

Algebraic relations for reciprocal sums of Fibonacci numbers

by

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1. Introduction. Let $\{F_n\}_{n \geq 0}$ and $\{L_n\}_{n \geq 0}$ be Fibonacci numbers and Lucas numbers defined by

$$\begin{aligned} F_0 &= 0, & F_1 &= 1, & F_{n+2} &= F_{n+1} + F_n & (n \geq 0), \\ L_0 &= 2, & L_1 &= 1, & L_{n+2} &= L_{n+1} + L_n & (n \geq 0). \end{aligned}$$

Duverney, Ke. Nishioka, Ku. Nishioka, and the last named author [9] (see also [8]) proved the transcendence of the numbers

$$\sum_{n=1}^{\infty} \frac{1}{F_n^{2s}}, \quad \sum_{n=1}^{\infty} \frac{1}{L_n^{2s}}, \quad \sum_{n=1}^{\infty} \frac{1}{F_{2n-1}^s}, \quad \sum_{n=1}^{\infty} \frac{1}{L_{2n}^s} \quad (s = 1, 2, 3, \dots)$$

by using Nesterenko's theorem on the Ramanujan functions $P(q)$, $Q(q)$, and $R(q)$ (see Section 3).

In this paper, we prove the algebraic independence of the numbers

$$\sum_{n=1}^{\infty} \frac{1}{F_n^{2s}}, \quad \sum_{n=1}^{\infty} \frac{1}{F_n^{4s}}, \quad \sum_{n=1}^{\infty} \frac{1}{F_n^{6s}} \quad \left(\text{respectively, } \sum_{n=1}^{\infty} \frac{1}{L_n^{2s}}, \quad \sum_{n=1}^{\infty} \frac{1}{L_n^{4s}}, \quad \sum_{n=1}^{\infty} \frac{1}{L_n^{6s}} \right)$$

and write each

$$\sum_{n=1}^{\infty} \frac{1}{F_n^{2s}} \quad \left(\text{respectively, } \sum_{n=1}^{\infty} \frac{1}{L_n^{2s}} \right) \quad (s = 4, 5, 6, \dots)$$

as a rational (respectively, algebraic) function of these three numbers over \mathbb{Q} (see Theorems 1, 3, and Example 1). Similar results are obtained for the alternating sums

$$\sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{F_n^{2s}} \quad \left(\text{respectively, } \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{L_n^{2s}} \right) \quad (s = 1, 2, 3, \dots)$$

(see Theorems 2 and 4).

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It is interesting to compare our results with some arithmetical properties of the values of the Riemann zeta function $\zeta(s) = \sum_{n=1}^{\infty} n^{-s}$:

(a) Apéry [3] proved the irrationality of $\zeta(3)$ (cf. [4], [6], [15]). It is still unknown whether $\zeta(5)$ is irrational or not. Zudilin [20] showed that at least one of $\zeta(5)$, $\zeta(7)$, $\zeta(9)$, $\zeta(11)$ is irrational.

(b) Euler's formula

$$\zeta(2k) = \frac{(-1)^{k-1} 2^{2k-1} B_{2k}}{(2k)!} \pi^{2k} \quad (k \in \mathbb{N}),$$

where B_{2k} are the Bernoulli numbers, implies the algebraic dependence of $\zeta(2k)$ on $\zeta(2) = \pi^2/6$ for any integer $k \geq 2$.

For Fibonacci numbers we consider the Fibonacci zeta function

$$\zeta_{\mathbb{F}}(s) = \sum_{n=1}^{\infty} \frac{1}{F_n^s} \quad (\operatorname{Re}(s) > 0),$$

which extends meromorphically to the whole complex plane (cf. [14]).

(a') André-Jeannin [2] proved the irrationality of $\zeta_{\mathbb{F}}(1)$ (see also [5], [7]). The arithmetical nature of $\zeta_{\mathbb{F}}(3)$ is unknown.

(b') Our results in this paper imply that the values $\zeta_{\mathbb{F}}(2)$, $\zeta_{\mathbb{F}}(4)$, $\zeta_{\mathbb{F}}(6)$ are algebraically independent, and that for any integer $s \geq 4$,

$$\zeta_{\mathbb{F}}(2s) - 5^{s-2} r_s \zeta_{\mathbb{F}}(4) \in \mathbb{Q}(u, v), \quad u := \zeta_{\mathbb{F}}(2), \quad v := \zeta_{\mathbb{F}}(6),$$

with some $r_s \in \mathbb{Q}$ ($r_s = 0$ if and only if s is odd), where the rational function of u and v is explicit; for example,

$$\begin{aligned} \zeta_{\mathbb{F}}(8) - \frac{15}{14} \zeta_{\mathbb{F}}(4) &= \frac{1}{378(4u+5)^2} (256u^6 - 3456u^5 + 2880u^4 + 1792u^3v \\ &\quad - 11100u^3 + 20160u^2v - 10125u^2 + 7560uv + 3136v^2 - 1050v) \end{aligned}$$

(see Theorem 1, Example 1, and Table 1).

2. Statement of results. Suppose that $\alpha, \beta \in \mathbb{C}$ satisfy $|\beta| < 1$ and $\alpha\beta = -1$. We put

$$(1) \quad U_n = \frac{\alpha^n - \beta^n}{\alpha - \beta} \quad (n \geq 0),$$

$$(2) \quad V_n = \alpha^n + \beta^n \quad (n \geq 0).$$

If $\alpha + \beta = a \in \mathbb{Z}$, then $\{U_n\}_{n \geq 0}$ and $\{V_n\}_{n \geq 0}$ represent integer-valued, linearly independent binary recurrences satisfying the relation

$$X_{n+2} = aX_{n+1} + X_n \quad (n \geq 0)$$

with initial values $(X_0, X_1) = (0, 1)$ and $(2, a)$, respectively. In particular, if $\beta = (1 - \sqrt{5})/2$, we have the Fibonacci and Lucas numbers:

$$U_n = F_n, \quad V_n = L_n \quad (n \geq 0).$$

In what follows s always denotes a positive integer. Set $\sigma_0(s) = 1$, and for $s \geq 2$ let $\sigma_1(s), \dots, \sigma_{s-1}(s)$ be the elementary symmetric functions of the $s - 1$ numbers $-1, -2^2, \dots, -(s - 1)^2$ defined by

$$\sigma_i(s) = (-1)^i \sum_{1 \leq r_1 < \dots < r_i \leq s-1} r_1^2 \cdots r_i^2 \quad (1 \leq i \leq s - 1).$$

The coefficients of the expansions

$$\operatorname{cosec}^2 x = \frac{1}{x^2} + \sum_{j=0}^{\infty} a_j x^{2j}, \quad \sec^2 x = \sum_{j=0}^{\infty} b_j x^{2j}$$

are given by

$$a_{j-1} = \frac{(-1)^{j-1} (2j - 1) 2^{2j} B_{2j}}{(2j)!},$$

$$b_{j-1} = \frac{(-1)^{j-1} (2j - 1) 2^{2j} (2^{2j} - 1) B_{2j}}{(2j)!}$$

($j \geq 1$), where $B_2 = 1/6$, $B_4 = -1/30$, $B_6 = 1/42$, ... are the Bernoulli numbers. (In this paper, the symbol B_n is used for the Bernoulli numbers only here. It will denote a q -series in Sections 3 and 4.)

THEOREM 1. *Let $\{U_n\}_{n \geq 1}$ be defined by (1) with $\alpha, \beta \in \overline{\mathbb{Q}}$ satisfying $|\beta| < 1$ and $\alpha\beta = -1$, and set*

$$\Phi_{2s} := (\alpha - \beta)^{-2s} \sum_{n=1}^{\infty} \frac{1}{U_n^{2s}}.$$

Then the numbers Φ_2, Φ_4, Φ_6 are algebraically independent, and for any integer $s \geq 4$ the number Φ_{2s} is written as

$$\Phi_{2s} = \frac{1}{(2s - 1)!}$$

$$\times \left(\sigma_{s-1}(s) \mu_s - \sum_{j=1}^{s-1} \frac{(-1)^j (2j)!}{2^{2j+3}} \sigma_{s-j-1}(s) (\varphi_j - (-1)^s \psi_j - a_j) \right)$$

with

$$\mu_s = \begin{cases} \Phi_2 & (s \text{ odd}), \\ \frac{1}{3} \left(4\Phi_2^2 + 2\Phi_2 - 18\Phi_4 + \omega - \frac{5}{4} \right) & (s \text{ even}), \end{cases}$$

$$\varphi_1 = \frac{4}{3} \left(32\Phi_2^2 - 5\Phi_2 - \omega + \frac{13}{10} \right),$$

$$\begin{aligned}\varphi_2 &= -\frac{4}{63}(24\Phi_2 - 1)\left(112\Phi_2^2 - 21\Phi_2 - 5\omega + \frac{77}{12}\right), \\ \varphi_j &= \frac{3}{(j-2)(2j+3)}\sum_{i=1}^{j-2}\varphi_i\varphi_{j-i-1} \quad (j \geq 3), \\ \psi_1 &= \frac{4}{3}\left(16\Phi_2^2 - 13\Phi_2 - 5\omega + \frac{25}{4}\right), \\ \psi_2 &= \frac{4}{9}(24\Phi_2 - 1)\left(16\Phi_2^2 - 13\Phi_2 - 5\omega + \frac{25}{4}\right), \\ \psi_j &= \frac{1}{j(2j-1)}\left(2(24\Phi_2 - 1)\psi_{j-1} - 3\sum_{i=1}^{j-2}\psi_i\psi_{j-i-1}\right) \quad (j \geq 3),\end{aligned}$$

where

$$\omega = \frac{56\Phi_6 + 5/4}{4\Phi_2 + 1}.$$

REMARK 1. If $s \geq 4$, then $(1 + 4\Phi_2)^{\lfloor s/2 \rfloor}(\Phi_{2s} - r_s\Phi_4) \in \mathbb{Q}[\Phi_2, \Phi_6]$, and the total degree of this does not exceed $s + \lfloor s/2 \rfloor$, where $r_s \in \mathbb{Q}$ ($r_s = 0$ if and only if s is odd).

Table 1. Relations for Φ_{2s} in Theorem 1

	$x = \Phi_2$	$y = \Phi_4$	$z = \Phi_6$
$s = 4$	$\Phi_8 = \frac{3}{70}y + \frac{1}{1890(4x+1)^2}(1280x^6 - 3456x^5 + 576x^4 + 8960x^3z - 444x^3 + 20160x^2z - 81x^2 + 1512xz + 15680z^2 - 42z)$		
$s = 5$	$\Phi_{10} = \frac{1}{297(4x+1)^2}(512x^7 - 704x^6 + 162x^5 - 1600x^4z - 30x^4 + 2560x^3z - 15x^3 + 450x^2z + 4760xz^2 + 75xz + 700z^2 + 15z)$		
$s = 6$	$\Phi_{12} = \frac{1}{462}y + \frac{1}{162162(4x+1)^3}(1089536x^9 - 774144x^8 - 111168x^7 - 4558848x^6z + 53424x^6 + 774144x^5z - 40608x^5 + 1737792x^4z + 8558592x^3z^2 - 11556x^4 + 311472x^3z + 7547904x^2z^2 - 3258x^3 + 77112x^2z + 1202544xz^2 + 2458624z^3 - 351x^2 + 3780xz + 127260z^2 - 252z)$		

THEOREM 2. Let $\{U_n\}_{n \geq 1}$ be defined as in Theorem 1, and set

$$\Phi_{2s}^* := (\alpha - \beta)^{-2s} \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{U_n^{2s}}.$$

Then the numbers Φ_2^* , Φ_4^* , Φ_6^* are algebraically independent, and for any integer $s \geq 4$ the number Φ_{2s}^* is written as

$$\Phi_{2s}^* = \frac{1}{(2s-1)!} \times \left(\sigma_{s-1}(s)\mu_s + \sum_{j=1}^{s-1} \frac{(-1)^j(2j)!}{2^{2j+3}} \sigma_{s-j-1}(s)(\varphi_j + (-1)^s\psi_j - a_j) \right)$$

with

$$\begin{aligned} \mu_s &= \begin{cases} \Phi_2^* & (s \text{ odd}), \\ \frac{1}{24}(4\xi - 1) & (s \text{ even}), \end{cases} \\ \varphi_1 &= -\frac{4}{45} \left(180\Phi_4^* - 10\xi^2 + 5\xi - \frac{11}{8} \right), \\ \varphi_2 &= -\frac{16}{189} \xi \left(180\Phi_4^* - 6\xi^2 + 5\xi - \frac{11}{8} \right), \\ \varphi_j &= \frac{3}{(j-2)(2j+3)} \sum_{i=1}^{j-2} \varphi_i \varphi_{j-i-1} \quad (j \geq 3), \\ \psi_1 &= -\frac{4}{9} \left(180\Phi_4^* + 2\xi^2 + 5\xi - \frac{11}{8} \right), \\ \psi_2 &= \frac{16}{27} \xi \left(180\Phi_4^* + 2\xi^2 + 5\xi - \frac{11}{8} \right), \\ \psi_j &= -\frac{1}{j(2j-1)} \left(8\xi\psi_{j-1} + 3 \sum_{i=1}^{j-2} \psi_i \psi_{j-i-1} \right) \quad (j \geq 3), \end{aligned}$$

where $\xi = \xi(\Phi_2^*, \Phi_4^*, \Phi_6^*)$ is a number satisfying

$$(3) \quad 8\xi^3 + 5\xi^2 + (1440\Phi_4^* - 46)\xi - \left(252\Phi_2^* + 1260\Phi_4^* - 7560\Phi_6^* - \frac{177}{16} \right) = 0.$$

Table 2. Relations for Φ_{2s}^* in Theorem 2

	$x = \Phi_2^*$	$y = \Phi_4^*$
$s = 4$	$\Phi_8^* = \frac{1}{3}y^2 + \left(\frac{5}{189}\xi^2 - \frac{11}{378}\xi + \frac{23}{432} \right)y$	$+\frac{1}{6804}\xi^4 - \frac{1}{6804}\xi^3 - \frac{23}{18144}\xi^2 + \frac{71}{108864}\xi - \frac{143}{1741824}$
$s = 5$	$\Phi_{10}^* = \frac{1}{630}x - \left(\frac{80}{693}\xi - \frac{10}{63} \right)y^2 - \left(\frac{8}{2079}\xi^3 - \frac{5}{2079}\xi^2 + \frac{5}{756}\xi - \frac{5}{756} \right)y$	$-\frac{1}{56133}\xi^5 + \frac{1}{64152}\xi^4 + \frac{403}{4490640}\xi^3 - \frac{5}{40824}\xi^2 + \frac{1007}{3265920}\xi - \frac{40739}{574801920}$

THEOREM 3. *Let $\{V_n\}_{n \geq 1}$ be defined by (2) with $\alpha, \beta \in \overline{\mathbb{Q}}$ satisfying $|\beta| < 1$ and $\alpha\beta = -1$, and set*

$$\Psi_{2s} := \sum_{n=1}^{\infty} \frac{1}{V_n^{2s}}.$$

Then the numbers Ψ_2, Ψ_4, Ψ_6 are algebraically independent, and for any integer $s \geq 4$ the number Ψ_{2s} is written as

$$\Psi_{2s} = \frac{1}{(2s-1)!} \left(\sigma_{s-1}(s)\mu_s + \sum_{j=1}^{s-1} \frac{(-1)^j (2j)!}{2^{2j+3}} \sigma_{s-j-1}(s)(\psi_j - (-1)^s(\varphi_j - b_j)) \right)$$

with

$$\mu_s = \begin{cases} \Psi_2 & (s \text{ odd}), \\ 4\Psi_2^2 + \Psi_2 - 6\Psi_4 & (s \text{ even}), \end{cases}$$

$$\varphi_1 = \frac{1}{2} (8\Psi_2 + 1)(8\Psi_2 + \eta + 1),$$

$$\varphi_2 = \frac{1}{12} (8\Psi_2 + 1)(8\Psi_2 + \eta + 1)(24\Psi_2 + \eta + 3),$$

$$\varphi_j = \frac{1}{j(2j-1)} \left((24\Psi_2 + \eta + 3)\varphi_{j-1} + 3 \sum_{i=1}^{j-2} \varphi_i \varphi_{j-i-1} \right) \quad (j \geq 3),$$

$$\psi_1 = -\frac{1}{2} (8\Psi_2 + 1)(8\Psi_2 - \eta + 1),$$

$$\psi_2 = \frac{1}{12} (8\Psi_2 + 1)(8\Psi_2 - \eta + 1)(24\Psi_2 - \eta + 3),$$

$$\psi_j = -\frac{1}{j(2j-1)} \left((24\Psi_2 - \eta + 3)\psi_{j-1} + 3 \sum_{i=1}^{j-2} \psi_i \psi_{j-i-1} \right) \quad (j \geq 3),$$

Table 3. Relations for Ψ_{2s} in Theorem 3

	$x = \Psi_2$	$y = \Psi_4$
$s = 4$	$\Psi_8 = \frac{3}{70}y + \frac{8}{35}x^4 + \frac{4}{35}x^3 + \left(\frac{1}{315}\eta^2 + \frac{1}{45}\eta + \frac{2}{63} \right)x^2$ $+ \left(\frac{1}{1260}\eta^2 + \frac{1}{180}\eta + \frac{11}{2520} \right)x + \frac{1}{20160}\eta^2 + \frac{1}{2880}\eta - \frac{1}{2520}$	
$s = 5$	$\Psi_{10} = \frac{2}{15}x^5 + \frac{1}{12}x^4 + \left(\frac{17}{5040}\eta^2 + \frac{11}{504}\eta + \frac{41}{720} \right)x^3$ $+ \left(\frac{17}{13440}\eta^2 + \frac{11}{1344}\eta + \frac{31}{1920} \right)x^2$ $+ \left(\frac{1}{5806080}\eta^4 + \frac{1}{96768}\eta^3 + \frac{67}{193536}\eta^2 + \frac{625}{290304}\eta + \frac{2221}{645120} \right)x$ $+ \frac{1}{46448640}\eta^4 + \frac{1}{774144}\eta^3 + \frac{233}{7741440}\eta^2 + \frac{61}{331776}\eta - \frac{1111}{5160960}$	

where $\eta = \eta(\Psi_2, \Psi_6)$ is a number satisfying

$$(\eta + 5)^2 = -192\Psi_2^2 - 48\Psi_2 + 6 + \frac{3840\Psi_6 + 30}{8\Psi_2 + 1}.$$

THEOREM 4. *Let $\{V_n\}_{n \geq 1}$ be defined as in Theorem 3, and set*

$$\Psi_{2s}^* := \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{V_n^{2s}}.$$

Then the numbers Ψ_2^ , Ψ_4^* , Ψ_6^* are algebraically independent, and for any integer $s \geq 4$ the number Ψ_{2s}^* is written as*

$$\Psi_{2s}^* = \frac{1}{(2s-1)!} \left(\sigma_{s-1}(s)\mu_s + \sum_{j=1}^{s-1} \frac{(-1)^j (2j)!}{2^{2j+3}} \sigma_{s-j-1}(s)(\psi_j + (-1)^s(\varphi_j - b_j)) \right)$$

with

$$\mu_s = \begin{cases} \Psi_2^* & (s \text{ odd}), \\ \frac{1}{8}(\theta - 1) & (s \text{ even}), \end{cases}$$

$$\varphi_1 = -\frac{1}{2}(96\Psi_4^* - \theta^2 + 2\theta - 3),$$

$$\varphi_2 = \frac{1}{12\theta}(96\Psi_4^* - \theta^2 + 2\theta - 3)(96\Psi_4^* - 3\theta^2 + 2\theta - 3),$$

$$\varphi_j = -\frac{1}{j(2j-1)} \left((96\Psi_4^* - 3\theta^2 + 2\theta - 3) \frac{\varphi_{j-1}}{\theta} - 3 \sum_{i=1}^{j-2} \varphi_i \varphi_{j-i-1} \right) \quad (j \geq 3),$$

$$\psi_1 = -\frac{1}{2}(96\Psi_4^* + \theta^2 + 2\theta - 3),$$

$$\psi_2 = \frac{1}{12\theta}(96\Psi_4^* + \theta^2 + 2\theta - 3)(96\Psi_4^* + 3\theta^2 + 2\theta - 3),$$

$$\psi_j = -\frac{1}{j(2j-1)} \left((96\Psi_4^* + 3\theta^2 + 2\theta - 3) \frac{\psi_{j-1}}{\theta} + 3 \sum_{i=1}^{j-2} \psi_i \psi_{j-i-1} \right) \quad (j \geq 3),$$

Table 4. Relations for Ψ_{2s}^* in Theorem 4

	$x = \Psi_2^*$	$y = \Psi_4^*$
$s = 4$	$\Psi_8^* = \frac{1}{35840\theta^2}(98304y^3 - 1024(22\theta + 9)y^2 + 32(11\theta^4 + 32\theta^2 + 44\theta + 9)y - 2\theta^5 - 11\theta^4 + 38\theta^2 - 22\theta - 3)$	
$s = 5$	$\Psi_{10}^* = \frac{1}{630}x + \frac{1}{645120\theta}(786432y^3 - 73728(\theta + 1)y^2 + 576(\theta^4 + 3\theta^2 + 8\theta + 4)y - 3\theta^5 - 18\theta^4 - 16\theta^3 - 54\theta^2 + 115\theta - 24)$	

where $\theta = \theta(\Psi_2^*, \Psi_4^*, \Psi_6^*)$ is a number satisfying

$$\theta^2 - (192\Psi_4^* - 6)\theta + 1920\Psi_6^* - 64\Psi_2^* - 7 = 0.$$

REMARK 2. The quantities ξ , η , and θ in Theorems 2, 3, and 4, respectively, are algebraic functions of the corresponding sums for $1 \leq s \leq 3$. In applying these theorems, we have to choose appropriate branches of them depending on the parameter a (or β). This will be done in Section 5 (Theorems 5, 6, and 7).

EXAMPLE 1. The reciprocal sums of Fibonacci (respectively, Lucas) numbers $\sum_{n=1}^{\infty} F_n^{-2}$, $\sum_{n=1}^{\infty} F_n^{-4}$, $\sum_{n=1}^{\infty} F_n^{-6}$ (respectively, $\sum_{n=1}^{\infty} L_n^{-2}$, $\sum_{n=1}^{\infty} L_n^{-4}$, $\sum_{n=1}^{\infty} L_n^{-6}$) are algebraically independent, and for any integer $s \geq 4$ the number $5^{-s} \sum_{n=1}^{\infty} F_n^{-2s}$ (respectively, $\sum_{n=1}^{\infty} L_n^{-2s}$) is written by the formula in Theorem 1 (respectively, Theorem 3 with η in Theorem 6). The alternating sums $\sum_{n=1}^{\infty} (-1)^{n+1} F_n^{-2s}$ and $\sum_{n=1}^{\infty} (-1)^{n+1} L_n^{-2s}$ have similar properties mentioned in Theorems 2 and 4 with ξ and θ in Theorems 5 and 7, respectively.

EXAMPLE 2. The *Pell numbers* 1, 2, 5, 12, 29, ... defined by

$$P_0 = 0, \quad P_1 = 1, \quad P_{n+2} = 2P_{n+1} + P_n \quad (n \geq 0)$$

are expressible by (1) with $\beta = 1 - \sqrt{2}$ and $\alpha = -1/\beta$ satisfying $\alpha - \beta = \sqrt{8}$ (see [13]). To the numbers $\sum_{n=1}^{\infty} P_n^{-2s}$ and $\sum_{n=1}^{\infty} (-1)^{n+1} P_n^{-2s}$, we can apply Theorems 1 and 2 (with Theorem 5).

3. Preparation for the proofs. The reciprocal sums of $\{U_n\}_{n \geq 1}$ and $\{V_n\}_{n \geq 1}$ in our theorems are written as series of hyperbolic functions. In [19] Zucker gave a method of summing such series. He wrote them as q -series expressible in closed form in terms of K , E , and k , where K and E are the complete elliptic integrals of the first and second kind with the modulus $k \neq 0, \pm 1$ defined by

$$K = K(k) = \int_0^1 \frac{dt}{\sqrt{(1-t^2)(1-k^2t^2)}}, \quad E = \int_0^1 \sqrt{\frac{1-k^2t^2}{1-t^2}} dt$$

for $k^2 \in \mathbb{C} \setminus (\{0\} \cup \{z \mid z \geq 1\})$. The branch of each integrand is chosen so that it tends to 1 as $t \rightarrow 0$. The relation among q and these quantities is as follows:

$$q = e^{-\pi c}, \quad c = K'/K, \quad K' = K(k'), \quad k^2 + (k')^2 = 1.$$

We start with the following formulas. We choose $c = c(\beta)$ (or $q = q(\beta)$) so that $q = e^{-\pi c} = \beta^2$, $\beta = -e^{-\pi c/2}$. Then by [19, Tables 1(i), 1(ii), 1(iv), 1(iii)],

$$(4) \quad \Sigma_1 := 2^{-2s} \sum_{\nu=1}^{\infty} \operatorname{cosech}^{2s}(\nu\pi c) = \frac{1}{(2s-1)!} \sum_{j=0}^{s-1} \sigma_{s-j-1}(s) A_{2j+1}(\beta^2),$$

$$(5) \quad \Sigma_2 := 2^{-2s} \sum_{\nu=1}^{\infty} \operatorname{sech}^{2s}(\nu\pi c) = \frac{(-1)^{s-1}}{(2s-1)!} \sum_{j=0}^{s-1} \sigma_{s-j-1}(s) B_{2j+1}(\beta^2),$$

$$(6) \quad \Sigma_3 := 2^{-2s} \sum_{\nu=1}^{\infty} \operatorname{sech}^{2s} \frac{(2\nu-1)\pi c}{2} = \frac{(-1)^{s-1}}{(2s-1)!} \sum_{j=0}^{s-1} \sigma_{s-j-1}(s) D_{2j+1}(\beta^2),$$

$$(7) \quad \Sigma_4 := 2^{-2s} \sum_{\nu=1}^{\infty} \operatorname{cosech}^{2s} \frac{(2\nu-1)\pi c}{2} = \frac{1}{(2s-1)!} \sum_{j=0}^{s-1} \sigma_{s-j-1}(s) C_{2j+1}(\beta^2)$$

(note that $\sigma_i(s)$ denotes $\alpha_i(s)$ in [19]), where

$$A_{2j+1}(q) = \sum_{n=1}^{\infty} \frac{n^{2j+1} q^{2n}}{1-q^{2n}}, \quad B_{2j+1}(q) = \sum_{n=1}^{\infty} \frac{(-1)^{n+1} n^{2j+1} q^{2n}}{1-q^{2n}},$$

$$C_{2j+1}(q) = \sum_{n=1}^{\infty} \frac{n^{2j+1} q^n}{1-q^{2n}}, \quad D_{2j+1}(q) = \sum_{n=1}^{\infty} \frac{(-1)^{n+1} n^{2j+1} q^n}{1-q^{2n}}.$$

Our reciprocal sums are expressible by these series of hyperbolic functions:

$$(8) \quad \Phi_{2s} = (\alpha - \beta)^{-2s} \left(\sum_{\nu=1}^{\infty} \frac{1}{U_{2\nu-1}^{2s}} + \sum_{\nu=1}^{\infty} \frac{1}{U_{2\nu}^{2s}} \right) = \Sigma_3 + \Sigma_1,$$

$$(9) \quad \Psi_{2s} = \sum_{\nu=1}^{\infty} \frac{1}{V_{2\nu-1}^{2s}} + \sum_{\nu=1}^{\infty} \frac{1}{V_{2\nu}^{2s}} = \Sigma_4 + \Sigma_2,$$

$$(10) \quad \Phi_{2s}^* = (\alpha - \beta)^{-2s} \left(\sum_{\nu=1}^{\infty} \frac{1}{U_{2\nu-1}^{2s}} - \sum_{\nu=1}^{\infty} \frac{1}{U_{2\nu}^{2s}} \right) = \Sigma_3 - \Sigma_1,$$

$$(11) \quad \Psi_{2s}^* = \sum_{\nu=1}^{\infty} \frac{1}{V_{2\nu-1}^{2s}} - \sum_{\nu=1}^{\infty} \frac{1}{V_{2\nu}^{2s}} = \Sigma_4 - \Sigma_2.$$

The q -series A_{2j+1} , B_{2j+1} , C_{2j+1} , and D_{2j+1} are generated from Fourier expansions of the squares of Jacobian elliptic functions $\operatorname{ns}^2 z$, $\operatorname{nc}^2 z$, $\operatorname{dn}^2 z$, and $\operatorname{nd}^2 z$, respectively, where

$$\operatorname{ns} z = \frac{1}{\operatorname{sn} z}, \quad \operatorname{nc} z = \frac{1}{\sqrt{1 - \operatorname{sn}^2 z}}, \quad \operatorname{dn} z = \sqrt{1 - k^2 \operatorname{sn}^2 z}, \quad \operatorname{nd} z = \frac{1}{\operatorname{dn} z},$$

with $w = \operatorname{sn} z$ defined by

$$z = \int_0^w \frac{dw}{\sqrt{(1-w^2)(1-k^2 w^2)}}.$$

The power series expansions of these elliptic functions give the expressions

of the corresponding q -series in terms of K , E , and k (cf. [10]). For example, we find in [17] the expressions

$$(12) \quad \begin{cases} P(q^2) := 1 - 24A_1(q) = \left(\frac{2K}{\pi}\right)^2 \left(\frac{3E}{K} - 2 + k^2\right), \\ Q(q^2) := 1 + 240A_3(q) = \left(\frac{2K}{\pi}\right)^4 (1 - k^2 + k^4), \\ R(q^2) := 1 - 504A_5(q) = \left(\frac{2K}{\pi}\right)^6 \frac{1}{2}(1 + k^2)(1 - 2k^2)(2 - k^2) \end{cases}$$

with $q = e^{-\pi c}$, $c = K'/K$.

We recall here the theorem of Nesterenko and its corollary ([16]).

NESTERENKO'S THEOREM. *If $\varrho \in \mathbb{C}$ with $0 < |\varrho| < 1$, then*

$$\text{trans.deg}_{\mathbb{Q}} \mathbb{Q}(\varrho, P(\varrho), Q(\varrho), R(\varrho)) \geq 3.$$

COROLLARY. *If $\varrho \in \overline{\mathbb{Q}}$ with $0 < |\varrho| < 1$, then $P(\varrho)$, $Q(\varrho)$, and $R(\varrho)$ are algebraically independent.*

The corollary with (12) implies the following:

LEMMA 1. *If $q = e^{-\pi c} \in \overline{\mathbb{Q}}$ with $0 < |q| < 1$, then K/π , E/π , and k are algebraically independent.*

To get the formulas stated in the theorems, we use the recurrence relations satisfied by the coefficients of the power series expansions of Jacobian elliptic functions, which are given by the following lemmas.

LEMMA 2. *The coefficients of the expansion*

$$\text{ns}^2 z = \frac{1}{z^2} + \sum_{j=0}^{\infty} c_j z^{2j}$$

are given by

$$c_0 = \frac{1}{3}(1 + k^2), \quad c_1 = \frac{1}{15}(1 - k^2 + k^4), \quad c_2 = \frac{1}{189}(1 + k^2)(1 - 2k^2)(2 - k^2),$$

$$(j - 2)(2j + 3)c_j = 3 \sum_{i=1}^{j-2} c_i c_{j-i-1} \quad (j \geq 3).$$

Proof. Since $w = \text{sn } z$ is a solution of $(w')^2 = (1 - w^2)(1 - k^2 w^2)$ such that $w(0) = 0$, the function $u = \text{ns}^2 z = w^{-2}$ satisfies $(u')^2 = 4w^{-6}(w')^2 = 4w^{-2}(w^{-2} - 1)(w^{-2} - k^2)$, that is,

$$(13) \quad (u')^2 = 4u(u - 1)(u - k^2),$$

and differentiation of (13) leads us to

$$(14) \quad u'' = 6u^2 - 4(1 + k^2)u + 2k^2.$$

Substituting $u = z^{-2} + v$, $v = \sum_{j=0}^{\infty} c_j z^{2j}$ into (14), we have

$$6z^{-4} + \sum_{j=0}^{\infty} 2j(2j-1)c_j z^{2j-2} = 6\left(\sum_{j=0}^{\infty} c_j z^{2j}\right)^2 + 12z^{-2} \sum_{j=0}^{\infty} c_j z^{2j} - 4(1+k^2) \sum_{j=0}^{\infty} c_j z^{2j} + 6z^{-4} - 4(1+k^2)z^{-2} + 2k^2.$$

Comparing the coefficients of z^{-2} and the constant terms on both sides, we obtain $c_0 = (1+k^2)/3$ and $c_1 = (1-k^2+k^4)/15$. For $j \geq 2$, the coefficients of z^{2j-2} on both sides satisfy

$$2j(2j-1)c_j = 6 \sum_{i=0}^{j-1} c_i c_{j-i-1} + 12c_j - 4(1+k^2)c_{j-1}.$$

Since $1+k^2 = 3c_0$, this can be written in the form

$$(15) \quad (j-2)(2j+3)c_j = 3 \sum_{i=0}^{j-1} c_i c_{j-i-1} - 6c_0 c_{j-1} = 3 \sum_{i=1}^{j-2} c_i c_{j-i-1}.$$

When $j = 2$, however, both sides of (15) vanish, and c_2 is not determined uniquely. To compute c_2 , substitute $u = z^{-2} + c_0 + c_1 z^2 + c_2 z^4 + \dots$ into equation (13), and compare the constant terms. Then we have $c_2 = (1+k^2)(1-2k^2)(2-k^2)/189$. Once c_0 , c_1 , and c_2 are known, the coefficients c_j ($j \geq 3$) are uniquely determined by (15). ■

It is also possible to derive the recurrence formula from the proof of Satz 3 in [11, pp. 169–170] concerning the \wp -function by choosing $ns^2z = \wp(z) + (1+k^2)/3$, $g_2 = (4/3)(1-k^2+k^4)$, and $g_3 = (4/27)(1+k^2)(1-2k^2)(2-k^2)$.

LEMMA 3. *The coefficients of the expansion*

$$(1-k^2)(nc^2z-1) = \sum_{j=1}^{\infty} c_j z^{2j}$$

are given by

$$c_1 = 1 - k^2, \quad c_2 = \frac{1}{3}(1-k^2)(2-k^2),$$

$$j(2j-1)c_1 c_j = 6c_2 c_{j-1} + 3c_1 \sum_{i=1}^{j-2} c_i c_{j-i-1} \quad (j \geq 3).$$

Proof. Since $nc^2z = (1 - sn^2z)^{-1}$, the function $u = (1-k^2)(nc^2z-1)$ is a solution of $(u')^2 = 4u(u+1)(u+1-k^2)$ such that $u(0) = u'(0) = 0$, and hence

$$u'' = 6u^2 + 4(2-k^2)u + 2(1-k^2).$$

Substitution of $u = \sum_{j=1}^{\infty} c_j z^{2j}$ yields the equality

$$\sum_{j=1}^{\infty} 2j(2j-1)c_j z^{2j-2} = 6 \left(\sum_{j=1}^{\infty} c_j z^{2j} \right)^2 + 4(2-k^2) \sum_{j=1}^{\infty} c_j z^{2j} + 2(1-k^2).$$

Looking at the constant terms, we have $c_1 = 1 - k^2$. Furthermore, for $j \geq 2$,

$$(16) \quad 2j(2j-1)c_j = 4(2-k^2)c_{j-1} + 6 \sum_{i=1}^{j-2} c_i c_{j-i-1},$$

in particular, $c_2 = (2-k^2)c_1/3 = (1-k^2)(2-k^2)/3$. Multiplying both sides of (16) by c_1 , we obtain the desired recursive relation for $j \geq 3$. ■

LEMMA 4. *The coefficients of the expansion*

$$(1-k^2) \operatorname{nd}^2 z = 1 - k^2 + \sum_{j=1}^{\infty} c_j z^{2j}$$

are given by

$$\begin{aligned} c_1 &= k^2(1-k^2), & c_2 &= -\frac{1}{3} k^2(1-k^2)(1-2k^2), \\ j(2j-1)c_1 c_j &= 6c_2 c_{j-1} - 3c_1 \sum_{i=1}^{j-2} c_i c_{j-i-1} \quad (j \geq 3). \end{aligned}$$

Proof. The function $u = (1-k^2) \operatorname{nd}^2 z = (1-k^2)(1-k^2 \operatorname{sn}^2 z)^{-1}$, which satisfies $(u')^2 = 4u(1-u)(u - (1-k^2))$, is a solution of

$$u'' = -6u^2 + 4(2-k^2)u - 2(1-k^2)$$

with the initial condition $u(0) = 1 - k^2$, $u'(0) = 0$. Substitution of $u = 1 - k^2 + v$, $v = \sum_{j=1}^{\infty} c_j z^{2j}$ yields $c_1 = k^2(1-k^2)$ and, for $j \geq 2$,

$$j(2j-1)c_j = -2(1-2k^2)c_{j-1} - 3 \sum_{i=1}^{j-2} c_i c_{j-i-1}.$$

Multiplying both sides by c_1 and observing $c_2 = -(1-2k^2)c_1/3$, we obtain the desired formula for $j \geq 3$. ■

LEMMA 5. *The coefficients of the expansion*

$$\operatorname{dn}^2 z = 1 + \sum_{j=1}^{\infty} c_j z^{2j}$$

are given by

$$c_1 = -k^2, \quad c_2 = \frac{1}{3}k^2(1+k^2),$$

$$j(2j-1)c_1c_j = 6c_2c_{j-1} - 3c_1 \sum_{i=1}^{j-2} c_i c_{j-i-1} \quad (j \geq 3).$$

Proof. By definition, $u = \operatorname{dn}^2 z$ satisfies $(u')^2 = 4u(1-u)(u - (1-k^2))$, $u(0) = 1$. This lemma is verified by the same argument as in the proof of Lemma 4. ■

4. Proofs of the theorems

4.1. Proof of Theorem 1. It follows from (4), (6), and (8) that

$$(17) \quad \Phi_{2s} = \frac{1}{(2s-1)!} \left(\sigma_{s-1}(s)\mu_s + \sum_{j=1}^{s-1} \sigma_{s-j-1}(s)(A_{2j+1} - (-1)^s D_{2j+1}) \right),$$

where $\mu_s = (A_1 - (-1)^s D_1)$. In particular

$$(18) \quad \Phi_2 = A_1 + D_1, \quad 6\Phi_4 = (A_3 - D_3) - (A_1 - D_1),$$

$$(19) \quad 120\Phi_6 = (A_5 + D_5) - 5(A_3 + D_3) + 4(A_1 + D_1).$$

Recall the generating functions of A_{2j+1} and D_{2j+1} given by

$$\left(\frac{2K}{\pi}\right)^2 \operatorname{ns}^2\left(\frac{2Kx}{\pi}\right) = \frac{4K(K-E)}{\pi^2} + \operatorname{cosec}^2 x - 8 \sum_{j=0}^{\infty} (-1)^j A_{2j+1} \frac{(2x)^{2j}}{(2j)!},$$

$$\left(\frac{2K}{\pi}\right)^2 (1-k^2) \operatorname{nd}^2\left(\frac{2Kx}{\pi}\right) = \frac{4KE}{\pi^2} - 8 \sum_{j=0}^{\infty} (-1)^j D_{2j+1} \frac{(2x)^{2j}}{(2j)!}$$

(cf. [19, Tables 1(i), 1(iv)], [12], [18, p. 535], [1]), which yield

$$(20) \quad \varphi_j := \left(\frac{2K}{\pi}\right)^{2j+2} c_j = a_j - (-1)^j \frac{2^{2j+3}}{(2j)!} A_{2j+1} \quad (j \geq 1),$$

$$(21) \quad \psi_j := \left(\frac{2K}{\pi}\right)^{2j+2} c'_j = (-1)^{j+1} \frac{2^{2j+3}}{(2j)!} D_{2j+1} \quad (j \geq 1).$$

Here c_j and c'_j are the coefficients given by Lemmas 2 and 4, respectively. Note that

$$(22) \quad \begin{cases} 8D_1 = \left(\frac{2K}{\pi}\right)^2 \left(\frac{E}{K} + k^2 - 1\right), \\ 16D_3 = \left(\frac{2K}{\pi}\right)^4 k^2(1-k^2), \\ 16D_5 = \left(\frac{2K}{\pi}\right)^6 k^2(1-k^2)(1-2k^2) \end{cases}$$

(cf. [19, Table 1(iv)] or (21) combined with Lemma 4). Then, by (12) with $q = \beta^2$ and (22),

$$(23) \quad A_1 + D_1 = \frac{1}{24} \left(1 - \left(\frac{2K}{\pi} \right)^2 (1 - 2k^2) \right),$$

$$(24) \quad A_1 - D_1 = \frac{1}{24} \left(1 - \left(\frac{2K}{\pi} \right)^2 \left(\frac{6E}{K} - 5 + 4k^2 \right) \right),$$

$$(25) \quad A_3 + D_3 = -\frac{1}{240} \left(1 - \left(\frac{2K}{\pi} \right)^4 (1 + 14k^2 - 14k^4) \right),$$

$$(26) \quad A_3 - D_3 = -\frac{1}{240} \left(1 - \left(\frac{2K}{\pi} \right)^4 (1 - 16k^2 + 16k^4) \right),$$

$$(27) \quad A_5 + D_5 = \frac{1}{504} \left(1 - \left(\frac{2K}{\pi} \right)^6 (1 - 2k^2)(1 - 31k^2 + 31k^4) \right),$$

$$(28) \quad A_5 - D_5 = \frac{1}{504} \left(1 - \left(\frac{2K}{\pi} \right)^6 (1 - 2k^2)(1 + 32k^2 - 32k^4) \right).$$

Substituting these quantities into (18) and (19), we have expressions of Φ_2 , Φ_4 , and Φ_6 in terms of K , E , and k ; and then we put

$$(29) \quad \begin{aligned} X &:= 24\Phi_2 - 1 = -\left(\frac{2K}{\pi} \right)^2 (1 - 2k^2), \\ Y &:= 1440\Phi_4 + 11 \\ &= 10 \left(\frac{2K}{\pi} \right)^2 \left(\frac{6E}{K} - 5 + 4k^2 \right) + \left(\frac{2K}{\pi} \right)^4 (1 - 16k^2 + 16k^4), \\ Z &:= -120960\Phi_6 + 4032\Phi_2 + 23 \\ &= 2 \left(\frac{2K}{\pi} \right)^6 (1 - 2k^2)(1 - 31k^2 + 31k^4) \\ &\quad + 21 \left(\frac{2K}{\pi} \right)^4 (1 + 14k^2 - 14k^4). \end{aligned}$$

Let us set

$$(30) \quad x := \left(\frac{2K}{\pi} \right)^2, \quad y := \left(\frac{2K}{\pi} \right)^2 \frac{E}{K}, \quad z := \left(\frac{2K}{\pi} \right)^2 (1 - 2k^2).$$

(Note that the same symbols x , y , and z denote different quantities in Section 2.) Then the quantities above are written as follows:

$$\begin{aligned} X &= -z, \quad Y = 10(6y - 3x - 2z) + 4z^2 - 3x^2, \\ Z &= \frac{z}{2} (31z^2 - 27x^2) + \frac{21}{2} (9x^2 - 7z^2). \end{aligned}$$

Solving these equations, we obtain

$$x^2 = \Omega := \frac{31X^3 + 147X^2 + 2Z}{27(X+7)},$$

$$30(2y - x) = Y - 4X^2 - 20X + 3\Omega, \quad z = -X.$$

Hence x, y, z are algebraic over $\mathbb{Q}(X, Y, Z) = \mathbb{Q}(\Phi_2, \Phi_4, \Phi_6)$, and the algebraic independence of Φ_2, Φ_4, Φ_6 follows from Lemma 1.

Substituting (20) and (21) into (17), we obtain the formula for Φ_{2s} stated in the theorem. Observe that

$$\frac{\Omega}{16} = \frac{124}{3} \Phi_2^2 - \frac{22}{3} \Phi_2 + \frac{103}{48} - \frac{5}{3} \omega \quad \text{with} \quad \omega = \frac{56\Phi_6 + 5/4}{4\Phi_2 + 1}.$$

The expression of μ_s is obtained by writing (23) and (24) in terms of X, Y , and Ω , namely, Φ_2, Φ_4 , and ω . Furthermore, (20) and (21) for $j = 1, 2$ together with (12) and (22) imply

$$(31) \quad \begin{cases} \varphi_1 = \frac{1}{15} \left(\frac{2K}{\pi} \right)^4 (1 - k^2 + k^4) = \frac{1}{60} (3x^2 + z^2), \\ \varphi_2 = \frac{1}{189} \left(\frac{2K}{\pi} \right)^6 (1 - 2k^2)(2 + k^2 - k^4) = \frac{z}{756} (9x^2 - z^2), \\ \psi_1 = \left(\frac{2K}{\pi} \right)^4 (k^2 - k^4) = \frac{1}{4} (x^2 - z^2), \\ \psi_2 = -\frac{1}{3} \left(\frac{2K}{\pi} \right)^2 (1 - 2k^2) \psi_1 = -\frac{z}{3} \psi_1 \end{cases}$$

with $x^2 = \Omega, z = -X$, from which the desired expressions follow. Multiplying both sides of the recurrence formula in Lemma 2 by $(2K/\pi)^{2j+2}$, and using (20), we obtain the formula for φ_j ($j \geq 3$). The recurrence formula for ψ_j is obtained from Lemma 4 by using $\psi_2/\psi_1 = X/3 = (24\Phi_2 - 1)/3$.

4.2. Proof of Remark 1. Denoting the denominator of ω by $T := 4\Phi_2 + 1$, we have the polynomials $T^\nu \varphi_{2\nu-1}, T^\nu \varphi_{2\nu}, T^\nu \psi_{2\nu-1}, T^\nu \psi_{2\nu}$ ($\nu \geq 1$) in (Φ_2, Φ_6) of total degrees, say, $\delta_{2\nu-1}, \delta_{2\nu}, \delta'_{2\nu-1}, \delta'_{2\nu}$, respectively. Then the recurrence formulas concerning φ_j and ψ_j imply, for $\nu \geq 1$,

$$\begin{aligned} \delta_{2\nu-1} &\leq \max_{1 \leq i \leq 2\nu-3} \{1 + \delta_{2i} + \delta_{2(\nu-i-1)}, \delta_{2i-1} + \delta_{2(\nu-i)-1}\}, \\ \delta_{2\nu} &\leq \max_{1 \leq i \leq 2\nu-2} \{\delta_{2i} + \delta_{2(\nu-i)-1}\}, \\ \delta'_{2\nu-1} &\leq \max_{1 \leq i \leq 2\nu-3} \{1 + \delta'_{2i} + \delta'_{2(\nu-i-1)}, \delta'_{2i-1} + \delta'_{2(\nu-i)-1}, 2 + \delta'_{2\nu-2}\}, \\ \delta'_{2\nu} &\leq \max_{1 \leq i \leq 2\nu-2} \{\delta'_{2i} + \delta'_{2(\nu-i)-1}, 1 + \delta'_{2\nu-1}\}. \end{aligned}$$

Using these relations, by induction on ν we can verify $\delta_{2\nu-1}, \delta'_{2\nu-1} \leq 3\nu$ and $\delta_{2\nu}, \delta'_{2\nu} \leq 3\nu + 1$, which yields the estimate for the total degree of the polynomial $T^{\lfloor s/2 \rfloor} (\Phi_{2s} - r_s \Phi_4)$.

4.3. Proof of Theorem 2. From (4), (6), and (10) we have

$$(32) \quad \Phi_{2s}^* = \frac{1}{(2s-1)!} \times \left(\sigma_{s-1}(s)\mu_s - \sum_{j=1}^{s-1} \sigma_{s-j-1}(s)((-1)^s D_{2j+1}(\beta^2) + A_{2j+1}(\beta^2)) \right)$$

with $\mu_s = -((-1)^s D_1(\beta^2) + A_1(\beta^2))$. In particular,

$$(33) \quad \Phi_2^* = D_1 - A_1, \quad 6\Phi_4^* = -(D_3 + A_3) + (D_1 + A_1),$$

$$(34) \quad 120\Phi_6^* = (D_5 - A_5) - 5(D_3 - A_3) + 4(D_1 - A_1).$$

Substituting (23)–(26) and (28) into (33) and (34), we obtain

$$\begin{aligned} Y &:= 1440\Phi_4^* - 11 = -10 \left(\frac{2K}{\pi} \right)^2 (1 - 2k^2) - \left(\frac{2K}{\pi} \right)^4 (1 + 14k^2 - 14k^4), \\ Z &:= 120960\Phi_6^* - 4032\Phi_2^* + 23 \\ &= 21 \left(\frac{2K}{\pi} \right)^4 (1 - 16k^2 + 16k^4) + 2 \left(\frac{2K}{\pi} \right)^6 (1 - 2k^2)(1 + 32k^2 - 32k^4), \end{aligned}$$

which are written in the form

$$Y = -10z - (9x^2 - 7z^2)/2, \quad Z = 21(4z^2 - 3x^2) + 2z(9x^2 - 8z^2)$$

with x and z given by (30). From these equations, we obtain

$$8\xi^3 + 5\xi^2 + (Y - 35)\xi - (14Y - Z)/16 = 0, \quad \xi = z/4,$$

which is written as (3). It is easy to see that K/π , E/π , and k are algebraic functions of Φ_2^* , Φ_4^* , Φ_6^* over \mathbb{Q} . This fact combined with Lemma 1 implies the algebraic independence of Φ_2^* , Φ_4^* , Φ_6^* .

Formula (32) together with (20) and (21) implies the expression of Φ_{2s}^* in terms of φ_j and ψ_j . By (23) and (24), we can write μ_s in terms of Φ_2^* and ξ . Substituting $z = 4\xi$ and $9x^2 = 16(7\xi^2 - 5\xi) - 2Y$ into (31), we obtain φ_1, \dots, ψ_2 expressed in terms of Φ_4^* and ξ . The recurrence formula for ψ_j follows from that of Lemma 4 combined with (21) and the relation $\psi_2/\psi_1 = -4\xi/3$.

4.4. Proof of Theorem 3. From (5), (7), and (9) we have

$$(35) \quad \Psi_{2s} = \frac{1}{(2s-1)!} \left(\sigma_{s-1}(s)\mu_s + \sum_{j=1}^{s-1} \sigma_{s-j-1}(s)(C_{2j+1} - (-1)^s B_{2j+1}) \right),$$

where $\mu_s = (C_1 - (-1)^s B_1)$. In particular,

$$\Psi_2 = C_1 + B_1, \quad 6\Psi_4 = (C_3 - B_3) - (C_1 - B_1),$$

$$120\Psi_6 = (C_5 + B_5) - 5(C_3 + B_3) + 4(C_1 + B_1).$$

The formulas

$$\begin{aligned} \left(\frac{2K}{\pi}\right)^2 (1-k^2) \left(\operatorname{nc}^2\left(\frac{2Kx}{\pi}\right) - 1\right) \\ = -\frac{4KE}{\pi^2} + \sec^2 x + 8 \sum_{j=0}^{\infty} (-1)^j B_{2j+1} \frac{(2x)^{2j}}{(2j)!}, \\ \left(\frac{2K}{\pi}\right)^2 \operatorname{dn}^2\left(\frac{2Kx}{\pi}\right) = \frac{4KE}{\pi^2} + 8 \sum_{j=0}^{\infty} (-1)^j C_{2j+1} \frac{(2x)^{2j}}{(2j)!} \end{aligned}$$

(cf. [19, Tables 1(ii), 1(iii)]) imply the relations

$$(36) \quad \varphi_j := \left(\frac{2K}{\pi}\right)^{2j+2} c_j = b_j + (-1)^j \frac{2^{2j+3}}{(2j)!} B_{2j+1} \quad (j \geq 1),$$

$$(37) \quad \psi_j := \left(\frac{2K}{\pi}\right)^{2j+2} c'_j = (-1)^j \frac{2^{2j+3}}{(2j)!} C_{2j+1} \quad (j \geq 1),$$

where c_j and c'_j are the coefficients given in Lemmas 3 and 5, respectively. From [19, Tables 1(ii), 1(iii)] or Lemmas 3 and 5, we obtain

$$\begin{aligned} 1 + 8B_1 &= \frac{4KE}{\pi^2}, & 8C_1 &= \left(\frac{2K}{\pi}\right)^2 \left(1 - \frac{E}{K}\right), \\ 1 - 16B_3 &= \left(\frac{2K}{\pi}\right)^4 (1 - k^2), & 16C_3 &= \left(\frac{2K}{\pi}\right)^4 k^2, \\ 1 + 8B_5 &= \left(\frac{2K}{\pi}\right)^6 \frac{1}{2} (1 - k^2)(2 - k^2), & 16C_5 &= \left(\frac{2K}{\pi}\right)^6 k^2(1 + k^2). \end{aligned}$$

Hence

$$(38) \quad C_1 + B_1 = -\frac{1}{8} + \frac{1}{8} \left(\frac{2K}{\pi}\right)^2,$$

$$(39) \quad C_1 - B_1 = \frac{1}{8} + \frac{1}{8} \left(\frac{2K}{\pi}\right)^2 \left(1 - \frac{2E}{K}\right),$$

$$(40) \quad C_3 + B_3 = \frac{1}{16} - \frac{1}{16} \left(\frac{2K}{\pi}\right)^4 (1 - 2k^2),$$

$$(41) \quad C_3 - B_3 = -\frac{1}{16} + \frac{1}{16} \left(\frac{2K}{\pi}\right)^4,$$

$$(42) \quad C_5 + B_5 = -\frac{1}{8} + \frac{1}{8} \left(\frac{2K}{\pi}\right)^6 (1 - k^2 + k^4),$$

$$(43) \quad C_5 - B_5 = \frac{1}{8} - \frac{1}{8} \left(\frac{2K}{\pi}\right)^6 (1 - 2k^2).$$

So we can express Ψ_2 , Ψ_4 , and Ψ_6 in terms of K , E , and k . Put

$$(44) \quad \begin{aligned} X &:= 8\Psi_2 + 1 = \left(\frac{2K}{\pi}\right)^2, \\ Y &:= 96\Psi_4 + 3 = \left(\frac{2K}{\pi}\right)^4 - 2\left(\frac{2K}{\pi}\right)^2\left(1 - \frac{2E}{K}\right), \\ Z &:= 1920\Psi_6 + 15 \\ &= 2\left(\frac{2K}{\pi}\right)^6(1 - k^2 + k^4) + 5\left(\frac{2K}{\pi}\right)^4(1 - 2k^2) + 8\left(\frac{2K}{\pi}\right)^2. \end{aligned}$$

By the same argument as in Sections 4.1 and 4.3 with x , y , z given by (30), we have

$$(45) \quad \begin{aligned} x &= X, & 4y &= Y - X^2 + 2X, \\ Xz^2 + 10Xz + (3X^3 + 16X - 2Z) &= 0, \end{aligned}$$

implying that x , y , and z are algebraic over $\mathbb{Q}(X, Y, Z) = \mathbb{Q}(\Psi_2, \Psi_4, \Psi_6)$.

Substitution of (36) and (37) into (35) yields the desired formula for Ψ_{2s} in terms of φ_j and ψ_j . Furthermore, using (36) and (37) for $j = 1, 2$ together with Lemmas 3 and 5, respectively, we have

$$(46) \quad \left\{ \begin{aligned} \varphi_1 &= \left(\frac{2K}{\pi}\right)^4(1 - k^2) = \frac{x}{2}(z + x), \\ \varphi_2 &= \frac{1}{3}\left(\frac{2K}{\pi}\right)^2(2 - k^2)\varphi_1 = \frac{1}{6}(z + 3x)\varphi_1, \\ \psi_1 &= -\left(\frac{2K}{\pi}\right)^4k^2 = \frac{x}{2}(z - x), \\ \psi_2 &= -\frac{1}{3}\left(\frac{2K}{\pi}\right)^2(1 + k^2)\psi_1 = \frac{1}{6}(z - 3x)\psi_1. \end{aligned} \right.$$

The right hand sides of these equalities as well as μ_s can be written in terms of X and z satisfying the quadratic equation (45). The recurrence formulas for φ_j and ψ_j are derived from those of Lemmas 3 and 5 combined with $\varphi_2/\varphi_1 = (3X + z)/6$ and $\psi_2/\psi_1 = -(3X - z)/6$, respectively. Replacing z by η , we obtain the theorem.

4.5. Proof of Theorem 4. Relation (11) together with (5) and (7) implies

$$(47) \quad \begin{aligned} \Psi_{2s}^* &= \frac{1}{(2s-1)!} \\ &\times \left(\sigma_{s-1}(s)\mu_s + \sum_{j=1}^{s-1} \sigma_{s-j-1}(s)(C_{2j+1}(\beta^2) + (-1)^s B_{2j+1}(\beta^2)) \right) \end{aligned}$$

with $\mu_s = C_1(\beta^2) + (-1)^s B_1(\beta^2)$. In particular,

$$(48) \quad \Psi_2^* = C_1 - B_1, \quad 6\Psi_4^* = (C_3 + B_3) - (C_1 + B_1),$$

$$(49) \quad 120\Psi_6^* = (C_5 - B_5) - 5(C_3 - B_3) + 4(C_1 - B_1).$$

Substituting (38)–(41) and (43) into (48) and (49), we obtain

$$Y := 96\Psi_4^* - 3 = -\left(\frac{2K}{\pi}\right)^4 (1 - 2k^2) - 2\left(\frac{2K}{\pi}\right)^2,$$

$$Z := 1920\Psi_6^* - 64\Psi_2^* - 7 = -2\left(\frac{2K}{\pi}\right)^6 (1 - 2k^2) - 5\left(\frac{2K}{\pi}\right)^4,$$

which are written as $Y = -xz - 2x$, $Z = -2x^2z - 5x^2$ with x and z given by (30). Then we have $x^2 - 2Yx + Z = 0$, $z = -2 - Y/x$. The quantities K/π , E/π , and k are algebraic functions of Ψ_2^* , Ψ_4^* , Ψ_6^* over \mathbb{Q} , which implies the algebraic independence of Ψ_2^* , Ψ_4^* , Ψ_6^* .

The desired expression of $\Psi_{2^s}^*$ is obtained by substituting (36) and (37) into (47). By (38) and (39) the quantity μ_s is expressible in terms of Ψ_2^* and x . Using (46), we have the desired expressions of φ_1, \dots, ψ_2 in terms of Ψ_4^* and x . Observing the relations $\varphi_2/\varphi_1 = (3x + z)/6$ and $\psi_2/\psi_1 = -(3x - z)/6$, we derive from Lemmas 3 and 5 the recurrence formulas for φ_j and ψ_j , respectively. Replacing x by θ in these expressions, we obtain the theorem.

5. Branches of the related algebraic functions. In what follows we put $a = \alpha + \beta \in \mathbb{C}$, namely $\beta = (a/2)(1 - \sqrt{1 + 4a^{-2}})$, where the branch is chosen so that $\beta(a) = O(a^{-1})$ as $a \rightarrow \infty$. Then each reciprocal sum in the theorems is treated as a function of a or β .

THEOREM 5. *Under the assumptions of Theorem 2, we have the following:*

- (i) *The function $\xi = \xi(a)$ is holomorphic in the domain $|a| > 8.146$, and is expressible in the form*

$$(50) \quad \xi(a) = -\frac{5}{24} + \varrho_0 + \gamma_1/\varrho_0, \quad \varrho_0 = \varrho_0(a) := \sqrt[3]{\gamma_2 + \sqrt{\gamma_1^3 - \gamma_2^2}} i$$

with $\gamma_1 = \gamma_1(a)$, $\gamma_2 = \gamma_2(a)$ expressed as

$$(51) \quad \gamma_1 := \frac{1129}{24^2} - 60\Phi_4^*, \quad \gamma_2 := -\frac{17963}{24^3} + \frac{63}{4}\Phi_2^* + \frac{195}{2}\Phi_4^* - \frac{945}{2}\Phi_6^*$$

satisfying $\gamma_1(a) = 1129 \cdot 24^{-2} + O(a^{-2})$ and $\gamma_2(a) = -17963 \cdot 24^{-3} + O(a^{-2})$ as $a \rightarrow \infty$. Here the branches of the square and cube roots are taken so that $\sqrt{\gamma_1(\infty)^3 - \gamma_2(\infty)^2} = (7/192)\sqrt{4395}$ and that

$$(52) \quad \varrho_0(\infty) = \frac{1}{48} (11 - \sqrt{4395} i).$$

- (ii) For any real number $a \geq 2.6$, the function $\xi(a)$ admits the expression (50) with the same branches as above, where

$$\gamma_1^3 - \gamma_2^2 > 0, \quad \gamma_2 + \sqrt{\gamma_1^3 - \gamma_2^2} i \neq 0.$$

- (iii) For $a = 1$ (respectively, $a = 2$) corresponding to the Fibonacci (respectively, Pell) numbers, the quantity $\xi(1) = -2.66158\dots$ (respectively, $\xi(2) = -0.72378\dots$) is expressible as (50) with

$$\varrho_0(1) = -1.22662\dots + (0.57598\dots)i, \quad \gamma_1(1)^3 - \gamma_2(1)^2 > 0$$

(respectively,

$$\varrho_0(2) = -0.25772\dots - (1.00662\dots)i, \quad \gamma_1(2)^3 - \gamma_2(2)^2 > 0).$$

THEOREM 6. Under the assumptions of Theorem 3, we have the following:

- (i) The function $\eta = \eta(a)$ is holomorphic in the domain $|a| > 5.431$, and is expressible in the form

$$\eta(a) = -5 + \sqrt{\chi(a)}$$

with

$$\chi(a) = -192\Psi_2^2 - 48\Psi_2 + 6 + (3840\Psi_6 + 30)/(8\Psi_2 + 1)$$

satisfying $\chi(a) = 36 + O(a^{-2})$ as $a \rightarrow \infty$. Here the branch is taken so that $\sqrt{\chi(\infty)} = 6$.

- (ii) For $a = 1$ corresponding to the Lucas numbers,

$$\eta(1) = -5 - \sqrt{\chi(1)} \quad (< -5),$$

and for any integer $a \geq 2$

$$(53) \quad \eta(a) = -5 + \sqrt{\chi(a)} \quad (> -5).$$

Furthermore, the last equality holds for any real number $a \geq 2.4$.

THEOREM 7. Under the assumptions of Theorem 4, we have the following:

- (i) The function $\theta = \theta(a)$ is holomorphic in the domain $|a| > 5.819$, and is expressible in the form

$$\theta(a) = \chi_1(a) + \sqrt{\chi_2(a)}$$

with

$$\chi_1(a) = 96\Psi_4^* - 3,$$

$$\chi_2(a) = 64\Psi_2^* + 9216(\Psi_4^*)^2 - 576\Psi_4^* - 1920\Psi_6^* + 16$$

satisfying $\chi_1(a) = -3 + O(a^{-4})$ and $\chi_2(a) = 16 + O(a^{-2})$ as $a \rightarrow \infty$. Here the branch of the square root is taken so that $\sqrt{\chi_2(\infty)} = 4$.

- (ii) For $a = 1$ corresponding to the Lucas numbers,

$$\theta(1) = \chi_1(1) - \sqrt{\chi_2(1)} \quad (< \chi_1(1)),$$

and for any integer $a \geq 2$,

$$\theta(a) = \chi_1(a) + \sqrt{\chi_2(a)} \quad (> \chi_1(a)).$$

Furthermore, the last equality holds for any real number $a \geq 2.5$.

5.1. Proof of Theorem 6. By definition we have $|V_n| = |\beta^{-n} + \beta^n| \geq |\beta|^{-n}(1 - |\beta|^{2n})$, so that $|V_n|^{-1} \leq (1 - |\beta|^2)^{-1}|\beta|^n$. This implies, for $s \geq 1$,

$$(54) \quad |\Psi_{2s}| \leq \sum_{n \geq 1} |V_n|^{-2s} \leq (1 - |\beta|^2)^{-2s} \sum_{n \geq 1} |\beta|^{2sn} \leq (1 - |\beta|^2)^{-2s-1} |\beta|^{2s},$$

and hence Ψ_{2s} ($s \in \mathbb{N}$) are holomorphic for $|\beta| < 1$. Set

$$(55) \quad \Omega(\beta) = (8\Psi_2 + 1)\chi(a) = 12(3 - 128\Psi_2^3 - 48\Psi_2^2 + 320\Psi_6).$$

Under the assumption $|\beta| < 0.3$, we deduce from (54) that $|\Psi_2| \leq 1.1^3|\beta|^2 \leq 0.12$, and that $|\Psi_6| \leq 1.1^7|\beta|^6 \leq 0.002$. By these inequalities, for $|\beta| < 0.3$ we get $|\Omega(\beta)|/12 \geq 3 - 128|\Psi_2|^3 - 48|\Psi_2|^2 - 320|\Psi_6| > 1$, and $|8\Psi_2 + 1| \geq 1 - 8|\Psi_2| > 0.03$. Hence, in the domain $|\beta| < 0.3$, the function $\sqrt{\chi(a)} \neq 0, \infty$ is holomorphic. Note that $\beta = \beta(a)$ conformally maps the domain $\mathbb{C} \setminus \{a = yi \mid -2 \leq y \leq 2\}$ to the disc $|\beta| < 1$. To examine the corresponding domain around $a = \infty$, observe that for $|z| < 1/7$,

$$|\sqrt{1+z} - 1| \leq \left| \int_0^z \frac{2}{3} (1+t)^{3/2} dt \right| \leq \frac{2}{3} \left(\frac{8}{7} \right)^{3/2} |z| \leq 0.8146|z|.$$

Using this estimate, we have $|2\beta| = |a(1 - \sqrt{1 + 4a^{-2}})| \leq 0.8146|4a^{-1}| = 3.2584|a|^{-1}$, provided that $|4a^{-2}| < 1/7$. This fact implies that the image of the domain $|a| > 5.431$ is contained in the disc $|\beta(a)| < 0.3$. By (54) and (55) we see

$$(56) \quad \chi(a) = 36 + O(\beta^2) = 36 + O(a^{-2}) \quad \text{as } a \rightarrow \infty.$$

To determine the branch of η , recall that

$$(57) \quad \eta = (2K/\pi)^2(1 - 2k^2) = 1 - 24\Phi_2$$

(cf. Section 4.4 and (29)). Since $|\alpha - \beta| |U_n| \geq |\beta|^{-n}(1 - |\beta|^2)$, we have $\Phi_2 = O(\beta^2) = O(a^{-2})$, so that $\eta(a) = 1 + O(a^{-2})$ as $a \rightarrow \infty$. This fact combined with (56) implies the assertion (i).

To determine the sign of $\eta(1)$, we use the following numerical values for $a = 1$:

$$\Psi_2 = 1.2072919\dots, \quad \Psi_6 = 1.0016249\dots, \quad \Phi_2 = 0.485264\dots$$

Then we have $\chi(1) = 31.88115\dots = (5.6463\dots)^2$. On the other hand, $\eta(1) = -10.64633\dots$ by (57). Consideration of these values leads to the sign $-$ of the square root. Similarly, for $a = 2$, using the values

$$\Psi_2 = 0.2839243\dots, \quad \Psi_6 = 0.0156465\dots, \quad \Phi_2 = 0.1622974\dots,$$

we get $\eta(2) = -2.89513\dots$ and $\chi(2) = (2.1048\dots)^2$, which yields the sign + in (53). Finally, we consider the case $a \geq 2.4$. Since $V_2 = a^2 + 2 = a(a + 2a^{-1}) \geq 3.23a$, we have $V_n \geq 3.23a^{n-1}$ ($n \geq 2$), which yields $\Psi_2 = \sum_{n \geq 1} V_n^{-2} \leq a^{-2} + 3.23^{-2}a^{-2} \sum_{j \geq 0} a^{-2j} < 1.116a^{-2} < 0.194$ for $a \geq 2.4$. Hence (55) implies $\chi(a) > 1$ for $a \geq 2.4$. Combining this fact with (56) and $\eta(a) = 1 + O(a^{-2})$ ($a \rightarrow \infty$), and taking the continuity of $\eta(a)$ into account, we deduce formula (53) for $a \geq 2.4$, which completes the proof.

5.2. Proof of Theorem 7. Since

$$|\Psi_{2s}^*| \leq \sum_{n \geq 1} |V_n|^{-2s} \leq (1 - |\beta|^2)^{-2s-1} |\beta|^{2s},$$

we have $|\Psi_{2s}^*| \leq 1.1^{2s+1} |\beta|^{2s}$ for $|\beta| < 0.3$, so that

$$|\chi_2(a)| \geq 16 - 64|\Psi_2^*| - 9216|\Psi_4^*|^2 - 576|\Psi_4^*| - 1920|\Psi_6^*| > 0.7$$

for $|\beta| < 0.28$. This implies that $\sqrt{\chi_2(a)}$ is holomorphic for $|\beta| < 0.28$. By the same argument as in Section 5.1, it is holomorphic for $|a| > 5.819$. Recall the relation

$$(58) \quad \theta = (2K/\pi)^2 = 8\Psi_2 + 1$$

(cf. (44)). Observing the estimates $\Psi_2 = O(a^{-2})$, $\chi_1(a) = -3 + O(a^{-4})$, and $\chi_2(a) = 16 + O(a^{-2})$ as $a \rightarrow \infty$, we have $\theta(a) = \chi_1(a) + \sqrt{\chi_2(a)}$ for $|a| > 5.819$, where the branch is taken so that $\sqrt{\chi_2(\infty)} = 4$. In this way the assertion (i) is verified.

When $a = 1$, by numerical computation we have

$$\Psi_2^* = 0.9370204\dots, \quad \Psi_4^* = 0.9912040\dots, \quad \Psi_6^* = 0.9988644\dots;$$

and relation (58) together with the numerical value of Ψ_2 (cf. Section 5.1) yields $\theta(1) = 10.65833\dots$. Similarly when $a = 2$,

$$\Psi_2^* = 0.2265861\dots, \quad \Psi_4^* = 0.0617537\dots, \quad \Psi_6^* = 0.0156036\dots,$$

and $\theta(2) = 3.27139\dots$. Using these numerical values, we can check the equalities for $a = 1, 2$. For $a \geq 2.5$, observe that $V_n \geq 3.3a^{n-1}$ ($n \geq 2$). Then for $a \geq 2.5$ we have $\Psi_6^* \leq \sum_{m \geq 1} V_{2m-1}^{-6} \leq a^{-6} + 3.3^{-6}a^{-12} \sum_{j \geq 0} a^{-12j} \leq 1.00001a^{-6} < 0.03a^{-2}$ and $\Psi_2^* \geq V_1^{-2} - V_2^{-2} \geq (a^2 + 2)^{-1} > 0.75a^{-2}$, so that $\chi_2(a) \geq 7 + 64\Psi_2^* - 1920\Psi_6^* > 7 - 9.6a^{-2} > 5$, which implies the assertion.

5.3. Proof of Theorem 5. Put $\xi = -5/24 + \lambda$. Then (3) is written in the form

$$(59) \quad \lambda^3 - 3\gamma_1\lambda - 2\gamma_2 = 0,$$

where γ_1 and γ_2 are as given in the theorem, which may be regarded as functions of $a \in \mathbb{C}$ or β . Since $|\alpha - \beta| |U_n| \geq |\beta|^{-n} (1 - |\beta|^2)$, we infer for $s \geq 1$ that $|\Phi_{2s}^*| \leq (1 - |\beta|^2)^{-2s-1} |\beta|^{2s}$. Hence Φ_{2s}^* ($s \in \mathbb{N}$) are holomorphic for $|\beta| < 1$. In particular, for $|\beta| < 0.2$, we have $|\Phi_2^*| < 0.046$, $|\Phi_4^*| < 0.002$

and $|\Phi_6^*| < 0.0001$. Using (51), we obtain $|\gamma_1| > 1.8$, $|\gamma_2| < 2.3$ for $|\beta| < 0.2$, and $\gamma_1 \rightarrow 1129 \cdot 24^{-2}$, $\gamma_2 \rightarrow -17963 \cdot 24^{-3}$ as $\beta \rightarrow 0$ (or $a \rightarrow \infty$).

To determine the branch of ξ , we note that $\xi(a) = \eta(a)/4$ (cf. (57)), which implies $\xi \rightarrow 1/4$ as $a \rightarrow \infty$. Let us find the solution of (59) corresponding to ξ . Such a solution may be expressed in the form

$$\lambda_0 = \gamma_1/\varrho_0 + \varrho_0, \quad \varrho_0 = \sqrt[3]{\gamma_2 + \sqrt{\gamma_1^3 - \gamma_2^2}i}$$

around $\beta = 0$, where the branches should be chosen so that $24^6(\gamma_1^3 - \gamma_2^2) \rightarrow 1129^3 - 17963^2 > 0$, $\xi \rightarrow 1/4$, namely, $\lambda_0 \rightarrow 11/24$ as $\beta \rightarrow 0$. The second condition is satisfied if (52) is valid. Furthermore, provided $|\beta| < 0.2$, it is easy to see that $|\gamma_1^3 - \gamma_2^2| \geq |\gamma_1^3| - |\gamma_2^2| > 1.8^3 - 2.3^2 > 0$, and that $\gamma_2 + \sqrt{\gamma_1^3 - \gamma_2^2}i \neq 0$. Hence $\varrho_0(\beta)$ is holomorphic for $|\beta| < 0.2$. By the same numerical argument as in Section 5.1, we can verify that the image of the domain $|a| > 8.146$ is contained in the disc $|\beta| < 0.2$. In this way the assertion (i) is proved.

If $a \geq 2.6$, then $\Phi_{2s}^* \leq 10.76^{-s}$, because $(\alpha - \beta)^2 = a^2 + 4 \geq 10.76$. Hence we have $\gamma_1 \geq 1.44$ and $-1.68 \leq \gamma_2 \leq 1.01$, implying $\gamma_1^3 - \gamma_2^2 > 0.16$ for $a \geq 2.6$. This fact combined with the continuity of ξ with respect to a implies the assertion (ii).

When $a = 1$, by numerical computation, we obtain

$$\Phi_2^* = 0.0335078\dots, \quad \Phi_4^* = 0.0020616\dots, \quad \Phi_6^* = 0.0001145\dots$$

Then we have $\gamma_1(1) = 1.836371\dots$, $\gamma_2(1) = -0.624754\dots$, and $\xi(1) = \eta(1)/4 = -2.66158\dots$ (cf. Section 5.1). Similarly, when $a = 2$, we have

$$\Phi_2^* = 0.0980088\dots, \quad \Phi_4^* = 0.0146727\dots, \quad \Phi_6^* = 0.0019227\dots,$$

implying $\gamma_1(2) = 1.079707\dots$, $\gamma_2(2) = 0.766330\dots$, and $\xi(2) = \eta(2)/4 = -0.72378\dots$. Using these numerical values, we can determine the branches as in the assertion (iii).

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