delta} \rightarrow \mathbb{R}$  by

$$z(x, t) = \int_0^{x-(1-t/T)} \frac{1}{\gamma} \varphi(s) ds,$$
$$v(x, t) = \int_0^x u(s, t) ds.$$

Furthermore, we define an auxiliary function  $w : \bar{\Delta} \rightarrow \mathbb{R}$  by the formula

$$w(x, t) = e^{-t}(z(x, t) - v(x, t)).$$

Since  $u$  is a classical solution of problem (1.1), we see  $w \in C^2(\Delta) \cap C(\bar{\Delta})$ .

A direct computation yields

$$w(x, t) = e^{-t}(0 - v(x, t)) \leq 0 \quad \text{for } x - (1 - t/T) = 0.$$

Using inequality (2.1) and the comparison principle, we have

$$w(1, t) = e^{-t} \left( \frac{1}{\gamma} \int_0^{t/T} \varphi(s) ds - \int_0^1 u(x, t) dx \right)$$
$$\leq e^{-t} \left( \frac{1}{\gamma} M_c - \int_0^1 \underline{u}(x, t) dx \right)$$
$$\leq e^{-t} \left( \frac{1}{\gamma} M_c - M \right) \leq 0, \quad 0 \leq t \leq T.$$

Since  $\lim_{x \rightarrow 0^+} \varphi(x) = +\infty$ , using L'Hospital's rule we derive

$$w_x(0, T) = \lim_{x \rightarrow 0^+} \frac{w(x, T) - w(0, T)}{x} = \lim_{x \rightarrow 0^+} e^{-T} \left( \frac{1}{\gamma} \varphi(x) - u(x, T) \right) = +\infty.$$

So  $w \leq 0$  in  $\partial\Delta \setminus \{(x, T) \mid 0 < x < 1\}$ . In addition, from the conditions  $w(0, T) = 0$  and  $w_x(0, T) = +\infty$ , we see that  $w(x, t)$  attains its positive maximum in the set  $\Delta \cup \{(x, T) \mid 0 < x < 1\}$  at some point  $(x_0, t_0)$ . This

implies that at  $(x_0, t_0)$ , the following relations hold true:

$$(2.7) \quad w(x_0, t_0) > 0,$$

$$(2.8) \quad w_t(x_0, t_0) \geq 0,$$

$$(2.9) \quad w_x(x_0, t_0) = 0,$$

$$(2.10) \quad w_{xx}(x_0, t_0) \leq 0.$$

Next, we show that for our choice of  $\gamma$  and  $T$  we have  $w_t(x_0, t_0) < 0$ , which contradicts (2.8). Indeed, by (2.9), we have  $z_x(x_0, t_0) = v_x(x_0, t_0)$ . Thus  $\frac{1}{\gamma}\varphi(\xi) = u(x_0, t_0)$ , where  $\xi = x_0 - (1 - t_0/T)$ . Therefore

$$\begin{aligned} v_t(x_0, t_0) &= \int_0^{x_0} u_t(x, t_0) dx = \int_0^{x_0} [(u^m)_{xx}(x, t_0) + (u^p)_x(x, t_0)] dx \\ &= (u^m)_x(x_0, t_0) + u^p(x_0, t_0) \\ &= m\left(\frac{1}{\gamma}\varphi(\xi)\right)^{m-1} u_x(x_0, t_0) + \left(\frac{1}{\gamma}\varphi(\xi)\right)^p. \end{aligned}$$

On the other hand, it follows from (2.10) that  $\frac{1}{\gamma}\varphi'(\xi) - u_x(x_0, t_0) \leq 0$ , which implies that

$$v_t(x_0, t_0) \geq m\left(\frac{1}{\gamma}\varphi\right)^{m-1} \frac{1}{\gamma}\varphi'(\xi) + \left(\frac{1}{\gamma}\varphi\right)^p(\xi).$$

Recalling that  $\varphi$  satisfies  $(\varphi^m)' + \varphi^p = 0$ , we deduce

$$v_t(x_0, t_0) \geq \left(\frac{1}{\gamma^p} - \frac{1}{\gamma^m}\right)\varphi^p(\xi),$$

which gives

$$\begin{aligned} z_t(x_0, t_0) - v_t(x_0, t_0) &\leq \frac{1}{\gamma T}\varphi(\xi) - \left(\frac{1}{\gamma^p} - \frac{1}{\gamma^m}\right)\varphi^p(\xi) \\ &= \frac{\varphi(\xi)}{\gamma} \left(\frac{1}{T} - \left(\frac{1}{\gamma^{p-1}} - \frac{1}{\gamma^{m-1}}\right)\varphi^{p-1}(\xi)\right). \end{aligned}$$

Since  $\varphi$  is decreasing, and since  $1/\gamma > 1$  and  $p > m + 1 \geq 2$ , we have

$$z_t(x_0, t_0) - v_t(x_0, t_0) \leq \frac{\varphi(\xi)}{\gamma} \left(\frac{1}{T} - \left(\frac{1}{\gamma^{p-1}} - \frac{1}{\gamma^{m-1}}\right)\varphi^{p-1}(1)\right).$$

Consequently, by the choice of  $T$  in (2.6), we obtain

$$(2.11) \quad z_t(x_0, t_0) - v_t(x_0, t_0) \leq 0.$$

Moreover, by inequality (2.7), we have

$$\begin{aligned} w_t(x_0, t_0) &= -w(x_0, t_0) + e^{-t_0}(z_t(x_0, t_0) - v_t(x_0, t_0)) \\ &< e^{-t_0}(z_t(x_0, t_0) - v_t(x_0, t_0)). \end{aligned}$$

This estimate together with (2.11) leads to the inequality  $w_t(x_0, t_0) < 0$ , which contradicts (2.8).

This contradiction implies that the solution  $u(x, t)$ , which was assumed to be global, has to blow up in finite time, no later than at time

$$\varphi^{1-p}(1) / \left( \frac{1}{\gamma^{p-1}} - \frac{1}{\gamma^{m-1}} \right) < T$$

(cf. (2.6)). ■

**Acknowledgements.** The authors would like to thank Professor Grzegorz Karch for carefully reading the manuscript and many valuable comments which improved the presentation considerably. Z. Wang is grateful to Professor Changxing Miao for his hospitality and encouragement.

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*Received 27 August 2012;*  
*revised 20 September 2012*

(5747)