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## A STUDY OF A UNILATERAL AND ADHESIVE CONTACT PROBLEM WITH NORMAL COMPLIANCE

Abstract. The aim of this paper is to study a quasistatic unilateral contact problem between an elastic body and a foundation. The constitutive law is nonlinear and the contact is modelled with a normal compliance condition associated to a unilateral constraint and Coulomb's friction law. The adhesion between contact surfaces is taken into account and is modelled with a surface variable, the bonding field, whose evolution is described by a first-order differential equation. We establish a variational formulation of the mechanical problem and prove an existence and uniqueness result in the case where the friction coefficient is small enough. The technique of proof is based on time-dependent variational inequalities, differential equations and the Banach fixed-point theorem. We also study a penalized and regularized problem which admits at least one solution and prove its convergence to the solution of the model when the penalization and regularization parameter tends to zero.

1. Introduction. Contact problems involving deformable bodies are quite frequent in industry as well as in daily life and play an important role in structural and mechanical systems. Contact processes involve complicated surface phenomena, and are modelled by highly nonlinear initial boundary value problems. Taking into account various contact conditions associated with more and more complex behavior laws leads to the introduction of new and nonstandard models, expressed with the aid of evolution variational inequalities. An early attempt to study contact problems within the framework of variational inequalities was made in 10. The mathematical, mechanical and numerical state of the art can be found in [23] where we find detailed mathematical and numerical studies of adhesive contact prob-

[^0]lems. Unilateral frictional contact problems involving Signorini's condition with or without adhesion were studied by several authors: see for instance the references in [1, 2, [5, 6, 8, 11, 17, 20, 27, 30, 32, 31].

In this paper, we study a mathematical model which describes a frictional unilateral contact problem with adhesion between a nonlinear elastic body and a deformable foundation. Following [16, 28] the contact is modelled with a normal compliance condition associated to unilateral constraint, where the penetration is limited. Recall that models for dynamic or quasistatic processes of frictionless adhesive contact between a deformable body and a foundation have been studied in [3, 4, 12, 17, 22, 23, 24, 28]. Also recently dynamic or quasistatic frictional contact problems with adhesion were studied in [7, 29]. Here as in [13, 14] we use the bonding field as an additional state variable $\beta$, defined on the contact surface of the boundary. The variable is restricted to values $0 \leq \beta \leq 1$; when $\beta=0$ all the bonds are severed and there are no active bonds, when $\beta=1$ all the bonds are active; when $0<\beta<1$ it measures the fraction of active bonds and partial adhesion takes place. We refer the reader to the extensive bibliography on


In this work we extend the result established in [27] to the unilateral contact problem with a normal compliance condition associated to Coulomb's friction law. We establish a variational formulation of the mechanical problem for which we prove the existence of a unique weak solution if the friction coefficient is small enough and obtain a partial regularity result for the solution. We also consider a penalized and regularized problem which has at least one solution and prove its convergence to the solution of the model when the parameter of penalization and regularization tends to zero.

The paper is structured as follows. In Section 2 we present some notations and give the variational formulation. In Section 3 we state and prove our main existence and uniqueness result, Theorem 3.1. In Section 4 we establish a convergence result, Theorem 4.6.
2. Problem statement and variational formulation. We consider a nonlinear elastic body which occupies a domain $\Omega \subset \mathbb{R}^{d}(d=2,3)$ and assume that its boundary $\Gamma$ is regular and partitioned into three measurable and disjoint parts $\Gamma_{1}, \Gamma_{2}, \Gamma_{3}$ such that meas $\left(\Gamma_{1}\right)>0$. The body is acted upon by a volume force of density $\varphi_{1}$ on $\Omega$ and a surface traction of density $\varphi_{2}$ on $\Gamma_{2}$. On $\Gamma_{3}$ the body is in unilateral and adhesive contact following Coulomb's friction law with a foundation.

Thus, the classical formulation of the mechanical problem is as follows.
Problem $P_{1}$. Find a displacement $u: \Omega \times[0, T] \rightarrow \mathbb{R}^{d}$ and a bonding field $\beta: \Gamma_{3} \times[0, T] \rightarrow \mathbb{R}$ such that

$$
\left.\begin{array}{ll}
\left.\begin{array}{ll}
\operatorname{div} \sigma+\varphi_{1}=0 & \text { in } \Omega \times(0, T) \\
\sigma=F \varepsilon(u) & \text { in } \Omega \times(0, T) \\
u=0 & \text { on } \Gamma_{1} \times(0, T) \\
\sigma \nu=\varphi_{2} & \text { on } \Gamma_{2} \times(0, T) \\
u_{\nu} \leq g, \sigma_{\nu}+p\left(u_{\nu}\right)-c_{\nu} \beta^{2} R_{\nu}\left(u_{\nu}\right) \leq 0 \\
\left(\sigma_{\nu}+p\left(u_{\nu}\right)-c_{\nu} \beta^{2} R_{\nu}\left(u_{\nu}\right)\right)\left(u_{\nu}-g\right)=0
\end{array}\right\} \text { on } \Gamma_{3} \times(0, T) \\
\left|\sigma_{\tau}+c_{\tau} \beta^{2} R_{\tau}\left(u_{\tau}\right)\right| \leq \mu p\left(u_{\nu}\right) \\
\left|\sigma_{\tau}+c_{\tau} \beta^{2} R_{\tau}\left(u_{\tau}\right)\right|<\mu p\left(u_{\nu}\right) \Rightarrow u_{\tau}=0 \\
\left|\sigma_{\tau}+c_{\tau} \beta^{2} R_{\tau}\left(u_{\tau}\right)\right|=\mu p\left(u_{\nu}\right) \Rightarrow \\
\exists \lambda \geq 0, \mu p\left(u_{\nu}\right)=-\left(\sigma_{\tau}+c_{\tau} \beta^{2} R_{\tau}\left(u_{\tau}\right)\right)
\end{array}\right\} \text { on } \Gamma_{3} \times(0, T),
$$

Equation (2.1) represents the equilibrium equation. Equation (2.2) represents the elastic constitutive law of the material in which $F$ is a given function and $\varepsilon(u)$ denotes the strain tensor; (2.3) and (2.4) are the displacement and traction boundary conditions, respectively, in which $\nu$ denotes the unit outward normal vector on $\Gamma$ and $\sigma \nu$ represents the Cauchy stress vector. The condition (2.5) represents the unilateral contact with adhesion in which $p$ and $-c_{\nu} \beta^{2} R_{\nu}\left(u_{\nu}\right)$ are the normal contact functions. Here $c_{\nu}$ is a given adhesion coefficient and $R_{\nu}$ is a truncation operator defined by

$$
R_{\nu}(s)= \begin{cases}L & \text { if } s<-L \\ -s & \text { if }-L \leq s \leq 0 \\ 0 & \text { if } s>0\end{cases}
$$

Here $L>0$ is the characteristic length of the bond, beyond which the bond has no additional traction (see [23]) and $p$ is a normal compliance function which satisfies assumption (2.16) below. We denote by $g \geq 0$ the maximum value of the penetration. When $u_{\nu}<0$, i.e. when there is separation between the body and the foundation, then the condition (2.5) combined with hypotheses (2.16) on the function $p$ shows that $\sigma_{\nu}=c_{\nu} \beta^{2} R_{\nu}\left(u_{\nu}\right)$ and by assumption (2.17) below, $\sigma_{\nu}$ does not exceed the value $L g$. When $g>0$, the body may interpenetrate into the foundation, but the penetration is limited, that is, $u_{\nu} \leq g$. In this case of penetration (i.e. $u_{\nu} \geq 0$ ), when $0 \leq u_{\nu}<g$ then $-\sigma_{\nu}=p\left(u_{\nu}\right)$, which means that the reaction of the foundation is uniquely determined by the normal displacement and $\sigma_{\nu} \leq 0$. Since $p$ is an increasing function, the reaction of the foundation increases with the penetration, and when $u_{\nu}=g$, then $-\sigma_{\nu} \geq p(g)$ and $\sigma_{\nu}$ is not uniquely determined. When $g>0$ and $p=0$, condition (2.5) becomes the Signorini contact condition
with adhesion with a gap function,

$$
u_{\nu} \leq g, \quad \sigma_{\nu}-c_{\nu} \beta^{2} R_{\nu}\left(u_{\nu}\right) \leq 0, \quad\left(\sigma_{\nu}-c_{\nu} \beta^{2} R_{\nu}\left(u_{\nu}\right)\right)\left(u_{\nu}-g\right)=0 .
$$

When $g=0$, (2.5) combined with assumption (2.16) becomes the Signorini contact condition with adhesion with a zero gap function, given by

$$
u_{\nu} \leq 0, \quad \sigma_{\nu}-c_{\nu} \beta^{2} R_{\nu}\left(u_{\nu}\right) \leq 0, \quad\left(\sigma_{\nu}-c_{\nu} \beta^{2} R_{\nu}\left(u_{\nu}\right)\right) u_{\nu}=0 .
$$

This contact condition was used in [6, 23, 24, 27, 28, 30, The condition (2.6) represents the frictional contact in which $c_{\tau} \beta^{2} R_{\tau}\left(u_{\tau}\right)$ is an adhesive where $c_{\tau}$ is an adhesion coefficient and $R_{\tau}$ is a truncation operator defined by

$$
R_{\tau}(v)= \begin{cases}v & \text { if }|v| \leq L, \\ L \frac{v}{|v|} & \text { if }|v|>L,\end{cases}
$$

where $L>0$ is the characteristic length of the bonds. Equation (2.7) is the ordinary differential equation which describes the evolution of the bonding field, where $\varepsilon_{a}$ is an adhesion coefficient and $\beta_{+}=\max (0, \beta)$. Since $\dot{\beta} \leq 0$ on $\Gamma_{3} \times(0, T)$, once debonding occurs bonding cannot be reestablished and, indeed, the adhesive process is irreversible. Also from [18] it must be pointed out that condition (2.7) does not allow for complete debonding in finite time. Finally, (2.8) is the initial condition, in which $\beta_{0}$ denotes the initial bonding field. In (2.7) a dot above a variable represents its derivative with respect to time. We denote by $S_{d}$ the space of second order symmetric tensors on $\mathbb{R}^{d}$ $(d=2,3)$, and $|\cdot|$ represents the Euclidean norm on $\mathbb{R}^{d}$ and $S_{d}$. Thus, for every $u, v \in \mathbb{R}^{d}$, u.v $=u_{i} v_{i},|v|=(v \cdot v)^{1 / 2}$, and for every $\sigma, \tau \in S_{d}$, $\sigma . \tau=\sigma_{i j} \tau_{i j},|\tau|=(\tau . \tau)^{1 / 2}$. Here and below, the indices $i$ and $j$ run between 1 and $d$ and the summation convention over repeated indices is adopted.

Now, to proceed with the variational formulation, we need the following function spaces:

$$
\begin{aligned}
& H=\left(L^{2}(\Omega)\right)^{d}, \quad H_{1}=\left(H^{1}(\Omega)\right)^{d}, \quad Q=\left\{\sigma=\left(\sigma_{i j}\right): \sigma_{i j}=\sigma_{j i} \in L^{2}(\Omega)\right\}, \\
& Q_{1}=\{\sigma \in Q: \operatorname{div} \sigma \in H\} .
\end{aligned}
$$

Note that $H$ and $Q$ are real Hilbert spaces endowed with the respective canonical inner products

$$
(u, v)_{H}=\int_{\Omega} u_{i} v_{i} d x, \quad(\sigma, \tau)_{Q}=\int_{\Omega} \sigma_{i j} \tau_{i j} d x
$$

The strain tensor is

$$
\varepsilon(u)=\left(\varepsilon_{i j}(u)\right)=\frac{1}{2}\left(u_{i, j}+u_{j, i}\right) ;
$$

$\operatorname{div} \sigma=\left(\sigma_{i j, j}\right)$ is the divergence of $\sigma$. For every $v \in H_{1}$ we denote by $v_{\nu}$ and $v_{\tau}$ the normal and tangential components of $v$ on the boundary $\Gamma$ given by

$$
v_{\nu}=v . \nu, \quad v_{\tau}=v-v_{\nu} \nu .
$$

We also denote by $\sigma_{\nu}$ and $\sigma_{\tau}$ the normal and the tangential traces of a function $\sigma \in Q_{1}$, and when $\sigma$ is a regular function then

$$
\sigma_{\nu}=(\sigma \nu) \cdot \nu, \quad \sigma_{\tau}=\sigma \nu-\sigma_{\nu} \nu
$$

and the following Green formula holds:

$$
(\sigma, \varepsilon(v))_{Q}+(\operatorname{div} \sigma, v)_{H}=\int_{\Gamma} \sigma \nu \cdot v d a \quad \forall v \in H_{1}
$$

where $d a$ is the surface measure element.
Now, let $V$ be the closed subspace of $H_{1}$ defined by

$$
V=\left\{v \in H_{1}: v=0 \text { on } \Gamma_{1}\right\}
$$

and let the convex subset of admissible displacements be given by

$$
K=\left\{v \in V: v_{\nu} \leq g \text { a.e. on } \Gamma_{3}\right\} .
$$

Since meas $\left(\Gamma_{1}\right)>0$, the following Korn inequality holds [10]:

$$
\begin{equation*}
\|\varepsilon(v)\|_{Q} \geq c_{\Omega}\|v\|_{H_{1}} \quad \forall v \in V \tag{2.9}
\end{equation*}
$$

where $c_{\Omega}>0$ is a constant which depends only on $\Omega$ and $\Gamma_{1}$. We equip $V$ with the inner product

$$
(u, v)_{V}=(\varepsilon(u), \varepsilon(v))_{Q}
$$

and $\|\cdot\|_{V}$ is the associated norm. It follows from (2.9) that the norms $\|\cdot\|_{H_{1}}$ and $\|\cdot\|_{V}$ are equivalent on $V$. Thus $\left(V,\|\cdot\|_{V}\right)$ is a real Hilbert space. Moreover by Sobolev's trace theorem, there exists $d_{\Omega}>0$ which only depends on the domain $\Omega, \Gamma_{1}$ and $\Gamma_{3}$ such that

$$
\begin{equation*}
\|v\|_{\left(L^{2}\left(\Gamma_{3}\right)\right)^{d}} \leq d_{\Omega}\|v\|_{V} \quad \forall v \in V \tag{2.10}
\end{equation*}
$$

For $p \in[1, \infty]$, we use the standard norm of $L^{p}(0, T ; V)$. We also use the Sobolev space $W^{1, \infty}(0, T ; V)$ equipped with the norm

$$
\|v\|_{W^{1, \infty}(0, T ; V)}=\|v\|_{L^{\infty}(0, T ; V)}+\|\dot{v}\|_{L^{\infty}(0, T ; V)}
$$

For every real Banach space $\left(X,\|\cdot\|_{X}\right)$ and $T>0$ we use the notation $C([0, T] ; X)$ for the space of continuous functions from $[0, T]$ to $X$; recall that $C([0, T] ; X)$ is a real Banach space with the norm

$$
\|x\|_{C([0, T] ; X)}=\max _{t \in[0, T]}\|x(t)\|_{X}
$$

We suppose that the body forces and surface tractions have the regularity

$$
\begin{equation*}
\varphi_{1} \in W^{1, \infty}(0, T ; H), \quad \varphi_{2} \in W^{1, \infty}\left(0, T ;\left(L^{2}\left(\Gamma_{2}\right)\right)^{d}\right) \tag{2.11}
\end{equation*}
$$

and denote by $f(t)$ the element of $V$ defined by

$$
\begin{equation*}
(f(t), v)_{V}=\int_{\Omega} \varphi_{1}(t) \cdot v d x+\int_{\Gamma_{2}} \varphi_{2}(t) \cdot v d a, \quad \forall v \in V, t \in[0, T] \tag{2.12}
\end{equation*}
$$

Using (2.11) and (2.12) yields

$$
f \in W^{1, \infty}(0, T ; V)
$$

In the study of the mechanical problem $P_{1}$ we assume that the nonlinear elasticity operator $F: \Omega \times S_{d} \rightarrow S_{d}$ satisfies:

$$
\left\{\begin{align*}
\text { (a) } & \text { there exists } M>0 \text { such that } \\
& \left|F\left(x, \varepsilon_{1}\right)-F\left(x, \varepsilon_{2}\right)\right| \leq M\left|\varepsilon_{1}-\varepsilon_{2}\right|, \\
& \forall \varepsilon_{1}, \varepsilon_{2} \in S_{d}, \text { a.e. } x \in \Omega ; \\
\text { (b) } & \text { there exists } m>0 \text { such that }  \tag{2.13}\\
& \left(F\left(x, \varepsilon_{1}\right)-F\left(x, \varepsilon_{2}\right)\right) \cdot\left(\varepsilon_{1}-\varepsilon_{2}\right) \geq m\left|\varepsilon_{1}-\varepsilon_{2}\right|^{2}, \\
& \forall \varepsilon_{1}, \varepsilon_{2} \in S_{d}, \text { a.e. } x \in \Omega ; \\
\text { (c) } & \text { the mapping } x \mapsto F(x, \varepsilon) \text { is Lebesgue measurable on } \Omega, \\
& \text { for any } \varepsilon \in S_{d} ; \\
\text { (d) } & F(x, 0)=0 \text { for a.e. } x \in \Omega .
\end{align*}\right.
$$

The adhesion coefficients are assumed to satisfy

$$
\begin{equation*}
c_{\nu}, c_{\tau} \in L^{\infty}\left(\Gamma_{3}\right), \quad \varepsilon_{a} \in L^{2}\left(\Gamma_{3}\right), \quad c_{\nu}, c_{\tau}, \varepsilon_{a} \geq 0 \quad \text { a.e. on } \Gamma_{3}, \tag{2.14}
\end{equation*}
$$

and the friction coefficient $\mu$ satisfies

$$
\begin{equation*}
\mu \in L^{\infty}\left(\Gamma_{3}\right) \quad \text { and } \quad \mu \geq 0 \text { a.e. on } \Gamma_{3} . \tag{2.15}
\end{equation*}
$$

Next we define the functional $j_{\mathrm{ad}}: L^{2}\left(\Gamma_{3}\right) \times V \times V \rightarrow \mathbb{R}$ by

$$
\begin{aligned}
j_{\mathrm{ad}}(\beta, u, v)=\int_{\Gamma_{3}}\left[\left(p\left(u_{\nu}\right)-c_{\nu} \beta^{2} R_{\nu}\left(u_{\nu}\right)\right) v_{\nu}\right. & \left.+c_{\tau} \beta^{2} R_{\tau}\left(u_{\tau}\right) \cdot v_{\tau}\right] d a \\
& \forall(\beta, u, v) \in L^{2}\left(\Gamma_{3}\right) \times V \times V,
\end{aligned}
$$

and the functional $j_{\mathrm{fr}}: V \times V \rightarrow \mathbb{R}_{+}$by

$$
j_{\mathrm{fr}}(u, v)=\int_{\Gamma_{3}} \mu p\left(u_{\nu}\right)\left|v_{\tau}\right| d a, \quad \forall(u, v) \in V \times V .
$$

We assume that the normal compliance function $p: \Gamma_{3} \times \mathbb{R} \rightarrow \mathbb{R}_{+}$satisfies:

$$
\begin{cases}\text { (a) } & \text { There exists } L_{p}>0 \text { such that } \\ & \left|p\left(x, r_{1}\right)-p\left(x, r_{2}\right)\right| \leq L_{p}\left|r_{1}-r_{2}\right|, \\ & \forall r_{1}, r_{2} \in \mathbb{R}, \text { a.e. } x \in \Gamma_{3} ; \\ \text { (b) } & \left(p\left(x, r_{1}\right)-p\left(x, r_{2}\right)\right)\left(r_{1}-r_{2}\right) \geq 0,  \tag{2.16}\\ & \forall r_{1}, r_{2} \in \mathbb{R}, \text { a.e. } x \in \Gamma_{3} ; \\ \text { (c) } & \text { the mapping } x \mapsto p_{\nu}(x, r) \text { is Lebesgue measurable on } \Gamma_{3}, \\ & \text { for any } r \in \mathbb{R} ; \\ \text { (d) } & p(x, r)=0, \forall r \leq 0, \text { a.e. } x \in \Gamma_{3} .\end{cases}
$$

We also assume that the initial bonding field satisfies

$$
\begin{equation*}
\beta_{0} \in L^{2}\left(\Gamma_{3}\right), \quad 0 \leq \beta_{0} \leq 1 \quad \text { a.e. on } \Gamma_{3}, \tag{2.17}
\end{equation*}
$$

and we introduce the set

$$
B=\left\{\theta:[0, T] \rightarrow L^{2}\left(\Gamma_{3}\right) ; 0 \leq \theta(t) \leq 1, \forall t \in[0, T], \text { a.e. on } \Gamma_{3}\right\} .
$$

Finally, assuming that the solution is sufficiently regular and applying Green's formula, we deduce the following variational formulation of the mechanical problem $P_{1}$.

Problem $P_{2}$. Find a displacement field $u:[0, T] \rightarrow V$ and a bonding field $\beta:[0, T] \rightarrow L^{2}\left(\Gamma_{3}\right)$ such that

$$
\begin{align*}
& u(t) \in K, \\
& (F \varepsilon(u(t)), \varepsilon(v)-\varepsilon(u(t)))_{Q}+j_{\mathrm{ad}}(\beta(t), u(t), v-u(t))  \tag{2.18}\\
& +j_{\mathrm{fr}}(u(t), v)-j_{\mathrm{fr}}(u(t), u(t)) \geq(f(t), v-u(t))_{V}, \forall v \in K, t \in[0, T], \\
& \dot{\beta}(t)=-\left[\beta(t)\left(c_{\nu}\left(R_{\nu}\left(u_{\nu}(t)\right)\right)^{2}+c_{\tau}\left|R_{\tau}\left(u_{\tau}(t)\right)\right|^{2}\right)-\varepsilon_{a}\right]_{+} \text {a.e. } t \in(0, T),  \tag{2.19}\\
& \beta(0)=\beta_{0} . \tag{2.20}
\end{align*}
$$

3. Existence and uniqueness result. Our main result is the following theorem.

Theorem 3.1. Let (2.11) and (2.13)-(2.17) hold. Then Problem $P_{2}$ has a unique solution, which satisfies

$$
\begin{equation*}
u \in W^{1, \infty}(0, T ; V) \cap C([0, T] ; K) \tag{3.1}
\end{equation*}
$$

and

$$
\begin{equation*}
\beta \in W^{1, \infty}\left(0, T ; L^{2}\left(\Gamma_{3}\right)\right) \cap B, \tag{3.2}
\end{equation*}
$$

if

$$
\|\mu\|_{L^{\infty}\left(\Gamma_{3}\right)}<m / L_{p} d_{\Omega}^{2} .
$$

The proof of Theorem 3.1 is carried out in several steps. In the first step, let $k>0$ and consider the closed subset $X$ of $C\left([0, T] ; L^{2}\left(\Gamma_{3}\right)\right)$ defined as

$$
X=\left\{\theta \in C\left([0, T] ; L^{2}\left(\Gamma_{3}\right)\right) \cap B: \theta(0)=\beta_{0}\right\}
$$

where the Banach space $C\left([0, T] ; L^{2}\left(\Gamma_{3}\right)\right)$ is endowed with the norm

$$
\|\theta\|_{X}=\max _{t \in[0, T]}\left[\exp (-k t)\|\theta(t)\|_{L^{2}\left(\Gamma_{3}\right)}\right] \quad \text { for } \theta \in C\left([0, T] ; L^{2}\left(\Gamma_{3}\right)\right) .
$$

Next for a given $\beta \in X$, we consider the following variational problem.
Problem $P_{1 \beta}$. Find $u_{\beta}:[0, T] \rightarrow V$ such that

$$
\begin{align*}
& u_{\beta}(t) \in K, \\
& \left(F \varepsilon\left(u_{\beta}(t)\right), \varepsilon\left(v-u_{\beta}(t)\right)\right)_{Q}+j_{\mathrm{ad}}\left(\beta(t), u_{\beta}(t), v-u_{\beta}(t)\right)  \tag{3.3}\\
& +j_{\mathrm{fr}}\left(u_{\beta}(t), v\right)-j_{\mathrm{fr}}\left(u_{\beta}(t), u_{\beta}(t)\right) \geq\left(f(t), v-u_{\beta}(t)\right)_{V} \\
& \quad \forall v \in K, t \in[0, T] .
\end{align*}
$$

We have the following result.

Proposition 3.2. Problem $P_{1 \beta}$ has a unique solution

$$
\begin{equation*}
u_{\beta} \in C([0, T] ; K) \tag{3.4}
\end{equation*}
$$

if

$$
\|\mu\|_{L^{\infty}\left(\Gamma_{3}\right)}<m / L_{p} d_{\Omega^{2}}^{2} .
$$

The proof of Proposition 3.2 will be established in several steps. In the first step for each $t \in[0, T]$ and a given $\eta \in K$, we consider the following intermediate problem.

Problem $P_{\beta \eta}$. Find $u_{\beta \eta}(t) \in K$ such that

$$
\begin{align*}
\left(F \varepsilon\left(u_{\beta \eta}(t)\right), \varepsilon(v-\right. & \left.\left.u_{\beta \eta}(t)\right)\right)_{Q}+j_{\mathrm{ad}}\left(\beta(t), u_{\beta \eta}(t), v-u_{\beta \eta}(t)\right)+j(\eta, v)  \tag{3.5}\\
& -j_{\mathrm{fr}}\left(\eta, u_{\beta \eta}(t)\right) \geq\left(f(t), v-u_{\beta \eta}(t)\right)_{V}, \quad \forall v \in K .
\end{align*}
$$

Lemma 3.3. Problem $P_{\beta \eta}$ has a unique solution.
Proof. Let the operator $A_{\beta(t)}: V \rightarrow V$ be defined by

$$
\left(A_{\beta(t)} u, v\right)_{V}=(F \varepsilon(u), \varepsilon(v))_{Q}+j_{\mathrm{ad}}(\beta(t), u, v), \quad \forall u, v \in V .
$$

We use (2.10), (2.13)(a), (2.13)(b), (2.16)(b) and (2.16)(c) to show that the operator $A_{\beta(t)}$ is strongly monotone and Lipschitz continuous.

Moreover, the functional $j(\eta, \cdot): V \rightarrow \mathbb{R}_{+}$is a continuous seminorm. Hence by a standard existence and uniqueness result for elliptic variational inequalities (see [25]), there exists a unique element $u_{\beta \eta}(t) \in K$ which satisfies (3.5) since $K$ is a non-empty, closed convex subset of $V$.

Now, in the second step, for a fixed $t \in[0, T]$ we consider the map $T_{t}: K \rightarrow K$ defined as

$$
T_{t}(\eta)=u_{\beta \eta}(t) .
$$

We have the following lemma.
Lemma 3.4. The map $T_{t}$ has a unique fixed point $\eta^{*}$ and $u_{\beta \eta^{*}}(t)$ is a unique solution of the inequality (3.3).

Proof. Let $\eta_{1}, \eta_{2} \in K$. In (3.5) satisfied by $u_{\eta_{1}}(t)$ set $v=u_{\eta_{2}}(t)$, and also in the same inequality satisfied by $u_{\eta_{2}}(t)$ take $v=u_{\eta_{1}}(t)$. Using (2.10), (2.13)(c) and (2.16), it follows after adding the resulting inequalities that

$$
\left\|T_{t}\left(\eta_{1}\right)-T_{t}\left(\eta_{2}\right)\right\|_{V} \leq \frac{\|\mu\|_{L^{\infty}\left(\Gamma_{3}\right)} L_{p} d_{\Omega}^{2}}{m}\left\|\eta_{1}-\eta_{2}\right\|_{V}, \quad \forall \eta_{1}, \eta_{2} \in K .
$$

Then for $\|\mu\|_{L^{\infty}\left(\Gamma_{3}\right)} L_{p} d_{\Omega}^{2} / m<1$, the map $T_{t}$ is a contraction; so it has a unique fixed point $\eta^{*}$ and $u_{\beta \eta^{*}}(t)$ is a unique solution of inequality (3.3). Next, denote $u_{\beta \eta^{*}}(t)=u_{\beta}(t)$ for each $t \in[0, T]$. As in [31], to show that $u_{\beta} \in C([0, T] ; K)$, it suffices to see from (3.3) that there exists a constant
$c>0$ such that

$$
\begin{align*}
& \left\|u_{\beta}\left(t_{1}\right)-u_{\beta}\left(t_{2}\right)\right\|_{V}  \tag{3.6}\\
& \quad \leq \frac{c}{m-\|\mu\|_{L^{\infty}\left(\Gamma_{3}\right)} L_{p} d_{\Omega}^{2}}\left(\left\|f\left(t_{1}\right)-f\left(t_{2}\right)\right\|_{V}+\left\|\beta\left(t_{1}\right)-\beta\left(t_{2}\right)\right\|_{L^{2}\left(\Gamma_{3}\right)}\right)
\end{align*}
$$

for all $t_{1}, t_{2} \in[0, T]$. Therefore, as $f \in C([0, T] ; V)$ and $\beta \in C\left([0, T] ; L^{2}\left(\Gamma_{3}\right)\right)$, we immediately deduce (3.4).

In the second step, we consider the following problem.
Problem $P_{2 \beta}$. Find $\chi_{\beta}:[0, T] \rightarrow L^{2}\left(\Gamma_{3}\right)$ such that

$$
\begin{align*}
& \dot{\chi}_{\beta}(t)=  \tag{3.7}\\
& -\left[\chi_{\beta}(t)\left(c_{\nu}\left(R_{\nu}\left(u_{\beta \nu}(t)\right)\right)^{2}+c_{\tau}\left|R_{\tau}\left(u_{\beta \tau}(t)\right)\right|^{2}\right)-\varepsilon_{a}\right]_{+} \quad \text { a.e. } t \in(0, T) \\
& \chi_{\beta}(0)=\beta_{0} \tag{3.8}
\end{align*}
$$

We obtain the following result.
Lemma 3.5. Problem $P_{2 \beta}$ has a unique solution $\chi_{\beta}$ which satisfies

$$
\chi_{\beta} \in W^{1, \infty}\left(0, T ; L^{2}\left(\Gamma_{3}\right)\right) \cap B
$$

Proof. Consider the mapping $F_{\beta}(t, \theta):[0, T] \times L^{2}\left(\Gamma_{3}\right) \rightarrow L^{2}\left(\Gamma_{3}\right)$ defined by

$$
F_{\beta}(t, \theta)=-\left[\theta\left(c_{\nu}\left(R_{\nu}\left(u_{\beta \nu}(t)\right)\right)^{2}+c_{\tau}\left|R_{\tau}\left(u_{\beta \tau}(t)\right)\right|^{2}\right)-\varepsilon_{a}\right]_{+}
$$

It follows from the properties of the truncation operators $R_{\nu}$ and $R_{\tau}$, that $F_{\beta}$ is Lipschitz continuous with respect to the second argument, uniformly in time. Moreover, for any $\theta \in L^{2}\left(\Gamma_{3}\right)$, the mapping $t \mapsto F_{\beta}(t, \theta)$ belongs to $L^{\infty}\left(0, T ; L^{2}\left(\Gamma_{3}\right)\right)$. Then, from a version of the Cauchy-Lipschitz theorem, we deduce the existence of a unique fonction $\chi_{\beta} \in W^{1, \infty}\left(0, T ; L^{2}\left(\Gamma_{3}\right)\right)$, which satisfies (3.7), (3.8). The regularity $\chi_{\beta} \in B$ follows from (3.7), (3.8) and (2.17) (see [22, 24, 26]). Therefore, from Lemma 3.5, we deduce that for all $\beta \in X$, the solution $\chi_{\beta}$ of Problem $P_{2 \beta}$ belongs to $X$.

Next, we define the mapping $\Lambda: X \rightarrow X$ by

$$
\Lambda \beta=\chi_{\beta}
$$

The third step consists of the following lemma.
Lemma 3.6. The mapping $\Lambda$ has a unique fixed point $\beta^{*}$.
Proof. We have

$$
\Lambda \beta(t)=\beta_{0}-\int_{0}^{t}\left[\chi_{\beta}(s)\left(c_{\nu}\left(R_{\nu}\left(u_{\beta \nu}(s)\right)\right)^{2}+c_{\tau}\left|R_{\tau}\left(u_{\beta \tau}(s)\right)\right|^{2}\right)-\varepsilon_{a}\right]_{+} d s
$$

where $u_{\beta}$ is the solution of Problem $P_{1 \beta}$. Then for $\beta_{1}, \beta_{2} \in X$, using (2.19)(a) and the properties of $R_{\nu}$ and $R_{\tau}$ (see [23]), there exists a constant $c_{1}>0$ such that

$$
\left|\chi_{\beta_{1}}(t)-\chi_{\beta_{2}}(t)\right| \leq c_{1} \int_{0}^{t}\left(\left|\chi_{\beta_{1}}(s)-\chi_{\beta_{2}}(s)\right|+\left|u_{\beta_{1} \tau}(s)-u_{\beta_{2} \tau}(s)\right|\right) d s .
$$

Applying Gronwall's inequality and using (2.10) yields

$$
\left\|\chi_{\beta_{1}}(t)-\chi_{\beta_{2}}(t)\right\|_{L^{2}\left(\Gamma_{3}\right)} \leq c_{2} \int_{0}^{t}\left\|u_{\beta_{1}}(s)-u_{\beta_{2}}(s)\right\|_{V} d s
$$

for some constant $c_{2}>0$. Now let $t \in[0, T]$. Then, using the inequalities (3.3), (2.13), (2.16) and $\|\mu\|_{L^{\infty}\left(\Gamma_{3}\right)} L_{p} d_{\Omega}^{2}<m$, we deduce that there exists a constant $c_{3}>0$ (see 31) such that

$$
\left\|u_{\beta_{1}}(t)-u_{\beta_{2}}(t)\right\|_{V} \leq c_{3}\left\|\beta_{1}(t)-\beta_{2}(t)\right\|_{L^{2}\left(\Gamma_{3}\right)} .
$$

Hence,

$$
\left\|\Lambda \beta_{1}(t)-\Lambda \beta_{2}(t)\right\|_{L^{2}\left(\Gamma_{3}\right)} \leq c_{4} \int_{0}^{t}\left\|\beta_{1}(s)-\beta_{2}(s)\right\|_{L^{2}\left(\Gamma_{3}\right)} d s, \quad \forall t \in[0, T],
$$

for some constant $c_{4}>0$. Therefore,

$$
\left\|\Lambda \beta_{1}-\Lambda \beta_{2}\right\|_{X} \leq \frac{c_{4}}{k}\left\|\beta_{1}-\beta_{2}\right\|_{X}, \quad \forall \beta_{1}, \beta_{2} \in X
$$

Thus, this previous inequality shows that for $k$ sufficiently large, $\Lambda$ is a contraction. Hence it has a unique fixed point $\beta^{*}$ which satisfies (3.7) and (3.8). On the other hand, from (3.6) we deduce that $u_{\beta^{*}} \in W^{1, \infty}(0, T ; V)$.

Proof of Theorem 3.1. Let $\beta=\beta^{*}$ and let $u_{\beta^{*}}$ be the solution to Problem $P_{1 \beta}$. We conclude by (3.3), (3.7) and (3.8) that ( $u_{\beta^{*}}, \beta^{*}$ ) is a solution of Problem $P_{2}$. Now to prove the uniqueness of the solution, suppose that $(u, \beta)$ is a solution of Problem $P_{2}$ which satisfies (2.18)-(2.20). It follows from (2.18) that $u$ is a solution of Problem $P_{1 \beta}$ and by Proposition 3.2 we get $u=u_{\beta}$. Taking $u=u_{\beta}$ in (2.18) and using the initial condition (2.20), we deduce that $\beta$ is a solution of Problem $P_{2 \beta}$. Finally, using Lemma 3.5, we obtain $\beta=\beta^{*}$ and so $\left(u_{\beta^{*}}, \beta^{*}\right)$ is a unique solution to Problem $P_{2}$ which satisfies (3.1), (3.2).
4. A convergence result. In this section we consider a frictional contact problem with normal compliance and adhesion with unlimited penetration. The contact condition (2.5) is replaced by the contact condition

$$
-\sigma_{\nu}=\frac{1}{\delta}\left(u_{\nu}-g\right)_{+}+p\left(u_{\nu}\right)-c_{\nu} \beta^{2} R_{\nu}\left(u_{\nu}\right) \quad \text { on } \Gamma_{3} \times(0, T)
$$

We recall that $\delta>0$ is the penalization and regularization parameter and $1 / \delta$ is interpreted as the stiffness coefficient of the foundation. We understand that when $\delta$ is small, the reaction of the foundation to the penetration is important; also when $\delta$ is large then the reaction is weaker. We study the behavior of the solution as $\delta \rightarrow 0$ and prove that in the limit we obtain the solution of the adhesive frictionless contact problem with normal compliance and finite penetration. Next we define the functionals $j_{\mathrm{ad} \delta}: L^{2}\left(\Gamma_{3}\right) \times V \times V \rightarrow$ $\mathbb{R}$ and $j_{\text {fr } \delta}: V \times V \rightarrow \mathbb{R}$ by

$$
\begin{aligned}
j_{\mathrm{ad} \delta}(\beta, u, v)= & \int_{\Gamma_{3}}\left(\frac{1}{\delta}\left(u_{\nu}-g\right)_{+}+p\left(u_{\nu}\right)-c_{\nu} \beta^{2} R_{\nu}\left(u_{\nu}\right)\right) v_{\nu} d a \\
& \forall(\beta, u, v) \in L^{2}\left(\Gamma_{3}\right) \times V \times V \\
j_{\mathrm{fr} \delta}(u, v) & =\int_{\Gamma_{3}} \mu p\left(u_{\nu}\right) \frac{u_{\tau}}{\sqrt{u_{\tau}^{2}+\delta^{2}}} v_{\tau} d a, \quad \forall u, v \in V
\end{aligned}
$$

With this notation, the variational formulation of the penalized and regularized problem with frictional contact and adhesion is the following.

Problem $P_{\delta}$. Find $u_{\delta}:[0, T] \rightarrow V$ and $\beta_{\delta}:[0, T] \rightarrow[0,1]$ such that

$$
\begin{align*}
\left(F \varepsilon\left(u_{\delta}(t)\right), \varepsilon(v)\right)_{Q}+j_{\mathrm{ad} \delta}\left(\beta_{\delta}(t), u_{\delta}(t)\right. & , v)+j_{\mathrm{fr} \delta}\left(u_{\delta}(t), v\right)  \tag{4.1}\\
& =(f(t), v)_{V}, \quad \forall v \in V, t \in[0, T] \tag{4.2}
\end{align*}
$$

We have the following result.
Theorem 4.1. Problem $P_{\delta}$ has a solution which satisfies

$$
u_{\delta} \in L^{\infty}(0, T ; V), \quad \beta_{\delta} \in W^{1, \infty}\left(0, T ; L^{2}\left(\Gamma_{3}\right)\right) \cap B
$$

Proof. As in [24], the proof is similar to the proof of Theorem 3.1 and it is carried out in several steps.
(i) For any $\beta \in X$, we consider the problem below.

Problem $P_{1 \delta}$. Find $u_{\delta}:[0, T] \rightarrow V$ such that

$$
\begin{align*}
\left(F \varepsilon\left(u_{\delta}(t)\right), \varepsilon(v)\right)_{Q}+j_{\mathrm{ad} \delta}(\beta(t) & \left., u_{\delta}(t), v\right)+j_{\mathrm{fr} \delta}\left(u_{\delta}(t), v\right)  \tag{4.4}\\
& =(f(t), v)_{V}, \quad \forall v \in V, t \in[0, T]
\end{align*}
$$

We have the following lemma.
Lemma 4.2. Problem $P_{1 \delta}$ has a solution $u_{\delta} \in L^{\infty}(0, T ; V)$.

Proof. As in [11], to prove Lemma 4.2 we use the theorem on pseudomonotone operators. For this, let us state some properties and abstract results for these operators, obtained in [33]. We start with the following definition.

Definition 4.3. Let $X$ be a reflexive Banach space and $X^{\prime}$ be its dual space. An operator $A: X \rightarrow X^{\prime}$ is called pseudomonotone if

$$
\begin{aligned}
u_{k} \rightarrow u \text { weakly in } & X, \limsup _{k \rightarrow \infty}\left\langle A u_{k}, u_{k}-u\right\rangle_{X^{\prime} \times X} \leq 0 \\
& \Rightarrow\langle A u, u-w\rangle_{X^{\prime} \times X} \leq \liminf _{k \rightarrow \infty}\left\langle A u_{k}, u_{k}-w\right\rangle_{X^{\prime} \times X}, \forall w \in X .
\end{aligned}
$$

In order to decide whether a given operator is pseudomonotone or not, we need Proposition 4.4 and Theorem 4.5 below.

Proposition 4.4. Let $X$ be a reflexive Banach space, $A: X \rightarrow X^{\prime}$ be strongly monotone and $B: X \rightarrow X^{\prime}$ be completely continuous. Then the sum $C=A+B$ is pseudomonotone.

A nonlinear operator $C: X \rightarrow X^{\prime}$ is called coercive if

$$
\lim _{\|u\|_{X} \rightarrow \infty} \frac{\langle A u, u\rangle_{X^{\prime} \times X}}{\|u\|_{X}}=+\infty
$$

The theorem on pseudomonotone operators that we use is the following.
TheOrem 4.5. Let $X$ be a real, reflexive and separable Banach space and $C: X \rightarrow X^{\prime}$ be pseudomonotone, bounded and coercive. Then for each $L \in X^{\prime}$ the equation

$$
C u=L
$$

has a solution.
In order to apply this theorem to our situation, we define the operators $A_{t}, P: V \rightarrow V^{\prime}$ by

$$
\begin{aligned}
\left\langle A_{t} u, v\right\rangle_{V^{\prime} \times V} & =\langle F \varepsilon(u), \varepsilon(v)\rangle_{Q}+j_{\mathrm{ad} \delta}(\beta(t), u, v) \\
\langle P u, v\rangle_{V^{\prime} \times V} & =\int_{\Gamma_{3}} \mu p\left(u_{\nu}\right) \frac{u_{\tau}}{\sqrt{u_{\tau}^{2}+\delta^{2}}} v_{\tau} d a, \quad \forall u, v \in V
\end{aligned}
$$

We use $(2.13),(2.14)$ and (2.16) and the definition of the operator $R_{\nu}$ to see that $A_{t}$ is strongly monotone and Lipschitz continuous. We also use (2.15) and (2.16), the compact embedding $H^{1 / 2}\left(\Gamma_{3}\right) \hookrightarrow L^{2}\left(\Gamma_{3}\right)$ and the Lebesgue dominated convergence theorem to show that the operator $P$ is completely continuous. We have

$$
\left|\langle P u, v\rangle_{V^{\prime} \times V}\right| \leq\|\mu\|_{L^{\infty}\left(\Gamma_{3}\right)} L_{p} d_{\Omega}^{2}\|u\|_{V}\|v\|_{V}, \quad \forall u, v \in V
$$

Moreover,

$$
\langle P v, v\rangle_{V^{\prime} \times V} \geq 0, \quad \forall v \in V
$$

Thus the operator $C_{t}=A_{t}+P$ is pseudomonotone, bounded and coercive. Consequently, by applying Theorem 4.5 we deduce that equation (4.4) has solution $u_{\delta}(t) \in V$. To prove that $u_{\delta} \in L^{\infty}(0, T ; V)$, it suffices to take $v=u_{\delta}(t)$ in (4.4); since

$$
j_{\mathrm{ad} \delta}\left(\beta_{\delta}(t), u_{\delta}(t), u_{\delta}(t)\right) \geq 0, \quad j_{\mathrm{fr} \delta}\left(u_{\delta}(t), u_{\delta}(t)\right) \geq 0
$$

it follows from (4.4) that

$$
\left(F \varepsilon\left(u_{\delta}(t)\right), \varepsilon\left(u_{\delta}(t)\right)\right)_{Q} \leq\left(f(t), u_{\delta}(t)\right)_{V}
$$

This inequality together with the assumption (2.13)(b) implies that

$$
\begin{equation*}
\left\|u_{\delta}(t)\right\|_{V} \leq\|f(t)\|_{V} / m \tag{4.5}
\end{equation*}
$$

Hence using the regularity $f \in C([0, T] ; V)$, we immediately conclude from (4.5) that $u_{\delta} \in L^{\infty}(0, T ; V)$.
(ii) There exists a unique $\beta_{\delta}$ such that

$$
\begin{align*}
& \beta_{\delta} \in W^{1, \infty}\left(0, T ; L^{2}\left(\Gamma_{3}\right)\right) \cap B  \tag{4.6}\\
& \dot{\beta}_{\delta}(t)= \\
& \quad-\left[\beta_{\delta}(t)\left(c_{\nu}\left(R_{\nu}\left(u_{\delta \nu}(t)\right)\right)^{2}+c_{\tau}\left|R_{\tau}\left(u_{\delta \tau}(t)\right)\right|^{2}\right)-\varepsilon_{a}\right]_{+} \text {a.e. } t \in(0, T) \tag{4.8}
\end{align*}
$$

$\beta_{\delta}(0)=\beta_{0}$
(iii) Let $\beta_{\delta}$ be as defined in (ii) and denote again by $u_{\delta}$ the function obtained in step (i) for $\beta=\beta_{\delta}$. Then, by using (4.4), (4.7) and (4.8) it is easy to see that $\left(u_{\delta}, \beta_{\delta}\right)$ is a solution to Problem $P_{\delta}$ and it satisfies

$$
\left(u_{\delta}, \beta_{\delta}\right) \in L^{\infty}(0, T ; V) \times W^{1, \infty}\left(0, T ; L^{2}\left(\Gamma_{3}\right)\right) \cap B
$$

Now, as in [23, 24], we study the convergence of the solution $\left(u_{\delta}, \beta_{\delta}\right)$ as $\delta \rightarrow 0$ in the following theorem.

Theorem 4.6. Assume that (2.13)-(2.16) hold. Then

$$
\begin{align*}
\lim _{\delta \rightarrow 0}\left\|u_{\delta}(t)-u(t)\right\|_{V}=0 & \text { for all } t \in[0, T]  \tag{4.9}\\
\lim _{\delta \rightarrow 0}\left\|\beta_{\delta}(t)-\beta(t)\right\|_{L^{2}\left(\Gamma_{3}\right)}=0 & \text { for all } t \in[0, T] \tag{4.10}
\end{align*}
$$

The proof is carried out in several steps. In the first step, we show the following lemma.

Lemma 4.7. For each $t \in[0, T]$, there exists $\bar{u}(t) \in K$ such that after passing to a subsequence still denoted $\left(u_{\delta}(t)\right)$ we have

$$
\begin{equation*}
u_{\delta}(t) \rightarrow \bar{u}(t) \quad \text { weakly in } V \tag{4.11}
\end{equation*}
$$

Proof. Let $t \in[0, T]$. Setting $v=u_{\delta}(t)$ in (4.1), we have

$$
\begin{align*}
\left(F \varepsilon\left(u_{\delta}(t)\right), \varepsilon\left(u_{\delta}(t)\right)\right)_{Q}+j_{\mathrm{ad} \delta} & \left(\beta_{\delta}(t), u_{\delta}(t), u_{\delta}(t)\right)  \tag{4.12}\\
& +j_{\mathrm{fr} \delta}\left(u_{\delta}(t), u_{\delta}(t)\right)=\left(f(t), u_{\delta}(t)\right)_{V}
\end{align*}
$$

Then it is easy to see from (4.12) that estimate (4.5) holds and so there exists an element $\bar{u}(t) \in V$ and a subsequence still denoted $u_{\delta}(t)$ such that

$$
u_{\delta}(t) \rightarrow \bar{u}(t) \quad \text { weakly on } V
$$

On the other hand, from (4.12) we also have

$$
\int_{\Gamma_{3}}\left(\frac{u_{\delta \nu}(t)-g}{\delta}\right)_{+}\left(u_{\delta \nu}(t)-g\right) d a \leq\left(f(t), u_{\delta}(t)\right)_{V}
$$

which implies that

$$
\left\|\left(u_{\delta \nu}(t)-g\right)_{+}\right\|_{L^{2}\left(\Gamma_{3}\right)}^{2} \leq \delta\|f(t)\|_{V}^{2} / m
$$

Then by (4.11), it follows that

$$
\begin{equation*}
\left\|\left(\bar{u}_{\nu}(t)-g\right)_{+}\right\|_{L^{2}\left(\Gamma_{3}\right)} \leq \liminf _{\delta \rightarrow 0}\left\|\left(u_{\delta \nu}(t)-g\right)_{+}\right\|_{L^{2}\left(\Gamma_{3}\right)}=0 \tag{4.13}
\end{equation*}
$$

Therefore we conclude by $(4.13)$ that $\left(\bar{u}_{\nu}(t)-g\right)_{+}=0$, i.e. $\bar{u}_{\nu}(t) \leq g$ a.e. on $\Gamma_{3}$, which shows that $\bar{u}(t) \in K$.

Now, we state the following problem.
Problem $P_{3}$. Find $\beta:[0, T] \rightarrow L^{2}\left(\Gamma_{3}\right)$ such that

$$
\begin{aligned}
& \dot{\beta}(t)=-\left[\beta(t)\left(c_{\nu}\left(R_{\nu}\left(\bar{u}_{\nu}(t)\right)\right)^{2}+c_{\tau}\left|R_{\tau}\left(\bar{u}_{\tau}(t)\right)\right|^{2}\right)-\varepsilon_{a}\right]_{+} \text {a.e. } t \in(0, T) \\
& \beta(0)=\beta_{0}
\end{aligned}
$$

As in [24, Lemma 3.2] we have the following result.
LEMMA 4.8. Problem $P_{3}$ has a unique solution $\beta \in W^{1, \infty}\left(0, T ; L^{2}\left(\Gamma_{3}\right)\right)$ $\cap B$.

We also show the following convergence result.
Lemma 4.9. Let $\beta$ be the solution to Problem $P_{3}$. Then

$$
\begin{equation*}
\lim _{\delta \rightarrow 0}\left\|\beta_{\delta}(t)-\beta(t)\right\|_{L^{2}\left(\Gamma_{3}\right)}=0 \quad \text { for all } t \in[0, T] \tag{4.14}
\end{equation*}
$$

Proof. As in [24, Lemma 3.2], we have

$$
\begin{align*}
\| \beta_{\delta}(t) & -\beta(t) \|_{L^{2}\left(\Gamma_{3}\right)}  \tag{4.15}\\
& \leq c \int_{0}^{t}\left(\left\|u_{\delta \nu}(s)-\bar{u}_{\nu}(s)\right\|_{L^{2}\left(\Gamma_{3}\right)}+\left\|u_{\delta \tau}(s)-\bar{u}_{\tau}(s)\right\|_{\left.\left(L^{2}\left(\Gamma_{3}\right)\right)^{d}\right)}\right) d s
\end{align*}
$$

for some $c>0$. Using (4.11) we deduce that $u_{\delta \nu}(t) \rightarrow \bar{u}_{\nu}(t)$ strongly in $L^{2}\left(\Gamma_{3}\right)$ and $u_{\delta \tau}(t) \rightarrow \bar{u}_{\tau}(t)$ strongly in $\left(L^{2}\left(\Gamma_{3}\right)\right)^{d}$ as $\delta \rightarrow 0$. On the other hand, we have

$$
\begin{aligned}
\left\|u_{\delta \nu}(t)-\bar{u}_{\nu}(t)\right\|_{L^{2}\left(\Gamma_{3}\right)}+ & \left\|u_{\delta \tau}(s)-\bar{u}_{\tau}(s)\right\|_{\left(L^{2}\left(\Gamma_{3}\right)\right)^{d}} \\
& \leq 2 d_{\Omega}\left\|u_{\delta}(t)-\bar{u}(t)\right\|_{V} \leq d_{\Omega}\left(\frac{\|f(t)\|_{V}}{m}+\|\bar{u}(t)\|_{V}\right)
\end{aligned}
$$

which implies that there exists a constant $c_{1}>0$ such that

$$
\left\|u_{\delta \nu}(t)-\bar{u}_{\nu}(t)\right\|_{L^{2}\left(\Gamma_{3}\right)}+\left\|u_{\delta \tau}(s)-\bar{u}_{\tau}(s)\right\|_{\left(L^{2}\left(\Gamma_{3}\right)\right)^{d}} \leq c_{1}
$$

Then it follows from Lebesgue's convergence theorem that

$$
\begin{equation*}
\lim _{\delta \rightarrow 0} \int_{0}^{t}\left(\left\|u_{\delta \nu}(s)-\bar{u}_{\nu}(s)\right\|_{L^{2}\left(\Gamma_{3}\right)}+\left\|u_{\delta \tau}(s)-\bar{u}_{\tau}(s)\right\|_{\left(L^{2}\left(\Gamma_{3}\right)\right)^{d}}\right) d s=0 \tag{4.16}
\end{equation*}
$$

The convergence result is now a consequence of (4.15) and (4.16).
Next, we prove the following lemma.
Lemma 4.10. We have $\bar{u}(t)=u(t)$ for all $t \in[0, T]$.
Proof. Let $v \in K$ and take $v-u_{\delta}(t)$ in (4.1) to obtain

$$
\begin{align*}
\left(F \varepsilon\left(u_{\delta}(t)\right),\right. & \left.\varepsilon\left(v-u_{\delta}(t)\right)\right)_{Q}+j_{\mathrm{ad} \delta}\left(\beta_{\delta}(t), u_{\delta}(t), v-u_{\delta}(t)\right)  \tag{4.17}\\
& +j_{\mathrm{fr} \delta}\left(u_{\delta}(t), v-u_{\delta}(t)\right) \geq\left(f(t), v-u_{\delta}(t)\right)_{V}, \quad \forall v \in K
\end{align*}
$$

Since

$$
\begin{aligned}
& j_{\mathrm{ad} \delta}\left(\beta_{\delta}(t), u_{\delta}(t), v-u_{\delta}(t)\right) \\
& \quad=\int_{\Gamma_{3}}\left(\left(\frac{u_{\delta \nu}(t)-g}{\delta}\right)_{+}+p\left(u_{\delta \nu}(t)\right)-c_{\nu} \beta_{\delta}^{2} R_{\nu}\left(u_{\delta \nu}(t)\right)\right)\left(v_{\nu}-u_{\delta \nu}(t)\right) d a \\
& \quad \leq \int_{\Gamma_{3}}\left(p\left(u_{\delta \nu}(t)\right)-c_{\nu} \beta_{\delta}^{2} R_{\nu}\left(u_{\delta \nu}(t)\right)\left(v_{\nu}-u_{\delta \nu}(t)\right)\right) d a
\end{aligned}
$$

we use (2.16), (4.10), (4.12), the properties of $R_{\nu}$ and the compact imbedding $H^{1 / 2}\left(\Gamma_{3}\right) \hookrightarrow L^{2}\left(\Gamma_{3}\right)$ to see that

$$
\int_{\Gamma_{3}}\left(p\left(u_{\delta \nu}(t)\right)-c_{\nu} \beta_{\delta}^{2} R_{\nu}\left(u_{\delta \nu}(t)\right)\right)\left(v_{\nu}-u_{\delta \nu}(t)\right) d a \rightarrow j_{\mathrm{ad}}(\beta(t), \bar{u}(t), v-\bar{u}(t))
$$

as $\delta \rightarrow 0$ and

$$
\lim _{\delta \rightarrow 0} j_{\mathrm{fr} \delta}\left(u_{\delta}(t), v-u_{\delta}(t)\right) \leq j_{\mathrm{fr}}(\bar{u}(t), v)-j_{\mathrm{fr}}(\bar{u}(t), \bar{u}(t))
$$

Therefore, passing to the limit in (4.17) as $\delta \rightarrow 0$, we obtain

$$
\begin{align*}
& (F \varepsilon(\bar{u}(t)), \varepsilon(v-\bar{u}(t)))_{Q}+j_{\mathrm{ad}}(\beta(t), \bar{u}(t), v-\bar{u}(t))+j_{\mathrm{fr}}(\bar{u}(t), v)  \tag{4.18}\\
& \quad-j_{\mathrm{fr}}(\bar{u}(t), \bar{u}(t)) \geq(f(t), v-\bar{u}(t))_{V}, \quad \forall v \in K .
\end{align*}
$$

Now, setting $v=u(t)$ in (4.18) and $v=\bar{u}(t)$ in (2.18) and adding the resulting inequalities, we obtain, by using the assumption (2.13)(b) on $F$,

$$
\begin{aligned}
& m\|\bar{u}(t)-u(t)\|_{V}^{2} \leq j_{\mathrm{ad}}(\beta(t), \bar{u}(t), u(t)-\bar{u}(t))+j_{\mathrm{ad}}(\beta(t), u(t), \bar{u}(t)-u(t)) \\
&+j_{\mathrm{fr}}(\bar{u}(t), u(t))-j_{\mathrm{fr}}(\bar{u}(t), \bar{u}(t))+j_{\mathrm{fr}}(u(t), \bar{u}(t))-j_{\mathrm{fr}}(u(t), u(t)) .
\end{aligned}
$$

Moreover using (2.16) and the properties of $R_{\nu}$ and $R_{\tau}$ we see that

$$
j_{\mathrm{ad}}(\beta(t), \bar{u}(t), u(t)-\bar{u}(t))+j_{\mathrm{ad}}(\beta(t), u(t), \bar{u}(t)-u(t)) \leq 0
$$

which implies that

$$
\begin{aligned}
& m\|\bar{u}(t)-u(t)\|_{V}^{2} \\
& \quad \leq j_{\mathrm{fr}}(\bar{u}(t), u(t))-j_{\mathrm{fr}}(\bar{u}(t), \bar{u}(t))+j_{\mathrm{fr}}(u(t), \bar{u}(t))-j_{\mathrm{fr}}(u(t), u(t))
\end{aligned}
$$

On the other hand using (2.10) and (2.16) we have

$$
\begin{aligned}
j_{\mathrm{fr}}(\bar{u}(t), u(t))-j_{\mathrm{fr}}(\bar{u}(t), \bar{u}(t))+j_{\mathrm{fr}}(u(t), \bar{u}(t))-j_{\mathrm{fr}}(u(t) & , u(t)) \\
& \leq L_{p} d_{\Omega}^{2}\|\bar{u}(t)-u(t)\|_{V}^{2}
\end{aligned}
$$

Hence we get

$$
\left(m-L_{p} d_{\Omega}^{2}\|\mu\|_{L^{\infty}\left(\Gamma_{3}\right)}\right)\|\bar{u}(t)-u(t)\|_{V}^{2} \leq 0
$$

and so as $m-L_{p} d_{\Omega}^{2}\|\mu\|_{L^{\infty}\left(\Gamma_{3}\right)}>0$, we obtain

$$
\begin{equation*}
\bar{u}(t)=u(t) \tag{4.19}
\end{equation*}
$$

Now, we have all the ingredients to prove Theorem 4.6. Indeed, from (4.19), we deduce immediately (4.10). To prove (4.9), we set $v=u(t)$ in (4.17) to obtain

$$
\begin{aligned}
m \| u_{\delta}(t)- & u(t) \|_{V}^{2} \\
\leq & j_{\mathrm{ad}}\left(\beta_{\delta}(t), u_{\delta}(t), u(t)-u_{\delta}(t)\right)-j_{\mathrm{ad}}\left(\beta(t), u(t), u(t)-u_{\delta}(t)\right) \\
& +j_{\mathrm{fr} \delta}\left(u_{\delta}(t), u(t)-u_{\delta}(t)\right)-j_{\mathrm{fr}}\left(u(t), u(t)-u_{\delta}(t)\right) \\
& +\left(F \varepsilon(u(t)), \varepsilon\left(u(t)-u_{\delta}(t)\right)\right)_{Q}+\left(f(t), u_{\delta}(t)-u(t)\right)_{V}
\end{aligned}
$$

Passing to the limit as $\delta \rightarrow 0$ and using the convergences

$$
\begin{aligned}
& j_{\mathrm{ad}}\left(\beta_{\delta}(t), u_{\delta}(t), u(t)-u_{\delta}(t)\right)-j_{\mathrm{ad}}\left(\beta(t), u(t), u(t)-u_{\delta}(t)\right) \rightarrow 0 \\
& j_{\mathrm{fr} \delta}\left(u_{\delta}(t), u(t)-u_{\delta}(t)\right)-j_{\mathrm{fr}}\left(u(t), u(t)-u_{\delta}(t)\right) \rightarrow 0 \\
& \left(F \varepsilon(u(t)), \varepsilon\left(u(t)-u_{\delta}(t)\right)\right)_{Q}+\left(f(t), u_{\delta}(t)-u(t)\right)_{V} \rightarrow 0
\end{aligned}
$$

we obtain the strong convergence

$$
\left\|u_{\delta}(t)-u(t)\right\|_{V} \rightarrow 0 \quad \text { for all } t \in[0, T]
$$

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