

MEAN VALUE DENSITIES FOR TEMPERATURES

BY

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Abstract. A positive measurable function K on a domain D in \mathbb{R}^{n+1} is called a mean value density for temperatures if $u(0, 0) = \iint_D K(x, t)u(x, t) dx dt$ for all temperatures u on \bar{D} . We construct such a density for some domains. The existence of a bounded density and a density which is bounded away from zero on D is also discussed.

1. Let D be a bounded domain in $(n + 1)$ -dimensional Euclidean space $\mathbb{R}^{n+1} = \{(x, t); x \in \mathbb{R}^n, t \in \mathbb{R}\}$. Suppose that $(0, 0) \in \bar{D}$. We say that a measurable function $K(x, t)$ on D is a *mean value density* (at the origin with respect to the heat equation) if $K > 0$ a.e. on D and

$$(1) \quad u(0, 0) = \iint_D K(x, t)u(x, t) dx dt$$

for every temperature u on \bar{D} , that is, for every function u which satisfies the heat equation on a neighborhood of \bar{D} .

An interesting example of such a density is the following function K on $\Omega(c)$:

$$(2) \quad K(x, t) := \frac{1}{2^{n+2}(\pi c)^{n/2}} \frac{\|x\|^2}{t^2}$$

(see [5]). Here $\Omega(c)$ is the heat ball defined by a level surface of the Gauss–Weierstrass kernel W , that is,

$$\Omega(c) := \{(x, t) \in \mathbb{R}^{n+1}; W(x, -t) > (4\pi c)^{-n/2}\}$$

with

$$W(x, t) := \begin{cases} (4\pi t)^{-n/2} \exp(-\|x\|^2/4t) & \text{if } t > 0, \\ 0 & \text{if } t \leq 0, \end{cases}$$

and $\|x\| = (x_1^2 + \dots + x_n^2)^{1/2}$.

In this paper, we consider the following problems:

- (i) Which domains have a mean value density?
- (ii) Which domains have a bounded mean value density?

(iii) Does there exist a mean value density that is bounded away from zero?

For the harmonic case, similar problems were discussed by Hansen and Netuka in [2]. They showed that, for every bounded domain U in \mathbb{R}^n that contains 0, there exists a bounded function $K > 0$ on U such that

$$h(0) = \int_U K(x)h(x) dx$$

for every bounded harmonic function h on U . Furthermore, for smooth domains they constructed such functions K with $\inf_{x \in U} K(x) > 0$. In our parabolic case the situation is more complicated. It is easily seen that if

$$(3) \quad \sup\{t; (x, t) \in D\} > 0$$

then D does not have a mean value density. Furthermore, there is no mean value density on a cone $\{\|x\| < -ct; -1 < t < 0\}$ (see Corollary 7(a) below). On the other hand, every rectangle $\{(x, t); |x_i| < c \text{ for all } i, -c^2 < t < 0\}$ has a bounded mean value density (see [1, p. 276]). A heat ball has a mean value density as above, but we shall see later that there is no bounded density there. Another example of a domain that has bounded mean value density is a modified heat ball, defined in [6]. Bounded mean value densities are useful for the monotone approximation of subtemperatures by smooth subtemperatures.

In Section 2 we construct mean value densities for certain domains. The argument is based on that in [2], but considerable modification of the details is necessary. In Section 3, we discuss the above problems (i)–(iii) for special domains.

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2. For a domain D in \mathbb{R}^{n+1} , we denote by $\partial_p D$ the parabolic boundary of D , that is, the set of boundary points which can be connected to some point of D by a curve in D having strictly increasing t -coordinate. Also for $(x_0, t_0) \in D$, $\Lambda(x_0, t_0; D)$ is the set of all points $(x, t) \in D \setminus \{(x_0, t_0)\}$ which can be connected to (x_0, t_0) by a polygonal line in D having strictly increasing t -coordinate. We write $\Omega(y, s; c)$ for the heat ball with centre (y, s) and radius $c > 0$, that is,

$$(4) \quad \Omega(y, s; c) := \{(x, t) \in \mathbb{R}^{n+1}; W(y - x, s - t) > (4\pi c)^{-n/2}\}.$$

Hence $\Omega(c) = \Omega(0, 0; c)$. Further, for $a > 0$ we put

$$(5) \quad K_a(x, t) := \frac{\|x\|^2}{(-t)^{(n+4-2a)/2}} \exp\left(\frac{(2a-n)\|x\|^2}{4n(-t)}\right) \quad (t < 0)$$

and define the constant $p(a, c)$ by

$$(6) \quad p(a, c) := \frac{a}{2^{n+1}n\pi^{n/2}c^a}.$$

Note that $p(n/2, c)K_{n/2}$ is the function K in (2). In view of [5], the functions $p(a, c)K_a$ are also mean value densities on $\Omega(c)$.

Regarding the existence of mean value densities, we have the following result.

THEOREM 1. *Let D be a bounded domain in \mathbb{R}^{n+1} such that $\Omega(c_0) \subset D$ for some $c_0 > 0$. Suppose that there exists a family $\{E_\alpha\}_{\alpha \in A}$ of subdomains satisfying the following conditions:*

- (a) *For each $\alpha \in A$, $\Omega(c_0/2) \subset E_\alpha \subset D$, $\overline{E_\alpha}^\circ = E_\alpha$, and for every $(y, s) \in E_\alpha$ there exists $(z, r) \in \Omega(c_0/2)$ such that $(y, s) \in \Lambda(z, r; E_\alpha)$.*
- (b) $\bigcup_{\alpha \in A} \overline{\partial_p E_\alpha} \supset D \setminus \Omega(2c_0/3)$.

Then there is a mean value density on D .

Proof. Fix a nonnegative, continuous function η on $[0, \infty)$ such that $\{t; \eta(t) > 0\} = [0, 1)$ and

$$\int_0^1 (4\pi t)^{n/2} \eta(t) dt = 1.$$

For each $(y, s) \in D$, put

$$\gamma(y, s) := \frac{1}{2} \sup\{c; \Omega(y, s; c) \subset D\},$$

and define

$$\begin{aligned} \tau_{(y,s)}(x, t) &:= \frac{1}{2n} K_{(n+2)/2}(y-x, s-t) \\ &\quad \times \gamma(y, s)^{-(n+2)/2} \eta\left(\frac{s-t}{\gamma(y, s)} \exp\left(\frac{\|y-x\|^2}{2n(s-t)}\right)\right) \end{aligned}$$

whenever $t < s$, and $\tau_{(y,s)}(x, t) := 0$ whenever $t \geq s$. Then $\tau_{(y,s)}$ is continuous on $\mathbb{R}^n \times (-\infty, s)$, and

$$\begin{aligned} &\{(x, t); \tau_{(y,s)}(x, t) > 0\} \\ &= \{(x, t); 0 < (s-t) \exp(\|y-x\|^2/(2n(s-t))) < \gamma(y, s), x \neq y\} \\ &= \Omega(y, s; \gamma(y, s)) \setminus (\{y\} \times \mathbb{R}). \end{aligned}$$

For every $\alpha \in A$, let $\mu_\alpha^{(z,r)}$ denote the parabolic measure at (z, r) for E_α , and put

$$w_\alpha(x, t) := p(n/2, c_0/2) \int_{\Omega(c_0/2)} \left(\int_{\partial E_\alpha} \tau_{(y,s)}(x, t) d\mu_\alpha^{(z,r)}(y, s) \right) \frac{\|z\|^2}{r^2} dz dr$$

for every $(x, t) \in \mathbb{R}^{n+1}$. By the minimum principle for temperatures, if $(y, s) \in \Lambda(z, r; E_\alpha)$ then $\mu_\alpha^{(y,s)}$ is absolutely continuous with respect to $\mu_\alpha^{(z,r)}$, so that condition (a) implies that

$$\bigcup \{\text{supp}(\mu_\alpha^{(y,s)}); (y, s) \in E_\alpha\} = \bigcup \{\text{supp}(\mu_\alpha^{(z,r)}); (z, r) \in \Omega(c_0/2)\}.$$

Hence, by [4, Theorem 1],

$$\overline{\partial_p E_\alpha} = \overline{\bigcup \{\text{supp}(\mu_\alpha^{(z,r)}); (z, r) \in \Omega(c_0/2)\}}.$$

We claim that

$$(7) \quad \{(x, t); w_\alpha(x, t) > 0\} = \bigcup \{\Omega(y, s; \gamma(y, s)) \setminus (\{y\} \times \mathbb{R}); (y, s) \in \overline{\partial_p E_\alpha}\}.$$

To prove (7), we first show that $w_\alpha(x, t) > 0$ if $(x, t) \in \Omega(y_0, s_0; \gamma(y_0, s_0)) \setminus (\{y_0\} \times \mathbb{R})$ for some $(y_0, s_0) \in \overline{\partial_p E_\alpha}$. Since $\tau_{(y_0, s_0)}(x, t) > 0$, there is an open neighbourhood B of (y_0, s_0) such that $\tau_{(y,s)}(x, t) > 0$ for all $(y, s) \in B$. In particular $\tau_{(y,s)}(x, t) > 0$ on $B \cap \partial E_\alpha$.

We consider two cases. First suppose that

$$\mu_\alpha^{(z,r)}(B \cap \partial E_\alpha) = 0$$

for all $(z, r) \in \Omega(c_0/2)$. Then $\text{supp}(\mu_\alpha^{(z,r)})$ is contained in $\partial E_\alpha \setminus B$, so that since $B \cap \partial E_\alpha$ is open in ∂E_α ,

$$(y_0, s_0) \in \overline{\partial_p E_\alpha} = \overline{\bigcup \{\text{supp}(\mu_\alpha^{(z,r)}); (z, r) \in \Omega(c_0/2)\}} \subset \partial E_\alpha \setminus B.$$

This contradicts the fact that $(y_0, s_0) \in B$.

Second, suppose that

$$\mu_\alpha^{(z_0, r_0)}(B \cap \partial E_\alpha) > 0$$

for some $(z_0, r_0) \in \Omega(c_0/2)$. Then, by the minimum principle,

$$\mu_\alpha^{(z,r)}(B \cap \partial E_\alpha) > 0$$

for all $(z, r) \in \Omega(c_0/2)$ with $r > r_0$. Since $\tau_{(y,s)}(x, t) > 0$ for $(y, s) \in B \cap \partial E_\alpha$, we have

$$\int_{\partial E_\alpha} \tau_{(y,s)}(x, t) d\mu_\alpha^{(z,r)}(y, s) > 0$$

for all $(z, r) \in \Omega(c_0/2)$ with $r > r_0$. This implies that $w_\alpha(x, t) > 0$.

Conversely, if $w_\alpha(x, t) > 0$, then

$$\int_{\partial E_\alpha} \tau_{(y,s)}(x, t) d\mu_\alpha^{(z_0, r_0)}(y, s) > 0$$

for some (z_0, r_0) , and hence

$$\tau_{(y_0, s_0)}(x, t) > 0$$

for some $(y_0, s_0) \in \text{supp}(\mu_\alpha^{(z_0, r_0)})$. Thus $(x, t) \in \Omega(y_0, s_0; \gamma(y_0, s_0)) \setminus (\{y_0\} \times \mathbb{R})$. Since

$$\text{supp}(\mu_\alpha^{(z_0, r_0)}) \subset \overline{\partial_p E_\alpha}$$

we have $(y_0, s_0) \in \overline{\partial_p E_\alpha}$. Thus (7) is established.

Next, for every temperature u on \overline{D} , we have

$$\begin{aligned} & \iint u(x, t) \tau_{(y, s)}(x, t) \, dx \, dt \\ &= \iint_{\Omega(y, s; \gamma(y, s))} u(x, t) K_{(n+2)/2}(y - x, s - t) \\ & \quad \times \gamma(y, s)^{-(n+2)/2} \eta\left(\frac{s - t}{\gamma(y, s)} \exp\left(\frac{\|y - x\|^2}{2n(s - t)}\right)\right) \, dx \, dt \\ &= \int_0^{\gamma(y, s)} d\ell \int_{\partial\Omega(y, s; \ell)} Q(y - \xi, s - \tau) u(\xi, \tau) \\ & \quad \times \gamma(y, s)^{-(n+2)/2} \eta\left(\frac{\ell}{\gamma(y, s)}\right) \, d\sigma(\xi, \tau) \\ &= \int_0^{\gamma(y, s)} \gamma(y, s)^{-(n+2)/2} \eta\left(\frac{\ell}{\gamma(y, s)}\right) (4\pi\ell)^{n/2} u(y, s) \, d\ell \\ &= u(y, s) \int_0^1 \gamma(y, s)^{-(n+2)/2} \eta(t) (4\pi t \gamma(y, s))^{n/2} \gamma(y, s) \, dt \\ &= u(y, s) \int_0^1 (4\pi t)^{n/2} \eta(t) \, dt = u(y, s), \end{aligned}$$

because $(s - \tau) \exp(\|y - \xi\|^2/2n(s - \tau)) = \ell$ on $\partial\Omega(y, s; \ell)$; here $Q(x, t) = \|x\|^2(4\|x\|^2 t^2 + (\|x\|^2 - 2nt)^2)^{-1/2}$ and σ is the surface area measure on $\partial\Omega((y, s); \ell)$ (see [6]). Hence

$$\begin{aligned} & \iint_D u(x, t) w_\alpha(x, t) \, dx \, dt \\ &= p(n/2, c_0/2) \\ & \quad \times \int_{\Omega(c_0/2)} \left(\int_{\partial E_\alpha} \left(\iint_D u(x, t) \tau_{(y, s)}(x, t) \, dx \, dt \right) d\mu_\alpha^{(z, r)}(y, s) \right) \frac{\|z\|^2}{r^2} \, dz \, dr \\ &= p(n/2, c_0/2) \int_{\Omega(c_0/2)} \left(\int_{\partial E_\alpha} u(y, s) \, d\mu_\alpha^{(z, r)}(y, s) \right) \frac{\|z\|^2}{r^2} \, dz \, dr \end{aligned}$$

$$= p(n/2, c_0/2) \int_{\Omega(c_0/2)} u(z, r) \frac{\|z\|^2}{r^2} dz dr = u(0, 0).$$

If $(x, t) \in D \setminus \Omega(c_0)$, then there is $(y, s) \in D \setminus \Omega(2c_0/3)$ such that $(x, t) \in \Omega(y, s; \gamma(y, s)) \setminus (\{y\} \times \mathbb{R})$. By condition (b), there is α such that $(y, s) \in \overline{\partial_p E_\alpha}$. So, by (7), $w_\alpha(x, t) > 0$ and $\{w_\alpha > 0\}$ is open. Thus the sets $\{w_\alpha > 0\}_{\alpha \in A}$ form an open cover for $D \setminus \Omega(c_0)$, so that the Lindelöf property ensures that we can choose a countable subcover $\{w_{\alpha_k} > 0\}_{k=1}^\infty$. Put

$$K(x, t) := \frac{p(n/2, c_0)}{2} \frac{\|x\|^2}{t^2} \chi_{\Omega(c_0)}(x, t) + \sum_{k=1}^\infty 2^{-k-1} w_{\alpha_k}(x, t),$$

where $\chi_{\Omega(c_0)}$ is the characteristic function of $\Omega(c_0)$. Then $K > 0$ a.e. on D . Also, for every temperature u on \overline{D} , we have

$$\begin{aligned} \iint_D u(x, t) K(x, t) dx dt &= \frac{p(n/2, c_0)}{2} \iint_{\Omega(c_0)} u(x, t) \frac{\|x\|^2}{t^2} dx dt \\ &\quad + \sum_{k=1}^\infty 2^{-k-1} \iint_D u(x, t) w_{\alpha_k}(x, t) dx dt \\ &= \frac{u(0, 0)}{2} + \sum_{k=1}^\infty 2^{-k-1} u(0, 0) = u(0, 0). \end{aligned}$$

This completes the proof of Theorem 1.

The class of domains which have bounded mean value densities is more restricted. The closure of such a domain contains every truncated heat ball, as we now show.

THEOREM 2. *Assume that there is a bounded mean value density K on a domain D . Then for every $c > 0$, there exists $t_c < 0$ such that*

$$(8) \quad \overline{D} \supset \Omega(c) \cap \{t > t_c\}.$$

Proof. Consider the function

$$(9) \quad v(y, s) := \iint_D K(x, t) W(x - y, t - s) dx dt.$$

Suppose that the assertion does not hold for some $c > 0$. Then we can choose points $\{(y_k, s_k)\}$ in $\Omega(c) \setminus \overline{D}$ such that $s_k > -1/k$ for all $k \geq 1$. Note that $(y_k, s_k) \rightarrow (0, 0)$ as $k \rightarrow \infty$. Since $(y_k, s_k) \notin \overline{D}$, we have

$$(10) \quad W(y_k, -s_k) = v(y_k, s_k)$$

by (1). Then $\liminf_{k \rightarrow \infty} W(y_k, -s_k) \geq (4\pi c)^{-n/2}$ because $(y_k, s_k) \in \Omega(c)$. On the other hand, the right hand side of (10) tends to zero as $k \rightarrow \infty$,

because the boundedness of K ensures that v is continuous on \mathbb{R}^{n+1} and $v(0,0) = 0$. This is a contradiction.

3. In this section we discuss mean value densities on domains of the form

$$D(\varphi) := \{(x, t) \in \mathbb{R}^{n+1}; \|x\| < \varphi(t), -1 < t < 0\},$$

where φ is a continuous function on $[-1, 0]$ with $\varphi > 0$ on $(-1, 0)$. For simplicity, we also assume that

- (*) there is $t_0 \in [-1, 0]$ such that φ is strictly decreasing on $[t_0, 0]$ and strictly increasing on $[-1, t_0]$.

The following remark will be useful below.

REMARK 3. If $D(\varphi)$ has a mean value density K , then whenever $(y, s) \notin \overline{D(\varphi)}$,

$$(11) \quad W(y, -s) = \iint_{D(\varphi)} K(x, t) W(x - y, t - s) dx dt.$$

Hence letting $(y, s) \rightarrow (y_0, s_0) \in \partial D(\varphi)$, we deduce from Fatou's lemma that

$$(12) \quad W(y_0, -s_0) \geq \iint_{D(\varphi)} K(x, t) W(x - y_0, t - s_0) dx dt.$$

Regarding the nonexistence of mean value densities, we have the following result.

THEOREM 4. *If the origin is a regular boundary point of $D(\varphi)$ with respect to the Dirichlet problem for the heat equation, then there is no mean value density on $D(\varphi)$.*

Proof. Under this hypothesis $t_0 < 0$. Let f be a continuous function on $\partial D(\varphi)$ such that $f(0,0) = 0$, $f(x, t) > 0$ if $t > t_0$, and $f(x, t) = 0$ if $t \leq t_0$. Let v be the solution of the Dirichlet problem on $D(\varphi)$ with boundary function f , and $v = f$ on $\partial D(\varphi)$. Then $v \geq 0$ and $v \not\equiv 0$. For $k \in \mathbb{N}$ such that $-1/k > t_0$, put

$$u_k(x, t) := \begin{cases} v(x, t - 1/k) & \text{if } t > t_0 + 1/k, \\ 0 & \text{if } t \leq t_0 + 1/k. \end{cases}$$

Now suppose that there is a mean value density K on $D(\varphi)$. Since u_k is a temperature on $\overline{D(\varphi)}$, we have

$$u_k(0, 0) = \iint_{D(\varphi)} K(x, t) u_k(x, t) dx dt.$$

Since $(0, 0)$ is regular, we have

$$\lim_{k \rightarrow \infty} u_k(0, 0) = \lim_{k \rightarrow \infty} v(0, -1/k) = f(0, 0) = 0.$$

On the other hand, Fatou's lemma implies that

$$\liminf_{k \rightarrow \infty} \iint_{D(\varphi)} K(x, t) u_k(x, t) dx dt \geq \iint_{D(\varphi)} K(x, t) v(x, t) dx dt > 0.$$

This is a contradiction.

REMARK 5. It is known that if φ satisfies

$$(13) \quad \varphi(t) < (4(-t) \log |\log(-t)|)^{1/2}$$

on a neighborhood of $t = 0$, then the origin is a regular boundary point (see [1, p. 339]). On the other hand, $(0, 0)$ is an irregular boundary point of $\Omega(c)$ (see [1, p. 340]). Let $m \geq 3$ be an integer. A modified heat ball $\Omega_m(c)$ is defined by

$$\Omega_m(c) := \{(x, t) \in \mathbb{R}^{n+1}; \|x\| < (2(m+n)(-t) \log(c/(-t)))^{1/2}, -c < t < 0\}.$$

The function

$$(14) \quad K(x, t) := c_0(2(m+n)(-t) \log(c/(-t)) - \|x\|^2)^{m/2} \\ \times \left(\frac{m(m+n)}{-t} \log(c/(-t)) + \frac{\|x\|^2}{t^2} \right)$$

is a bounded mean value density on $\Omega_m(c)$ (see [6]), where

$$c_0 := \frac{\omega_m}{2(m+2)(4\pi c)^{(m+n)/2}}$$

and ω_m is the volume of the unit ball in \mathbb{R}^m .

REMARK 6. In general, the regularity of the origin is not a sufficient condition for nonexistence of a mean value density. In fact, given an integer $m \geq 3$, let

$$D := \Omega_m(c) \setminus \{(x_1, \dots, x_{n-1}, 0, t); x_i \in \mathbb{R} \ (i = 1, \dots, n-1), -c/2 < t < 0\}.$$

Then $(0, 0)$ is a regular point of ∂D (see [3, p. 218] for the case $n = 1$); but (14) is a bounded mean value density for D , because $\overline{D}^o = \Omega_m(c)$. This example also shows that we cannot replace \overline{D} by D in (8).

For two functions φ and ψ on $[-1, 0]$, we write $\varphi \approx \psi$ if there exist positive constants c_1, c_2 such that $c_1\psi(t) \leq \varphi(t) \leq c_2\psi(t)$ on a neighbourhood of $t = 0$. We have the following results about the domains $D(\varphi)$.

COROLLARY 7. (a) *If $\varphi(t) \approx (-t)^\beta$ with $\beta \geq 1/2$, then there is no mean value density on $D(\varphi)$.*

(b) *If $\varphi(t) \approx (-t)^\beta$ with $\beta < 1/2$, then we can construct a mean value density on $D(\varphi)$.*

(c) *If there are $c_1 > 0$ and $t_1 < 0$ such that $(D(\varphi) \setminus \Omega(c_1)) \cap \{t > t_1\} = \emptyset$, then $D(\varphi)$ does not have a bounded mean value density. In particular, there is no bounded mean value density on a heat ball.*

(d) If $\varphi(t) \approx (-t)^\beta$, then $D(\varphi)$ has no mean value density that is bounded away from zero.

Proof. Part (a) follows from Theorem 4. To prove (b), we use Theorem 1. Choose $c_0 > 0$ such that $D(\varphi) \supset \Omega(c_0)$. For $0 < \alpha < 1$, put

$$E_\alpha := \{(x, t); \|x\| < \alpha\varphi(t), -1 < t < 0\} \cup \Omega(c_0/2).$$

Then $\{E_\alpha\}_{0 < \alpha < 1}$ satisfies the conditions of Theorem 1.

Part (c) follows from Theorem 2.

To show (d), we use the following assertion: There is a positive integer N , which depends only on the dimension n , and points $\{x_i\}_{i=1}^N$ in the unit sphere of \mathbb{R}^n such that for every $r > 0$,

$$(15) \quad B(0, r/2) \setminus \{0\} \subset \bigcup_{i=1}^N B(rx_i, r),$$

where $B(x, r)$ is the usual ball in \mathbb{R}^n with centre x and radius $r > 0$. The existence of N and $\{x_i\}_{i=1}^N$ is not difficult, because $B(0, 1/2) \cap B(x_1, 1)$ contains a truncated cone at the origin which has a positive aperture.

Now suppose that there is a mean value density K such that $K \geq c_0 > 0$ on $D(\varphi)$. Then $\beta < 1/2$. For $\{x_i\}_{i=1}^N$ as above and $-1 < s < 0$, we put

$$u_s(x, t) := \sum_{i=1}^N W(x - x_i\varphi(s), t - s).$$

Then u_s is a nonnegative temperature on $D(\varphi)$, and by (12) we have

$$(16) \quad u_s(0, 0) \geq \iint_{D(\varphi)} K(x, t)u_s(x, t) dx dt \geq c_0 \iint_{D(\varphi)} u_s(x, t) dx dt.$$

Note that

$$u_s(0, 0) = \sum_{i=1}^N W(x_i\varphi(s), -s) = \frac{N}{(4\pi(-s))^{n/2}} \exp\left(-\frac{\varphi(s)^2}{-4s}\right).$$

On the other hand, since

$$B(0, \varphi(s)/2) \setminus \{0\} \subset \bigcup_{i=1}^N B(x_i\varphi(s), \varphi(s))$$

by (15), we see that

$$\begin{aligned} \iint_{D(\varphi)} u_s(x, t) dx dt &= \int_s^0 \left(\int_{\|x\| < \varphi(s)} \sum_{i=1}^N W(x - x_i\varphi(s), t - s) dx \right) dt \\ &\geq \int_s^0 \left(\int_{\|x\| < \varphi(s)/2} W(x, t - s) dx \right) dt \end{aligned}$$

$$\begin{aligned}
&= n\omega_n \int_s^0 \left(\int_0^{\varphi(s)/2} \frac{1}{(4\pi(t-s))^{n/2}} \exp\left(-\frac{r^2}{4(t-s)}\right) r^{n-1} dr \right) dt \\
&= n\omega_n \pi^{-n/2} \int_s^0 \left(\int_0^{\varphi(s)/4\sqrt{t-s}} \tau^{n-1} \exp(-\tau^2) d\tau \right) dt \geq c_3(-s)
\end{aligned}$$

for all sufficiently small s , because $\varphi(s) \geq c_1(-s)^\beta$ and $\beta < 1/2$. Thus (16) implies that

$$\frac{N}{(4\pi(-s))^{n/2}} \exp(-c_1(-s)^{2\beta-1}) \geq \frac{N}{(4\pi(-s))^{n/2}} \exp\left(-\frac{\varphi(s)^2}{-s}\right) \geq c_0 c_3(-s).$$

This is a contradiction.

REMARK 8. We conjecture that there is no domain which has a mean value density bounded away from zero. Assertion (d) in Corollary 7 supports our conjecture.

REFERENCES

- [1] J. L. Doob, *Classical Potential Theory and its Probabilistic Counterpart*, Springer, 1984.
- [2] W. Hansen and I. Netuka, *Volume densities with the mean value property for harmonic functions*, Proc. Amer. Math. Soc. 123 (1995), 135–140.
- [3] R. Kaufman and J.-M. Wu, *Parabolic potential theory*, J. Differential Equations 43 (1982), 204–234.
- [4] N. Suzuki, *On the essential boundary and supports of harmonic measures for the heat equation*, Proc. Japan Acad. Ser. A Math. Sci. 56 (1980), 381–385.
- [5] N. A. Watson, *Volume mean values of subtemperatures*, Colloq. Math. 86 (2000), 253–258.
- [6] —, *Elementary proofs of some basic subtemperature theorems*, ibid. 94 (2002), 111–140.

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