# C OLLOQUIUM MATHEMATICUM 

# ON SOME CLASS OF NEARLY CONFORMALLY SYMMETRIC MANIFOLDS 

BY

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1. Introduction. Let $(M, g)$ and $(\bar{M}, \bar{g})$ be two Riemannian or pseudoRiemannian manifolds of class $C^{\infty}$. A mapping $\gamma:(M, g) \rightarrow(\bar{M}, \bar{g})$ is said to be geodesic if it preserves geodesics, i.e. maps geodesics of $(M, g)$ onto geodesics of $(\bar{M}, \bar{g})$. The metrics $g$ and $\bar{g}$ are then said to be geodesically corresponding.

Suppose that both $g$ and $\bar{g}$ are metrics on the same manifold $M$. Let $\mathfrak{F}(M)$ be the ring of differentiable functions and $\mathfrak{X}(M)$ the $\mathfrak{F}$-module of differentiable vector fields on $M$. Each of the conditions below is necessary and sufficient for the metrics $g$ and $\bar{g}$ to be geodesically corresponding:

$$
\begin{gather*}
\bar{\nabla}_{X} Y=\nabla_{X} Y+(X \psi) Y+(Y \psi) X  \tag{1}\\
\left(\nabla_{X} \bar{g}\right)(Y, Z)=2(X \psi) \bar{g}(Y, Z)+(Y \psi) \bar{g}(X, Z)+(Z \psi) \bar{g}(X, Y) \tag{2}
\end{gather*}
$$

for all $X, Y, Z \in \mathfrak{X}(M)$, where $\psi \in \mathfrak{F}(M)$, and $\nabla$ and $\bar{\nabla}$ are the Levi-Civita connections with respect to $g$ and $\bar{g}$.

A manifold $(M, g)$ admits a geodesic mapping if and only if there exist a function $\varphi \in \mathfrak{F}(M)$ and a symmetric non-singular bilinear form $a$ on $M$ satisfying

$$
\begin{equation*}
\left(\nabla_{X} a\right)(Y, Z)=(Y \varphi) g(X, Z)+(Z \varphi) g(X, Y) \tag{3}
\end{equation*}
$$

for all $X, Y, Z \in \mathfrak{X}(M)([9])$.
In a chart $(U, x)$ on $M$ the local components of $g, \bar{g}, a, X \varphi$ and $X \psi$ given by

$$
\begin{gathered}
g_{i j}=g\left(X_{i}, X_{j}\right), \quad \bar{g}_{i j}=\bar{g}\left(X_{i}, X_{j}\right), \quad a_{i j}=a\left(X_{i}, X_{j}\right), \\
\varphi_{i}=X_{i} \varphi, \quad \psi_{i}=X_{i} \psi
\end{gathered}
$$

satisfy

$$
\begin{equation*}
a_{i j}=\exp (2 \psi) \bar{g}^{s t} g_{s i} g_{t j}, \tag{4}
\end{equation*}
$$

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$$
\begin{equation*}
\varphi_{i}=-\exp (2 \psi) \bar{g}^{s t} g_{s i} \psi_{t}, \tag{5}
\end{equation*}
$$

where $X_{i}=\partial / \partial x^{i} \in \mathfrak{X}(U)$ and the $\bar{g}^{i j}$ are the components of $\left(\bar{g}_{i j}\right)^{-1}$.
A geodesic mapping is said to be non-trivial if it is non-affine, which is equivalent to $\varphi \neq$ const on $M$.

By (1), the curvature tensors and Ricci tensors of $(M, g)$ and $(M, \bar{g})$ are related by

$$
\begin{align*}
\bar{R}(X, Y) Z & =R(X, Y) Z+P(X, Z) Y-P(Y, Z) X  \tag{6}\\
\bar{S}(X, Y) & =S(X, Y)+(n-1) P(X, Y) \tag{7}
\end{align*}
$$

where

$$
\begin{equation*}
P(X, Y)=H_{\psi}(X, Y)-(X \psi)(Y \psi) \tag{8}
\end{equation*}
$$

and $H_{\psi}$ is the Hessian of $\psi$.
Following W. Roter ([7]), the manifold $(M, g)$ is said to be nearly conformally symmetric if the tensor

$$
L(X, Y)=\frac{1}{n-2}\left[S(X, Y)-\frac{r}{2(n-1)} g(X, Y)\right]
$$

is a Codazzi tensor, where $S$ is the Ricci tensor and $r$ denotes the scalar curvature.
N. S. Sinyukov ([9)] and E. N. Sinyukova ([10]) investigated manifolds whose Ricci tensor satisfies

$$
\begin{equation*}
\left(\nabla_{X} S\right)(Y, Z)=\sigma(X) g(Y, Z)+\nu(Y) g(X, Z)+\nu(Z) g(X, Y) \tag{9}
\end{equation*}
$$

where $\sigma$ and $\nu$ are some 1 -forms. Such manifolds are known under different names (see [9], [10], [3]). In what follows a Riemannian (or pseudoRiemannian) manifold satisfying (9) with non-constant scalar curvature will be called a Sinyukov manifold. Such manifolds always admit non-trivial geodesic mappings and every Sinyukov manifold is nearly conformally symmetric (see Lemma 1).

Let $(M, g)$ admit a non-trivial geodesic mapping onto the manifold $(M, \bar{g})$ defined by the 1-form $d \psi$ (see (2)). In [5] it was proved that $(M, g)$ ( $\operatorname{dim} M$ $\geq 3)$ is a conformally flat Sinyukov manifold if and only if $(M, \widetilde{g}=\exp (2 \psi) g)$ is of constant sectional curvature. In the present paper we prove that $(M, g)$ with nowhere vanishing Weyl conformal curvature tensor is a Sinyukov manifold if and only if $(M, \widetilde{g}=\exp (2 \psi) g)$ is either an Einstein manifold admitting non-trivial geodesic mappings or a Sinyukov manifold.

In [3] some properties of Sinyukov manifolds with non-null vector $\Phi$ defined by (24) below were investigated. In the present paper we deal with manifolds without any assumption on $\Phi$. Finally, the local structure theorem for Sinyukov manifolds is given.
2. Preliminaries. If $\widetilde{g}$ is a metric on $M$ and there exists $\lambda \in \mathfrak{F}(M)$ such that $\widetilde{g}=\exp (2 \lambda) g$, then $g$ and $\widetilde{g}$ are said to be conformally related or conformal to each other, and the transformation $g \rightarrow \widetilde{g}$ is called a conformal change. As is well-known, the Christoffel symbols, the curvature tensors and the Ricci tensors of the manifolds $(M, g)$ and $(M, \widetilde{g})$ are then related by

$$
\begin{equation*}
\widetilde{\nabla}_{X} Y=\nabla_{X} Y+(X \lambda) Y+(Y \lambda) X-g(X, Y) \Lambda \tag{10}
\end{equation*}
$$

where the vector field $\Lambda$ is defined by $g(X, \Lambda)=X \lambda$ for $X \in \mathfrak{X}(M)$;

$$
\begin{align*}
g(\widetilde{R}(X, Y) Z, V)= & g(R(X, Y) Z, V)  \tag{11}\\
& +Q(X, Z) g(Y, V)-Q(Y, Z) g(X, V) \\
& +g(X, Z) Q(Y, V)-g(Y, Z) Q(X, V)
\end{align*}
$$

for arbitrary $X, Y, Z, V \in \mathfrak{X}(M)$;

$$
\begin{equation*}
Q(X, Y)+L(X, Y)=\widetilde{L}(X, Y) \tag{12}
\end{equation*}
$$

where the tensor fields $Q$ and $L$ are given by

$$
\begin{align*}
Q(X, Y) & =H_{\lambda}(X, Y)-(X \lambda)(Y \lambda)+\frac{1}{2} g(\Lambda, \Lambda) g(X, Y)  \tag{13}\\
L(X, Y) & =\frac{1}{n-2}\left[S(X, Y)-\frac{r}{2(n-1)} g(X, Y)\right] \tag{14}
\end{align*}
$$

with $H_{\lambda}$ being the Hessian of $\lambda$, and $r$ standing for the scalar curvature of $(M, g)$. The tensor field $\widetilde{L}$ on $(M, \widetilde{g})$ is defined analogously.

The Weyl conformal curvature tensor $C$ satisfying
(15) $g(C(X, Y) Z, V)=g(R(X, Y) Z, V)+g(X, V) L(Y, Z)-g(Y, V) L(X, Z)$

$$
+L(X, V) g(Y, Z)-L(Y, V) g(X, Z)
$$

is invariant under conformal change, i.e. $\widetilde{C}=C$.
From (12) and (13) we get easily

$$
\begin{equation*}
g(C(X, Y) Z, \Lambda)=D(X, Y, Z)-\widetilde{D}(X, Y, Z) \tag{16}
\end{equation*}
$$

where

$$
\begin{equation*}
D(X, Y, Z)=\left(\nabla_{X} L\right)(Y, Z)-\left(\nabla_{Y} L\right)(X, Z) \tag{17}
\end{equation*}
$$

and the tensor field $\widetilde{D}$ on $(M, \widetilde{g})$ is defined in the same manner.
In the sequel we shall use the following theorem and lemmas.
Theorem 1 ([9]). If $(M, g)$ admits a non-trivial geodesic mapping onto a manifold $(M, \bar{g})$ defined by a 1 -form $d \psi$, then the manifold ( $M, a$ ), where a satisfies (3), admits a geodesic mapping onto $(M, \widetilde{g}=\exp (2 \psi) g)$ determined by the same 1 -form $d \psi$.

Lemma 1 ([12], [1]). On a Sinyukov manifold the tensor $D$ given by (17) vanishes, i.e. $L$ is a Codazzi tensor.

Lemma 2 ([9)]. If on $(M, g)$ relation (9) is satisfied at a point $p$, then

$$
\begin{equation*}
\sigma(X)=\frac{n}{(n-1)(n+2)}(X r), \quad \nu(X)=\frac{n-2}{2(n-1)(n+2)}(X r) \tag{18}
\end{equation*}
$$

for any $X \in T_{p}(M)$. Consequently, $(M, g)$ is a Sinyukov manifold if and only if the scalar curvature $r \neq$ const and the condition (9) holds everywhere on $M$.

We define $(1,1)$ tensor fields Ric and $A$ as follows:

$$
\begin{align*}
g(\operatorname{Ric}(X), Y) & =S(X, Y)  \tag{19}\\
g(A(X), Y) & =a(X, Y) \tag{20}
\end{align*}
$$

for all $X, Y \in \mathfrak{X}(M)$.
Lemma 3 ([10], [3]). If $(M, g)$ is a Sinyukov manifold and $d \varphi \neq 0$ at a point $p \in M$, then

$$
\begin{gather*}
a(\operatorname{Ric}(X), Y)=a(X, \operatorname{Ric}(Y)),  \tag{21}\\
a(X, N)-\frac{\operatorname{tr}(A)}{n} \nu(X)=S(X, \Phi)-\frac{r}{n}(X \varphi),  \tag{22}\\
(X \varphi)\left[S(Y, Z)-\frac{r}{n} g(Y, Z)\right]-(Y \varphi)\left[S(X, Z)-\frac{r}{n} g(X, Z)\right]  \tag{23}\\
=\nu(X)\left[a(Y, Z)-\frac{\operatorname{tr}(A)}{n} g(Y, Z)\right]-\nu(Y)\left[a(X, Z)-\frac{\operatorname{tr}(A)}{n} g(X, Z)\right]
\end{gather*}
$$

at $p$ for all $X, Y, Z \in T_{p}(M)$, where $N$ and $\Phi$ are given by

$$
\begin{equation*}
g(X, N)=\nu(X), \quad g(X, \Phi)=X \varphi \tag{24}
\end{equation*}
$$

3. Properties of conformal and geodesic mappings of Sinyukov manifolds. Let $p \in M$ be such that $d \varphi \neq 0$ and (3) hold at $p$. Choose a local coordinate system $(U, x)$ so that $p \in U$. By $R_{i j k}^{l}, S_{i j}, \varphi_{i j}$ we denote the components of the tensors $R, S$ and $H_{\varphi}$ in this coordinate system. Differentiating covariantly (3) and applying the Ricci identity we get

$$
\begin{equation*}
a_{i t} R_{j k l}^{t}+a_{t j} R_{i k l}^{t}=\varphi_{l i} g_{j k}+\varphi_{l j} g_{i k}-\varphi_{k i} g_{j l}-\varphi_{k j} g_{i l} \tag{25}
\end{equation*}
$$

Differentiating covariantly (25) with respect to $x^{m}$, contracting with $g^{l m}$ and applying the Ricci identity, by (3) and (9), we obtain
(26) $4 \varphi_{t} R^{t}{ }_{j k i}=\varphi_{k} S_{i j}-\varphi_{i} S_{j k}$

$$
+\frac{n+2}{n-2}\left[\nu^{t} a_{t k} g_{i j}-\nu^{t} a_{t i} g_{k j}+\nu_{i} a_{k j}-\nu_{k} a_{i j}\right]+b_{i} g_{j k}-b_{k} g_{i j}
$$

where the $\nu_{i}$ are the components of the 1 -form $\nu$, the $\nu^{i}$ are the components of the field $N$ (i.e. $\nu^{i}=g^{i t} \nu_{t}$ ) whereas $b_{i}=\varphi_{i t ; s} g^{t s}$ and the semicolon denotes covariant differentiation on $(M, g)$. Moreover, substituting (22) and
(23) into (26), we get

$$
\begin{align*}
4 \varphi_{t} R^{t}{ }_{j k i}= & \frac{4}{n-2}\left(\varphi_{i} S_{j k}-\varphi_{k} S_{i j}\right)  \tag{27}\\
& +\frac{n+2}{n-2}\left(\varphi^{t} S_{t k} g_{i j}-\varphi^{t} S_{t i} g_{k j}\right)+b_{i} g_{j k}-b_{k} g_{i j}
\end{align*}
$$

where $\varphi^{i}=\varphi_{t} g^{t i}$ are the components of the field $\Phi$.
Now, we shall prove
Proposition 1. If $(M, g)$ is a Sinyukov manifold and the Weyl conformal curvature tensor $C \neq 0$ and $d \varphi \neq 0$ at a point $p \in M$, then

$$
\begin{align*}
g(\Phi, C(X, Y) Z) & =0  \tag{28}\\
g(N, C(X, Y) Z) & =0 \tag{29}
\end{align*}
$$

on some neighbourhood $U_{1}$ of $p$, where $\Phi$ and $N$ are as in (24). So, for the metrics $\widetilde{g}_{1}=\exp (2 \varphi) g$ and $\widetilde{g}_{2}=\exp (2 \nu) g$, where $\nu \in \mathfrak{F}\left(U_{1}\right)$ and $X \nu=$ $\nu(X)$, the tensors $\widetilde{L}_{1}, \widetilde{L}_{2}$ defined by (14) are Codazzi tensors.

Proof. Transvecting (27) with $g^{j k}$ we get

$$
\begin{equation*}
b_{i}=\frac{n+6}{n-2} \varphi^{t} S_{t i}-\frac{4 r}{(n-1)(n-2)} \varphi_{i} \tag{30}
\end{equation*}
$$

Substituting (30) into (27), in view of (15), we find (28). Beginning with (9) and following the above argument we obtain (29). Thus the proposition is proved.

Proposition 2. If $(M, g)$ is a Sinyukov manifold, then on the set $U_{\varphi}=$ $\{p \in M: d \varphi \neq 0$ at $p\}$ the following identities hold:

$$
\begin{align*}
H_{\varphi}(X, Y)= & \frac{1}{n-2}\left[a(\operatorname{Ric}(X), Y)-\frac{r}{n(n-1)} a(X, Y)\right.  \tag{31}\\
& \left.-\frac{\operatorname{tr}(A)}{n} S(X, Y)+(n-2) \varrho_{1} g(X, Y)\right]+F(X \varphi)(Y \varphi), \\
H_{\psi}(X, Y)- & (X \psi)(Y \psi)  \tag{32}\\
= & -\frac{1}{n-2} S(X, Y)+\frac{r}{n(n-1)(n-2)} g(X, Y) \\
& -\frac{\operatorname{tr}(A)}{n(n-2)} \exp (-2 \psi) \bar{g}(\operatorname{Ric}(X), Y) \\
& +\bar{K}_{1} \bar{g}(X, Y)+F(X \psi)(Y \varphi), \\
\left(\nabla_{X} \nu\right)(Y)= & \frac{1}{n-2}\left[S(\operatorname{Ric}(X), Y)-\frac{r}{n-1} S(X, Y)\right.  \tag{33}\\
& \left.+(n-2) \varrho_{2} g(X, Y)\right]+G \nu(X) \nu(Y),
\end{align*}
$$

where $F, G \in \mathfrak{F}\left(U_{\varphi}\right), \bar{K}_{1}=-\varrho_{1} \exp (-2 \psi)-\exp (-2 \psi) g(\Phi, \Psi)$ and $\Psi$ is given by $g(X, \Psi)=X \psi$, and

$$
\begin{gather*}
\varrho_{1}=\frac{\Delta \varphi}{n}-\frac{1}{n(n-2)} g(\operatorname{Ric}(X), A(Y))+\frac{r \operatorname{tr}(A)}{n(n-1)(n-2)},  \tag{34}\\
\varrho_{2}=\frac{\Delta \nu}{n}-\frac{1}{n(n-2)}|S|^{2}+\frac{r^{2}}{n(n-1)(n-2)} \tag{35}
\end{gather*}
$$

with $\Delta \varphi, \Delta \nu$ standing for the traces of the $H e s s i a n ~ H_{\varphi}$ and $\nabla \nu$ with respect to $g$. If $\Phi$ or $N$ is non-null, then $F=0$ or $G=0$ respectively.

Proof. Transvecting (25) with $\varphi^{l}$ and applying (27) and (30) we obtain

$$
\begin{gather*}
\frac{1}{n-2}\left[\left(S_{k}^{t} a_{t i}-\frac{r}{n-1} a_{i k}\right) \varphi_{j}-\left(S_{i k}-\frac{r}{n-1} g_{i k}\right) a_{j t} \varphi^{t}\right.  \tag{36}\\
\quad+S_{j}^{t} \varphi_{t} a_{i k}-\varphi^{t} S_{t}^{s} a_{j s} g_{i k}+\left(S_{k}^{t} a_{t j}-\frac{r}{n-1} a_{j k}\right) \varphi_{i} \\
\left.\quad-\left(S_{j k}-\frac{r}{n-1} g_{j k}\right) a_{i t} \varphi^{t}+\varphi^{t} S_{t i} a_{j k}-\varphi^{t} S_{t}^{s} a_{i s} g_{j k}\right] \\
=\varphi_{k i} \varphi_{j}+\varphi_{k j} \varphi_{i}-\varphi_{t i} \varphi^{t} g_{j k}-\varphi_{t j} \varphi^{t} g_{i k}
\end{gather*}
$$

where $S_{i}^{j}=S_{i t} g^{t j}$. Transvecting (36) with $g^{j k}$ and making use of (21) we get

$$
\begin{align*}
\varphi_{i t} \varphi^{t}= & \varrho_{1} \varphi_{i}+\frac{1}{n-2} S_{i}^{t} a_{t s} \varphi^{s}-\frac{r}{n(n-1)(n-2)} a_{i t} \varphi^{t}  \tag{37}\\
& -\frac{\operatorname{tr}(A)}{n(n-2)} S_{i t} \varphi^{t}
\end{align*}
$$

where

$$
\varrho_{1}=\frac{\Delta \varphi}{n}-\frac{1}{n(n-2)} S^{t s} a_{t s}+\frac{r \operatorname{tr}(A)}{n(n-1)(n-2)}
$$

and $S^{i j}=S_{t}^{i} g^{t j}$. Substituting (37) into (36), in view of (22) and (23), we easily obtain (31). Hence, by metric contraction and the use of (34), we have either $F=0$, provided that $\Phi$ is non-null, or $F \neq 0$, provided that $\Phi$ is null. Moreover, (4) and (5) yield $\varphi_{i}=-a_{i t} \psi^{t}$, whence, by covariant differentiation and the use of (3) and (31), we get (32). Finally, beginning with (9) relations (33) and (35) can be obtained in a similar way to (31) and (34). This completes the proof.

Lemma 4. If $(M, g)$ is a Sinyukov manifold, then

$$
X \varphi=\omega \nu(X)
$$

on $U_{\varphi}$, where $\omega \in \mathfrak{F}\left(U_{\varphi}\right)$ and $X \in \mathfrak{X}\left(U_{\varphi}\right)$.

Proof. Consider the following two cases.
(i) The vector field $\Phi$ is null $($ see $(24))$. Since $\bar{g}(\operatorname{Ric}(X), Y)=\bar{g}(X, \operatorname{Ric}(Y))$ (cf. [8], p. 294), by (32), we get $X \varphi=-\tau(X \psi), \tau \in \mathfrak{F}\left(U_{\varphi}\right)$. It follows that if $\Phi$ is null, then so is $\Psi$. From (4) and (5) we have

$$
\begin{equation*}
a_{i t} \varphi^{t}=\tau \varphi_{i} . \tag{38}
\end{equation*}
$$

Moreover, (37) yields

$$
\begin{equation*}
\left(\tau-\frac{\operatorname{tr}(A)}{n}\right) S_{i t} \varphi^{t}=\left(\frac{r \tau}{n(n-1)}-(n-2) \varrho_{1}\right) \varphi_{i} \tag{39}
\end{equation*}
$$

In a local chart, (23) takes the form

$$
\begin{align*}
\varphi_{i}\left(S_{j k}-\frac{r}{n} g_{j k}\right)- & \varphi_{j}\left(S_{i k}-\frac{r}{n} g_{i k}\right)  \tag{40}\\
& =\nu_{i}\left(a_{j k}-\frac{\operatorname{tr}(A)}{n} g_{j k}\right)-\nu_{j}\left(a_{i k}-\frac{\operatorname{tr}(A)}{n} g_{i k}\right)
\end{align*}
$$

whence, transvecting with $\varphi^{k}$ and making use of (38) and (39), we get

$$
\begin{equation*}
S_{i t} \varphi^{t}=\tau_{1} \varphi_{i} \tag{41}
\end{equation*}
$$

where $\tau_{1} \in \mathfrak{F}\left(U_{\varphi}\right)$. Differentiating covariantly (41) and transvecting the resulting equation with $\varphi^{i}$ we find $\nu_{t} \varphi^{t}=0$. Finally, transvecting (40) with $\varphi^{j}$ we obtain

$$
\varphi_{i}\left(S_{t k} \varphi^{t}-\frac{r}{n} \varphi_{k}\right)=\nu_{i}\left(\tau-\frac{\operatorname{tr}(A)}{n}\right) \varphi_{k}
$$

Consider two cases.

1) $\tau=\operatorname{tr}(A) / n$. Then $S_{t k} \varphi^{t}=(r / n) \varphi_{k}$ at each point where $\varphi_{i} \neq 0$. Differentiating covariantly with respect to $x^{l}$ and alternating the resulting equation in $i, l$, in view of (9) and (18), we have $\varphi_{i} \nu_{l}=\varphi_{l} \nu_{i}$, and the result follows.
2) $\tau \neq \operatorname{tr}(A) / n$. Alternating the above result in $i, k$ and applying (41) we obtain the assertion.
(ii) The vector field $\Phi$ is non-null. Differentiating covariantly (22) and alternating the resulting equation, by (21), (18), (3), (31) and (33), we obtain
(42) $\frac{2(n+2)}{n}\left(\varphi_{i} \nu_{j}-\varphi_{j} \nu_{i}\right)+G\left(a_{i t} \nu^{t} \nu_{j}-a_{j t} \nu^{t} \nu_{i}\right)$

$$
=F\left(S_{i t} \varphi^{t} \varphi_{j}-S_{j t} \varphi^{t} \varphi_{i}\right)
$$

If $\Phi$ and $N$ are non-null, then the result follows from (42) and Proposition 2. Finally, let $N$ be a null vector field. Differentiating (33) covariantly, then applying the Ricci identity and comparing the resulting equation to (29), in
view of (9) and (18), we have

$$
\begin{equation*}
\varrho_{2 ; i}=\frac{2}{n-2} S_{i t} \nu^{t}+\varrho_{3} \nu_{i} \tag{43}
\end{equation*}
$$

where $\varrho_{3} \in \mathfrak{F}\left(U_{\varphi}\right)$. On the other hand, (33) gives

$$
S_{i}^{t} S_{t p} \nu^{p}-\frac{r}{n-1} S_{i t} \nu^{t}+(n-2) \varrho_{2} \nu_{i}=0 .
$$

Differentiating covariantly with respect to $x^{k}$, then transvecting with $\nu^{k}$, by the use of (43), we obtain $S_{t p} \nu^{t} \nu^{p}=0$. Hence, by transvection of (40) with $\nu^{j} \nu^{k}$, we have

$$
\begin{equation*}
-\varphi_{t} \nu^{t}\left(S_{i t} \nu^{t}-\frac{r}{n} \nu_{i}\right)=\nu_{i} a_{t p} \nu^{t} \nu^{p} \tag{44}
\end{equation*}
$$

Now, transvecting (40) with $\varphi^{i} \nu^{j}$ and applying the last result, we get $S_{k t} \nu^{t}=(r / n) \nu_{k}$ at each point where $\varphi_{t} \varphi^{t} \neq 0$ and $\varphi_{t} \nu^{t}=0$. Then transvection of (40) with $\nu^{k}$ results in $\nu_{i} a_{j t} \nu^{t}-\nu_{j} a_{i t} \nu^{t}=0$. Hence and from (42), by Proposition 2 we have $\varphi_{i}=\omega \nu_{i}$. On the other hand, if $\varphi_{t} \nu^{t} \neq 0$ in (44), then $S_{t i} \nu^{t}=\tau_{2} \nu_{i}, \tau_{2} \in \mathfrak{F}\left(U_{\varphi}\right)$. Therefore, transvecting (40) with $\varphi^{i} \nu^{j}$ and using (22), we obtain $a_{i t} \nu^{t}=\tau_{3} \nu_{i}, \tau_{3} \in \mathfrak{F}\left(U_{\varphi}\right)$, whence, by (42), we have $\varphi_{i}=\omega \nu_{i}$ again. From the above considerations it follows that the case when $\Phi$ is non-null but $N$ is null does not occur. This completes the proof.
4. Main results. From (23) and Lemma 4 it follows that

$$
\begin{equation*}
\omega\left(S_{i j}-\frac{r}{n} g_{i j}\right)=a_{i j}-\frac{\operatorname{tr}(A)}{n} g_{i j}+B \nu_{i} \nu_{j} \tag{45}
\end{equation*}
$$

in a local chart $(U, x)$, where $B \in \mathfrak{F}\left(U_{\varphi}\right)$. Transvecting (45) with $g^{i j}$ we get $B \nu_{t} \nu^{t}=0$. Hence and from Lemma 4 it follows that if the vector field $\Phi$ is non-null, then $B=0$. Now we shall prove

Proposition 3. Assuming that (3) and (9) are satisfied at a point $p \in$ $U_{\varphi}$ and $\Phi$ is a null vector field, we have $\nu(X)=0, X \in T_{p}(M)$, and $(U, g)$ is an Einstein manifold.

Proof. Suppose that the vector field $\Phi$ is null. Differentiating covariantly (45), then making use of (3), (9) and (18) we get

$$
\begin{equation*}
\omega_{k}\left(S_{i j}-\frac{r}{n} g_{i j}\right)=B_{k} \nu_{i} \nu_{j}+B\left(\nu_{i ; k} \nu_{j}+\nu_{i} \nu_{j ; k}\right) \tag{46}
\end{equation*}
$$

where $\omega_{k}=X_{k} \omega, B_{k}=X_{k} B$ and the semicolon stands for covariant differentiation on $(M, g)$. If $\omega=$ const, then $\nu_{i}=0$ is a consequence of the results of [10]. If $\omega_{k} \neq 0$ at $p$, then, by covariant differentiation of $\varphi_{i}=\omega \nu_{i}$, we obtain $X_{i} \omega=\omega_{1} \nu_{i}$ and $X_{i}\left(\omega_{1}\right)=\omega_{2} \nu_{i}$, where $\omega_{1}, \omega_{2} \in \mathfrak{F}(U)$. Moreover,
differentiating covariantly (46) and applying the Ricci identity, by (29), we find

$$
\begin{equation*}
\nu_{i} T_{j k l}+\nu_{j} T_{i k l}=0 \tag{47}
\end{equation*}
$$

where

$$
\begin{align*}
T_{j k l}= & {\left[\frac{B}{n-2} S_{j k}-\left(\frac{B r}{n(n-1)(n-2)}-\omega_{1}\right) g_{j k}\right] \nu_{l} }  \tag{48}\\
& -\left[\frac{B}{n-2} S_{j l}-\left(\frac{B r}{n(n-1)(n-2)}-\omega_{1}\right) g_{j l}\right] \nu_{k} .
\end{align*}
$$

If $\nu_{i} \neq 0$, then (47) results in $T_{j k l}=0$. Thus, by (48),

$$
\frac{B}{n-2} S_{i j}-\left(\frac{B r}{n(n-1)(n-2)}-\omega_{1}\right) g_{i j}=B_{1} \nu_{i} \nu_{j}, \quad B_{1} \in \mathfrak{F}(U)
$$

Hence, metric contraction with respect to $i, j$ gives

$$
\begin{equation*}
\omega_{1}=-\frac{B r}{n(n-1)} \tag{49}
\end{equation*}
$$

Therefore

$$
\begin{equation*}
S_{i j}-\frac{r}{n} g_{i j}=B_{2} \nu_{i} \nu_{j} \tag{50}
\end{equation*}
$$

where $B_{2}=(n-2) B_{1} / B$.
From (49) it follows that if $r \neq 0$, then $B_{i}=B_{3} \nu_{i}, B_{3} \in \mathfrak{F}(U)$. Substituting (50) into (46) and taking into account the above considerations, we obtain $\nu_{i} \nu_{j ; k}-\nu_{k} \nu_{j ; i}=0$ at each point where $B \neq 0$. Hence

$$
\begin{equation*}
\nu_{i ; j}=G_{1} \nu_{i} \nu_{j} \tag{51}
\end{equation*}
$$

where $G_{1} \in \mathfrak{F}(U)$. From (50) we obtain $\left(S_{i t}-(r / n) g_{i t}\right) \nu^{t}=0$, whence, by covariant differentiation and the use of (51) and (18), we have $\frac{n-2}{n} \nu_{i} \nu_{j}=0$. So, $\nu_{i}=0$ if $\Phi$ is a null vector field. Then (9) results in $\left(\nabla_{X} S\right)(Y, Z)=0$, which implies $\left(\nabla_{X} S\right)(Y, Z)-\left(\nabla_{Y} S\right)(X, Z)=0$ for $X, Y, Z \in T_{p}(M)$. Now the second part of our proposition is a consequence of the results of [11]. This completes the proof.

From Lemma 4, Proposition 3 and the results of [10] (cf. [3]) we obtain
ThEOREM 2. A manifold ( $M, g$ ) admitting a non-trivial geodesic mapping onto a pseudo-Riemannian manifold is a Sinyukov manifold if and only if both $r \neq$ const and the condition

$$
\begin{equation*}
a(X, Y)=\omega[S(X, Y)-(\sigma+c) g(X, Y)] \tag{52}
\end{equation*}
$$

holds everywhere, where $\omega=\mathrm{const} \neq 0, c=\mathrm{const}, \sigma \in \mathfrak{F}(M)$ and $X \sigma=$ $\sigma(X), X, Y \in \mathfrak{X}(M)$.

From Theorem 2 we obtain

Corollary 1. On a Sinyukov manifold $X \varphi=\omega \nu(X), \omega=\mathrm{const} \neq 0$.
Corollary 2 ([9]). A Sinyukov manifold $(M, g)$ always admits a nontrivial geodesic mapping onto a pseudo-Riemannian manifold.

Moreover, from Proposition 3 we have
Corollary 3. On a Sinyukov manifold the vector field $\Phi$ is non-null.
Now we shall prove
Proposition 4. Suppose that $(M, g)$ is a Sinyukov manifold and let $\bar{g}$ be a metric satisfying (2), i.e. $\bar{g}$ is geodesically corresponding to $g$. If $p \in U_{\varphi}$, then

$$
\begin{align*}
& H_{\psi}(X, Y)-(X \psi)(Y \psi)  \tag{53}\\
& \quad=-\frac{1}{(n-2) \omega} a(X, Y)+K g(X, Y)+\bar{K} \bar{g}(X, Y)
\end{align*}
$$

at $p$, where $X, Y \in T_{p}(M)$,

$$
K=-\frac{2}{n-2}(\sigma+c)+\frac{r}{(n-1)(n-2)}
$$

and

$$
\begin{equation*}
\bar{K}=\left[\frac{\operatorname{tr}(A)}{n(n-2)}(\sigma+c)-\varrho_{1}-g(\Psi, \Phi)\right] \exp (-2 \psi)=\text { const }, \tag{54}
\end{equation*}
$$

$H_{\psi}$ is the Hessian of the function $\psi, \varrho_{1}$ is given by (34) and $X(\sigma+c)=\sigma(X)$.
Proof. Equation (53) results immediately from (32) and (52). Differentiating covariantly (31) and applying the Ricci identity, by (3), (9), (52) and (28), we obtain

$$
X_{i}\left(\varrho_{1}\right)=\frac{2}{n-2} S_{i t} \varphi^{t}+\left[\frac{2 r}{(n-2)^{2}(n-1)}-\frac{2(n+2)}{n(n-2)^{2}}(\sigma+c)\right]\left(X_{i} \varphi\right)
$$

Then differentiating covariantly (54), by (3), (31), (32), (52) and the above identity, we easily find that $\bar{K}$ is constant on $U_{\varphi}$. Thus the proposition is proved.

Theorem 2 and Proposition 4 result in

## Corollary 4. On a Sinyukov manifold,

$$
\begin{equation*}
g(\Psi, C(X, Y) Z)=0 \tag{55}
\end{equation*}
$$

Moreover, on $(M, \widetilde{g}=\exp (2 \psi) g)$ the tensor field $\widetilde{L}$ given by (14) is a Codazzi tensor.

Proof. (52), (4) and (5) yield $-X \varphi=\omega[S(X, \Psi)-(\sigma+c)(X \psi)]$. Differentiating covariantly (53) and applying the Ricci identity, in view of
(3), (2) and the above equation, we have (55). Together with (16) and (17), this implies that $\widetilde{L}$ is a Codazzi tensor. This completes the proof.

Suppose that the 1-form $d \psi$ defines a geodesic mapping of a Sinyukov manifold $(M, g)$ onto a pseudo-Riemannian manifold $(M, \bar{g})$. Theorem 1 states that the manifold $(M, a)$, where $a$ is given by (3), admits a geodesic mapping onto the manifold $(M, \widetilde{g}=\exp (2 \psi) g)$ determined by the same 1-form $d \psi$.

Theorem 3 ([5]). A manifold $(M, g)(\operatorname{dim} M \geq 3)$ is a conformally flat Sinyukov manifold if and only if $(M, \widetilde{g}=\exp (2 \psi) g)$ is of constant sectional curvature.

From (2) and (10) we obtain
Lemma 5. On a manifold $(M, \widetilde{g}=\exp (2 \psi) g)$,

$$
\begin{equation*}
\left(\widetilde{\nabla}_{X} \bar{g}\right)(Y, Z)=\widetilde{\varphi}(Y) \widetilde{g}(X, Z)+\widetilde{\varphi}(Z) \widetilde{g}(X, Y) \tag{56}
\end{equation*}
$$

where $\widetilde{\varphi}(X)=\bar{g}(X, \Psi) \exp (-2 \psi)$. Thus, on $(M, \widetilde{g})$ the tensor $\bar{g}$ satisfies the same condition as does the tensor a on $(M, g)$.

THEOREM 4. Suppose that a manifold $(M, g)$ admits a non-trivial geodesic mapping onto a manifold $(M, \bar{g})$ defined by a 1-form $d \psi$. Let $U_{C}=$ $\{p \in M: C \neq 0$ at $p\}$, where $C$ is the Weyl conformal curvature tensor. Then $\left(U_{C}, g\right)$ is a Sinyukov manifold if and only if either
(i) $\left(U_{C}, \widetilde{g}=\exp (2 \psi) g\right)$ is an Einstein manifold which admits a geodesic mapping determined by the 1-form $-d \psi$, or
(ii) $\left(U_{C}, \widetilde{g}=\exp (2 \psi) g\right)$ is a Sinyukov manifold which admits a geodesic mapping determined by the 1-form $-d \psi$.

Proof. On $(M, \widetilde{g}=\exp (2 \psi) g)$, by (12)-(14), (52) and Proposition 4, we get

$$
\begin{equation*}
\widetilde{S}(X, Y)=(n-2) \bar{K} \bar{g}(X, Y)+\widetilde{K} \widetilde{g}(X, Y), \tag{57}
\end{equation*}
$$

where

$$
\begin{aligned}
\widetilde{K}= & \frac{\widetilde{r}}{2(n-1)}+\frac{r}{2(n-1)} \exp (-2 \psi) \\
& -(\sigma+c) \exp (-2 \psi)+\frac{n-2}{n} g(\Psi, \Psi) \exp (-2 \psi)
\end{aligned}
$$

Differentiating covariantly (57) and making use of (56) we have

$$
\begin{equation*}
\left(\widetilde{\nabla}_{Z} \widetilde{S}\right)(X, Y)=\widetilde{\nu}(X) \widetilde{g}(Y, Z)+\widetilde{\nu}(Y) \widetilde{g}(X, Z)+\widetilde{\sigma}(Z) \widetilde{g}(X, Y) \tag{58}
\end{equation*}
$$

where $\widetilde{\nu}(X)=(n-2) \bar{K} \widetilde{\varphi}(X), \widetilde{\sigma}(X)=X(\widetilde{K})$. As in [9], p. 131 (see also Lemma 2), one can prove that

$$
\begin{equation*}
\widetilde{\nu}(X)=\frac{n-2}{2(n-1)(n-2)}(X \widetilde{r}), \quad \widetilde{\sigma}(X)=\frac{n}{(n-1)(n+2)}(X \widetilde{r}) \tag{59}
\end{equation*}
$$

Consider the following two cases.
(i) The scalar curvature $\widetilde{r}$ of $(M, \widetilde{g})$ is constant. Since $(M, \widetilde{g})$ admits a non-trivial geodesic mapping onto $(M, a)$, we see, by the above considerations, that $\widetilde{r}=$ const if and only if $\bar{K}=0$. Then (57) implies that $(M, \widetilde{g})$ is an Einstein manifold. Conversely, if $(M, \widetilde{g})$ is an Einstein manifold which admits a geodesic mapping corresponding to $-d \psi$, then, as in [9], p. 130 (see also [12]), we easily conclude that $(M, g)$ is a Sinyukov manifold.
(ii) If $\widetilde{r}$ is not constant, then from (58) and (59) it follows that $(M, \widetilde{g})$ is a Sinyukov manifold. This completes the proof.

Notice that if $(M, \widetilde{g})$ is an Einstein manifold, then, by the results of [6], so is $(M, a)$. Hence and from Theorem 4 we have

Corollary 5. If $\widetilde{g}=\exp (2 \psi) g$ is an Einstein metric, then $\widetilde{a}=$ $\exp (2 \psi) a$ is a Sinyukov metric.
5. Local structure theorem. The local structure theorem for conformally flat Sinyukov manifolds is given in [5]. Let $a$ be a differentiable symmetric bilinear form on $U_{a} \subseteq M$ satisfying (3) and having $t$ different eigenvalues $\stackrel{1}{\lambda}, \ldots, \stackrel{t}{\lambda}$. From the very definition, at each point $p \in U_{a}$ they coincide with the eigenvalues of the endomorphism $A_{p}$ of the tangent space $T_{p}(M)$ corresponding to $a$, i.e. $g(A X, Y)=a(X, Y)$ for all $X, Y \in \mathfrak{X}\left(U_{a}\right)$. Let $(U, x)$ be a chart on $M$ such that $U \subseteq U_{a}$. Suppose that $\stackrel{\alpha}{v}$ is an eigenvector of the matrix $a_{i j}$ corresponding to the eigenvalue $\stackrel{\alpha}{\lambda}$, i.e. satisfying the condition

$$
\begin{equation*}
\left(a_{i j}-\stackrel{\alpha}{\lambda} g_{i j}\right) \stackrel{\alpha}{v}{ }^{j}=0 \tag{60}
\end{equation*}
$$

Following [4], one can prove that $\varphi_{t}{ }_{v}^{\alpha}$ t $=0$ and $\stackrel{\alpha}{v}_{i}=\stackrel{\alpha}{B}\left(X_{i} \stackrel{\alpha}{\lambda}\right)$, where $\stackrel{\alpha}{B} \in$ $\mathfrak{F}(U)$. Transvecting (60) with $\psi^{i}$ and making use of (4) and (5) we have $\psi_{i}{ }^{\alpha}{ }^{i}=0$. From [2], it follows that if $(M, g)$ admits a geodesic mapping then $\exp (-2 \psi)=\prod_{\alpha=1}^{t}\left(f_{\alpha}\right)^{\tau_{\alpha}}$, where $\tau_{\alpha}$ denotes the algebraic multiplicity of $\stackrel{\alpha}{\lambda}=f_{\alpha}\left(x^{n_{\alpha}+\tau_{\alpha}}\right), n_{1}=0, n_{\beta}=\tau_{1}+\ldots+\tau_{\beta-1}, \beta=2, \ldots, t$. Hence

$$
\begin{equation*}
\stackrel{\alpha}{v_{i_{\alpha}}}=\stackrel{\alpha}{F} \psi_{i_{\alpha}} \quad \text { and } \quad \stackrel{\alpha}{v_{j}}=0 \quad \text { for } j \neq i_{\alpha}, \tag{61}
\end{equation*}
$$

where $i_{\alpha}=n_{\alpha}+1, \ldots, n_{\alpha}+\tau_{\alpha}, \stackrel{\alpha}{F} \in \mathfrak{F}(U), \alpha=1, \ldots, t$.
Lemma 6. On a Sinyukov manifold the eigenvectors of the matrix $a_{i j}(p)$, $p \in U_{a}$, are non-null.

Proof. Suppose, to the contrary, that the eigenvector $\stackrel{\alpha}{v}$ corresponding to the eigenvalue $\stackrel{\alpha}{\lambda}$ is a null vector. Differentiating covariantly (61) with
respect to $x^{k}$, then transvecting the resulting equation with $\stackrel{\alpha}{v}^{i}$ and applying the relation $\psi_{j i_{\alpha}}=0$ for $j \neq j_{\alpha}$, we obtain $\psi_{k t}{ }_{v}^{\alpha}=0$. Therefore, from (53) and (60), we have

$$
-\frac{1}{(n-2) \omega}{ }^{\alpha} \lambda+K+\bar{K}\left(f_{\alpha}\right)^{-1} \prod_{\beta=1}^{t}\left(f_{\beta}\right)^{-\tau_{\beta}}=0
$$

Since

$$
X_{i} K=-\frac{2}{(n-2) \omega} \varphi_{i} \quad \text { and } \quad \varphi=\frac{1}{2} \sum_{\beta=1}^{t} \tau_{\beta} f_{\beta}
$$

(see [2]), it is easily seen that the above relation is false if the manifold admits a non-trivial geodesic mapping. This completes the proof.

Assume that a manifold $(M, g)$ admits a geodesic mapping onto a manifold $(M, \bar{g})$. If at $p \in M$ the eigenvectors of the matrix $a_{i j}(p)$ are non-null, then in some neighbourhood of $p$ there exists a coordinate system such that the components of the metric tensors $g$ and $\bar{g}$ take the form ([2])

$$
\begin{align*}
& g_{\mu \mu}=e_{\mu} \prod_{\substack{\beta=1 \\
\beta \neq \mu}}^{t}\left(f_{\beta}-f_{\mu}\right)^{\tau_{\beta}}, \quad \bar{g}_{\mu \mu}=\prod_{\beta=1}^{t}\left(f_{\beta}\right)^{-\tau_{\beta}}\left(f_{\mu}\right)^{-1} g_{\mu \mu} \\
& g_{i_{e} j_{e}}=\prod_{\substack{\beta=1 \\
\beta \neq \varrho}}^{t}\left(f_{\varrho}-f_{\beta}\right)^{\tau_{\beta}} \stackrel{g}{g}_{i_{e} j_{e}}, \quad \bar{g}_{i_{e} j_{\varrho}}=\prod_{\beta=1}^{t}\left(f_{\beta}\right)^{-\tau_{\beta}}\left(f_{\varrho}\right)^{-1} g_{i_{e} j_{\varrho}} \tag{62}
\end{align*}
$$

where $f_{\mu}=f_{\mu}\left(x^{\mu}\right), f_{\varrho}=\mathrm{const} \neq 0, e_{\mu}= \pm 1, \mu=1, \ldots, k, \varrho=k+1, \ldots, t$, $t \leq 2 k+1, \tau_{1}=\ldots=\tau_{k}=1, \tau_{\varrho}>1, i_{\varrho}, j_{\varrho}=n_{\varrho}+1, n_{\varrho}+2, \ldots, n_{\varrho}+\tau_{\varrho}$, $n_{1}=0, n_{\gamma}=\tau_{1}+\tau_{2}+\ldots+\tau_{\gamma-1}, \gamma=2, \ldots, t$ and $\stackrel{\varrho}{g}_{i_{e} j_{e}}\left(x^{n_{\varrho}+1}, \ldots, x^{n_{\varrho}+\tau_{\varrho}}\right)$ are metric tensors on $\tau_{\varrho}$-dimensional submanifolds $\stackrel{\varrho}{M}, \exp (-2 \psi)=\prod_{\alpha=1}^{t}\left(f_{\alpha}\right)^{\tau_{\alpha}}$.

The following lemma is a consequence of (25), (31), (15) and (52).
Lemma 7. If $(M, g)$ is a Sinyukov manifold and the Weyl conformal curvature tensor $C \neq 0$ at a point $p$ then, at $p$,

$$
\begin{equation*}
a(X, C(Y, Z) V)+a(V, C(Y, Z) X)=0 \tag{63}
\end{equation*}
$$

Taking into account (63) in the coordinate system in which the metric has the form (62) and applying the equality $a_{i_{\alpha} j_{\alpha}}=f_{\alpha} g_{i_{\alpha} j_{\alpha}}$ we find

Lemma 8. If $\stackrel{\varrho}{g}$ are metrics of one-dimensional manifolds, then the adjoint metric

$$
\stackrel{*}{g}=\sum_{\mu=1}^{k} \prod_{\substack{\beta=1 \\ \beta \neq \mu}}^{t}\left(f_{\beta}-f_{\mu}\right)^{\tau_{\beta}}\left(d x^{\mu}\right)^{2}+\sum_{\varrho=k+1}^{t} \prod_{\substack{\beta=1 \\ \beta \neq \varrho}}^{t}\left(f_{\varrho}-f_{\beta}\right)^{\tau_{\beta}}\left(d y^{\varrho}\right)^{2}
$$

is a metric of a conformally flat manifold. In particular, if $a_{i j}(p), p \in$ $U_{a}$, has $n$ distinct eigenvalues, then $\left(U_{a}, g\right)$ is a conformally flat Sinyukov manifold.

Theorem 5. Suppose that a 1-form $d \psi$ defines a geodesic mapping of a Sinyukov manifold $(M, g)$ with $C \neq 0$ everywhere on M. If $\widetilde{g}=\exp (2 \psi) g$ is a Sinyukov metric, then on a neighbourhood of each point $p \in M$ there exists a coordinate system such that the metrics $g$ and $\widetilde{g}$ take one of the following forms:
(i) if $k=1$ and $t=2$, then
(64) $\left.g=\frac{1}{4\left(c_{1}-x^{1}\right) W_{1}\left(x^{1}\right)}\left(d x^{1}\right)^{2}+\left(c_{1}-x^{1}\right)\right)_{\alpha \beta}^{2} d x^{\alpha} d x^{\beta}, \quad \widetilde{g}=\left(x^{1}\right)^{-1} g$, where

$$
W_{1}(z)=A_{2} z^{2}+A_{1} z+A_{0},
$$

$A_{0}, A_{2}, c_{1}=$ const $\neq 0, A_{1}=$ const, $\stackrel{2}{h}=\stackrel{2}{h}\left(x^{2}, \ldots, x^{n}\right)$ is an $(n-1)$ dimensional Einstein metric with the Ricci tensor

$$
\stackrel{2}{S}=-(n-2) W_{1}\left(c_{1}\right) \stackrel{2}{h}
$$

$\alpha, \beta=2, \ldots, n$;
(ii) if $k=1$ and $t=3$, then we have
(65) $g=\frac{(n-2) \omega}{W_{2}\left(x^{1}\right)}\left(d x^{1}\right)^{2}+\sum_{\varrho=2}^{3}\left(c_{\varrho}-x^{1}\right)^{\varrho}{\stackrel{\varrho}{i_{\varrho} j_{\varrho}}} d x^{i_{\varrho}} d x^{j_{\varrho}}, \quad \widetilde{g}=\left(x^{1}\right)^{-1} g$,
where

$$
W_{2}(z)=4\left(c_{2}-z\right)\left(c_{3}-z\right)\left(c_{1}+z\right)
$$

$c_{\varrho}, c_{1}, \omega=\mathrm{const} \neq 0, \varrho=2,3, \stackrel{2}{h}=\stackrel{2}{h}\left(x^{2}, \ldots, x^{\tau_{2}+1}\right)$ is a $\tau_{2}$-dimensional Einstein metric with the Ricci tensor

$$
\stackrel{2}{S}=\left(\tau_{2}-1\right)\left(c_{2}-c_{3}\right)\left(c_{1}+c_{2}\right) K \stackrel{2}{h}
$$

$\stackrel{3}{h}=\stackrel{3}{h}\left(x^{\tau_{2}+2}, \ldots, n\right)$ is a $\tau_{3}$-dimensional Einstein metric with te Ricci tensor

$$
\stackrel{3}{S}=\left(\tau_{3}-1\right)\left(c_{3}-c_{2}\right)\left(c_{1}+c_{3}\right) K \stackrel{3}{h}
$$

$K=\frac{1}{(n-2) \omega}, i_{2}, j_{2}=2, \ldots, \tau_{2}+1, i_{3}, j_{3}=\tau_{2}+2, \ldots, n, 1+\tau_{2}+\tau_{3}=n ;$
(iii) if $k>1$, then

$$
\begin{align*}
& g=\sum_{\substack{\mu=1}}^{k} \prod_{\substack{\eta=1 \\
\eta \neq \mu}}^{k} \frac{x^{\eta}-x^{\mu}}{W_{3}\left(x^{\mu}\right)}\left(d x^{\mu}\right)^{2}+\sum_{\varrho=k+1}^{t} \prod_{\mu=1}^{k}\left(f_{\varrho}-x^{\mu}\right)^{\varrho}{\stackrel{\varrho}{i_{e} j_{e}}} d x^{i_{e}} d x^{j_{e}},  \tag{66}\\
& \widetilde{g}=\left(x^{1} \ldots x^{k}\right)^{-1} g
\end{align*}
$$

where

$$
W_{3}(z)=(-1)^{k+1} 4 A_{k+2} z^{k+2}+A_{k+1} z^{k+1}+\ldots+A_{1} z+4 A_{0}
$$

$A_{0}, A_{1}, \ldots, A_{k+2}=\mathrm{const}, A_{0}, A_{k+2} \neq 0, f_{\varrho}=\mathrm{const} \neq 0$, and the $f_{\varrho}$ are roots of the polynomial $W_{3}, \stackrel{\varrho}{g}$ are $\tau_{\varrho}$-dimensional Einstein metrics with the Ricci tensors

$$
\stackrel{\varrho}{S}=\left(\tau_{\varrho}-1\right) K_{\varrho} \stackrel{\varrho}{g}
$$

and $K_{\varrho}=(-1)^{k+1} \frac{1}{4} W_{3}^{\prime}\left(f_{\varrho}\right), k>1, t \leq 2 k+1, \varrho=k+1, \ldots, t, i_{\varrho}, j_{\varrho}=$ $n_{\varrho}+1, \ldots, n_{\varrho}+\tau_{\varrho}, n_{1}=0, n_{\gamma}=\tau_{1}+\tau_{2}+\ldots+\tau_{\gamma-1}, \gamma=2, \ldots, t, \tau_{\varrho}>1$.

Proof. Solving (53) in the local coordinate system in which $g$ and $\bar{g}$ are of the form (62) and using the equality $a_{i_{\alpha} j_{\alpha}}=f_{\alpha} g_{i_{\alpha} j_{\alpha}}$, in the same way as in the proof of Theorem 3 of [1], we obtain our assertion.

Theorem 6. Let $\mathbb{R}^{n}$ be endowed with a metric of the form either (64) or (65) or (66), where $h, \stackrel{2}{h}$ or $\stackrel{3}{h}$ and at least one of the forms $\stackrel{\varrho}{g}$ are nonconformally flat Einstein metrics. Then $\left(\mathbb{R}^{n}, g\right)\left(\right.$ and $\left.\left(\mathbb{R}^{n}, \widetilde{g}\right)\right)$ is a nonconformally flat Sinyukov manifold.

Proof. By elementary computation one can easily verify that (9) holds on $\left(\mathbb{R}^{n}, g\right)$ (and the analogous condition is satisfied on $\left(\mathbb{R}^{n}, \widetilde{g}\right)$ ). The components of the 1-form $\sigma\left(\nu=\frac{n-2}{2 n} \sigma\right)$ are respectively:

1) for the metric (64):

$$
\sigma_{1}=-n A_{2}, \sigma_{\alpha}=0 \quad\left(\widetilde{\sigma}_{1}=\frac{-n A_{0} c_{1}}{\left(x^{1}\right)^{2}}, \widetilde{\sigma}_{\alpha}=0\right), \quad \alpha=2, \ldots, n
$$

2) for the metric (65):

$$
\begin{aligned}
& \sigma_{1}=\frac{n}{(n-2) \omega}, \sigma_{\alpha}=0 \quad\left(\widetilde{\sigma}_{1}=\frac{-n c_{1} c_{2} c_{3}}{(n-2) \omega\left(x^{1}\right)^{2}}, \widetilde{\sigma}_{\alpha}=0\right), \\
& \\
& \alpha=2, \ldots, n,
\end{aligned}
$$

3) for the metric (66):

$$
\sigma_{\mu}=n A_{k+2}, \sigma_{i_{e}}=0 \quad\left(\tilde{\sigma}_{\mu}=\frac{-n A_{0}}{\left(x^{\mu}\right)^{2}}, \widetilde{\sigma}_{i_{e}}=0\right), \quad \mu=1, \ldots, k
$$

Moreover, in the metrics (64), (65) and (66), the conformal curvature tensor $C \neq 0$ if and only if $\stackrel{2}{h}$ (resp. $\stackrel{2}{h}$ or $\stackrel{3}{h}$, resp. at least one of $\stackrel{\varrho}{g}$ ) is a nonconformally flat metric. This completes the proof.

Remark. In [4] the local structure theorem for Einstein manifolds admitting geodesic mappings is proved. If $\widetilde{g}=\exp (2 \psi) g$ is an Einstein manifold, then, by Theorem 4(i), Corollary 5 and the results of [4], the local structure of Sinyukov manifolds can be easily obtained. This, together
with Theorem 5, provides a complete description of the local structure of Sinyukov manifolds.

From Theorems 5, 6 and the results of [4] we have the following
Corollary 6. If $M$ is a Sinyukov manifold and $\operatorname{dim} M \leq 4$, then $M$ is conformally flat.

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