

ON A CUBIC HECKE ALGEBRA ASSOCIATED WITH THE QUANTUM GROUP $U_q(2)$

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Abstract. We define an operator α on $\mathbb{C}^3 \otimes \mathbb{C}^3$ associated with the quantum group $U_q(2)$, which satisfies the Yang-Baxter equation and a cubic equation $(\alpha^2 - 1)(\alpha + q^2) = 0$. This operator can be extended to a family of operators $h_j := I_j \otimes \alpha \otimes I_{n-2-j}$ on $(\mathbb{C}^3)^{\otimes n}$ with $0 \leq j \leq n-2$. These operators generate the cubic Hecke algebra $\mathcal{H}_{q,n}(2)$ associated with the quantum group $U_q(2)$. The purpose of this note is to present the construction.

1. Introduction. In [SLW3] Woronowicz considered the algebra of operators which intertwine the n -th tensor power of the fundamental representation f_N of the quantum group $SU_q(N)$ with itself. For $n = 2$ it is generated by

$$\sigma = I - \frac{\mu^{-2N+4}}{[N-2]_{q^2}!} (I \otimes E^*)(E \otimes I).$$

Here

$$[N]_{q^2}! = \prod_{k=1}^N \frac{1 - q^{2k}}{1 - q^2},$$

and $E : \mathbb{C} \rightarrow (\mathbb{C}^N)^{\otimes N}$ and $E^* : (\mathbb{C}^N)^{\otimes N} \rightarrow \mathbb{C}$ are defined as

$$E(1) = \sum_{k_1, \dots, k_N=1}^N E_{k_1, \dots, k_N} \cdot \varepsilon_{k_1} \otimes \dots \otimes \varepsilon_{k_N}, \quad E^*(\varepsilon_{k_1} \otimes \dots \otimes \varepsilon_{k_N}) = E_{k_1, \dots, k_N} \quad (1.1)$$

for the standard basis $\{\varepsilon_1, \dots, \varepsilon_N\}$ of \mathbb{C}^N . This operator σ can be written explicitly (see [SLW3], formula (4.13)) as

$$\sigma(\varepsilon_a \otimes \varepsilon_b) = \begin{cases} q \cdot \varepsilon_b \otimes \varepsilon_a & \text{for } a < b, \\ \varepsilon_b \otimes \varepsilon_a & \text{for } a = b. \\ q\varepsilon_b \otimes \varepsilon_a + (1 - q^2)\varepsilon_a \otimes \varepsilon_b & \text{for } a > b. \end{cases} \quad (1.2)$$

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It has interesting properties, in particular it is a self-adjoint Hecke operator and satisfies the Yang-Baxter equation ([SLW3], (4.14)–(4.16)):

$$\sigma^* = \sigma, \quad \sigma^2 = (1 - q^2)\sigma + q^2I, \quad (\sigma \otimes I)(I \otimes \sigma)(\sigma \otimes I) = (I \otimes \sigma)(\sigma \otimes I)(I \otimes \sigma) \quad (1.3)$$

Therefore it defines a Hecke algebra $H_{q,n}$, generated by elements $\{g_j : j = 1, \dots, n\}$, if one puts

$$g_j := I_j \otimes \sigma \otimes I_{n-j-2} \quad \text{for } j = 1, \dots, n - 2,$$

where I_k denotes the identity map on $(\mathbb{C}^N)^{\otimes k}$. This Hecke algebra is exactly the algebra of intertwining operators for $f_N^{\otimes n}$.

2. The Yang-Baxter operator associated with $U_q(2)$. In this note we are going to show an analogous construction for the quantum group $U_q(2)$. The paper [W] contains a construction of the quantum group $U_q(2)$, in which the crucial role is played by the function counting the number of cycles in permutations from the symmetric group S_3 . Namely, by considering the function $S_3 \ni \sigma \mapsto (-q)^{3-c(\sigma)}$, where $c(\sigma)$ is the number of cycles and $q > 0$, we constructed the following array:

$$\begin{aligned} E_{1,2,3} &= 1, & E_{1,3,2} &= E_{2,1,3} = E_{3,2,1} = -q, \\ E_{2,3,1} &= E_{3,1,2} = q^2, & E_{i,j,k} &= 0 \text{ if } \{i, j, k\} \not\subseteq \{1, 2, 3\}. \end{aligned}$$

This array defines an operator ρ on $\mathbb{C}^3 \otimes \mathbb{C}^3$ by

$$\rho : \mathbb{C}^3 \otimes \mathbb{C}^3 \ni (a, b) \mapsto \sum_{i,j,k=1}^3 E_{i,j,k} E_{k,a,b}(i, j) \in \mathbb{C}^3 \otimes \mathbb{C}^3 \quad (2.4)$$

where (a, b) denotes the standard basis element $\varepsilon_a \otimes \varepsilon_b$. In particular $\varepsilon_1 = (1, 0, 0)$, $\varepsilon_2 = (0, 1, 0)$ and $\varepsilon_3 = (0, 0, 1)$.

The definition of E implies that (2.4) simplifies to

$$\rho(a, b) = E_{a,b,k} E_{k,a,b}(a, b) + E_{b,a,k} E_{k,a,b}(b, a), \quad \text{where } \{a, b, k\} = \{1, 2, 3\} \quad (2.5)$$

for $a \neq b$ and $a, b = 1, 2, 3$. If $a = b$ then we get $\rho(a, a) = 0$. The formulas can be written explicitly as follows.

$$\begin{aligned} \rho(1, 2) &= E_{1,2,3} E_{3,1,2}(1, 2) + E_{2,1,3} E_{3,1,2}(2, 1) = q^2(1, 2) + q^3(2, 1), \\ \rho(2, 1) &= E_{2,1,3} E_{3,2,1}(2, 1) + E_{1,2,3} E_{3,2,1}(1, 2) = q^2(2, 1) + q(1, 2), \\ \rho(1, 3) &= E_{1,3,2} E_{2,1,3}(1, 3) + E_{3,1,2} E_{2,1,3}(3, 1) = q^2(1, 3) + q^3(3, 1), \\ \rho(3, 1) &= E_{3,1,2} E_{2,3,1}(3, 1) + E_{1,3,2} E_{2,3,1}(1, 3) = q^4(3, 1) + q^3(1, 3), \\ \rho(2, 3) &= E_{2,3,1} E_{1,2,3}(2, 3) + E_{3,2,1} E_{1,2,3}(3, 2) = q^2(2, 3) + q(3, 2), \\ \rho(3, 2) &= E_{3,2,1} E_{1,3,2}(3, 2) + E_{2,3,1} E_{1,3,2}(2, 3) = q^2(3, 2) + q^3(2, 3). \end{aligned}$$

Therefore, the operator $\alpha := I_2 - \frac{1}{q^2}\rho$ acts as: $\alpha(a, a) = (a, a)$ for $a = 1, 2, 3$ and

$$\begin{aligned} \alpha(1, 2) &= -q(2, 1), \\ \alpha(1, 3) &= -q(3, 1), \\ \alpha(3, 2) &= -q(2, 3), \\ \alpha(2, 1) &= -q^{-1}(1, 2), \\ \alpha(2, 3) &= -q^{-1}(3, 2), \\ \alpha(3, 1) &= (1 - q^2)(3, 1) - q(1, 3). \end{aligned} \quad (2.6)$$

This operator is not self-adjoint (which was the case for the Woronowicz’s operator σ – see (1.3)), but $\alpha^2 = (\alpha^2)^*$ is so, since

$$\begin{aligned}
 \alpha^2(1, 2) &= (2, 1), \\
 \alpha^2(2, 1) &= (2, 1), \\
 \alpha^2(2, 3) &= (3, 2), \\
 \alpha^2(3, 2) &= (2, 3), \\
 \alpha^2(1, 3) &= q^2(1, 3) - q(1 - q^2)(3, 1), \\
 \alpha^2(3, 1) &= (1 - q^2 + q^4)(3, 1) - q(1 - q^2)(1, 3).
 \end{aligned}
 \tag{2.7}$$

The first important property of α is that it is a Yang-Baxter operator.

PROPOSITION 2.1. *The operator α satisfies the Yang-Baxter equation*

$$(\alpha \otimes I)(I \otimes \alpha)(\alpha \otimes I) = (I \otimes \alpha)(\alpha \otimes I)(I \otimes \alpha).
 \tag{2.8}$$

Proof. Let $L = (\alpha \otimes I)(I \otimes \alpha)(\alpha \otimes I)$ be the left-hand side and $P = (I \otimes \alpha)(\alpha \otimes I)(I \otimes \alpha)$ be the right-hand side of (2.8). We have to show that $L(a, b, c) = P(a, b, c)$ for every $a, b, c \in \{1, 2, 3\}$ (with the notation: $(a, b, c) = \varepsilon_a \otimes \varepsilon_b \otimes \varepsilon_c$). This requires checking 27 cases. It is clear that $L(a, a, a) = P(a, a, a)$ for any $a = 1, 2, 3$. The direct calculation provides the following formulas for the other cases.

$$\begin{aligned}
 L(3, 2, 3) &= (3, 2, 3) = P(3, 2, 3), \\
 L(2, 3, 2) &= (2, 3, 2) = P(2, 3, 2), \\
 L(1, 2, 1) &= (1, 2, 1) = P(1, 2, 1), \\
 L(2, 1, 2) &= (2, 1, 2) = P(2, 1, 2), \\
 L(1, 2, 3) &= -q(3, 2, 1) = P(1, 2, 3), \\
 L(1, 3, 2) &= -q^3(2, 3, 1) = P(1, 3, 2), \\
 L(2, 1, 3) &= -q^{-1}(3, 1, 2) = P(3, 1, 2), \\
 L(3, 3, 2) &= q^2(2, 3, 3) = P(3, 3, 2), \\
 L(2, 2, 3) &= q^2(3, 2, 2) = P(2, 2, 3), \\
 L(3, 2, 2) &= q^2(2, 2, 3) = P(3, 2, 2), \\
 L(1, 1, 3) &= q^2(3, 1, 1) = P(1, 1, 3), \\
 L(1, 3, 3) &= q^2(3, 3, 1) = P(1, 3, 3), \\
 L(1, 1, 2) &= q^2(2, 1, 1) = P(1, 1, 2), \\
 L(1, 2, 2) &= q^2(2, 2, 1) = P(1, 2, 2), \\
 L(2, 3, 3) &= q^{-2}(3, 3, 2) = P(2, 3, 3), \\
 L(2, 1, 1) &= q^{-2}(1, 1, 2) = P(2, 1, 1), \\
 L(2, 2, 1) &= q^{-2}(1, 2, 2) = P(2, 2, 1).
 \end{aligned}
 \tag{2.9}$$

$$\begin{aligned}
 L(3, 2, 1) &= (1 - q^2)(3, 2, 1) - q(1, 2, 3) = P(3, 2, 1), \\
 L(3, 1, 2) &= q^2(1 - q^2)(2, 3, 1) - q^3(2, 1, 3) = P(3, 1, 2), \\
 L(2, 3, 1) &= q^{-2}(1 - q^2)(3, 1, 2) - q^{-1}(1, 3, 2) = P(2, 3, 1), \\
 L(1, 3, 1) &= -q(1 - q^2)(3, 1, 1) + q^2(1, 3, 1) = P(1, 3, 1), \\
 L(3, 1, 3) &= -q(1 - q^2)(3, 3, 1) + q^2(3, 1, 3) = P(3, 1, 3),
 \end{aligned}
 \tag{2.10}$$

$$\begin{aligned}
 L(3, 1, 1) &= (1 - q^2)(3, 1, 1) - q(1 - q^2)(1, 3, 1) + q^2(1, 1, 3) = P(3, 1, 1), \\
 L(3, 3, 1) &= (1 - q^2)(3, 3, 1) - q(1 - q^2)(3, 1, 3) + q^2(1, 3, 3) = P(3, 3, 1).
 \end{aligned}
 \tag{2.11}$$

From these formulas the Proposition follows. ■

3. The cubic Hecke algebra associated with $U_q(2)$. The second important property of the operator α is that, even though it is not a Hecke operator, it does satisfy a cubic equation, and thus it generates a *cubic Hecke algebra*. This notion has been introduced by Funar in [F], where the cubic equation $\alpha^3 - I = 0$ was considered.

PROPOSITION 3.1. *The operator α satisfies the cubic equation*

$$(\alpha^2 - I)(\alpha + q^2 \cdot I) = 0.
 \tag{3.12}$$

Proof. From the formulas (2.6), defining α it follows that it acts on the following subspaces by simple matricial formulas.

1. On the span of $(1, 2), (2, 1)$ as $\beta := \begin{pmatrix} 0 & \frac{-1}{q} \\ -q & 0 \end{pmatrix}$.
2. On the span of $(2, 3), (3, 2)$ as $\beta^* := \begin{pmatrix} 0 & -q \\ \frac{-1}{q} & 0 \end{pmatrix}$.
3. On the span of $(1, 3), (3, 1)$ as $\gamma := \begin{pmatrix} 0 & -q \\ -q & 1 - q^2 \end{pmatrix}$.
4. As identity on every (a, a) with $a = 1, 2, 3$.

It is straightforward to see that $\beta^2 - I = 0 = (\beta^*)^2 - I$. On the other hand, since

$$\gamma^2 = \begin{pmatrix} q^2 & -q(1 - q^2) \\ -q(1 - q^2) & 1 - q^2 + q^4 \end{pmatrix}$$

we obtain

$$(\gamma^2 - I)(\gamma + q^2 I) = (q^2 - 1) \begin{pmatrix} 1 & q \\ q & q^2 \end{pmatrix} \begin{pmatrix} q^2 & -q \\ -q & 1 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}.$$

Therefore both β and γ satisfy the equation (3.12), so the α does. ■

Let us define the elements

$$h_j := I_j \otimes \alpha \otimes I_{n-j-2} \quad \text{for } j = 1, \dots, n - 2,
 \tag{3.13}$$

where I_k denotes the identity map on $(\mathbb{C}^N)^{\otimes k}$. Then by Propositions 2.1 and 3.1 the elements h_1, \dots, h_n generate a cubic Hecke algebra, associated with the quantum group $U_q(2)$.

DEFINITION 3.2. The algebra $\mathcal{H}_{q,n}(2)$ generated by the elements $h_j, j = 1, 2, \dots, n$ defined by (3.13) will be called the *cubic Hecke algebra* associated with the quantum group $U_q(2)$.

The basic properties of this algebra are summarized in the following.

THEOREM 3.3. *The generators $\{h_j : 1 \leq j \leq n\}$ of $H_{q,n}(2)$ satisfy:*

$$\begin{aligned} h_j h_{j+1} h_j &= h_{j+1} h_j h_{j+1} && \text{for } j = 1, \dots, n-1, \\ h_j h_k &= h_k h_j && \text{for } |j-k| \geq 2, \\ ((h_j)^2 - 1)(h_j + q^2) &= 0 && \text{for } j = 1, \dots, n, \end{aligned} \quad (3.14)$$

The role of the Hecke algebra in the study of $SU_q(2)$ was that it was the intertwining algebra of the tensor powers of the fundamental co-representation. It is still to be checked whether the same role is played here by the $\mathcal{H}_{q,n}(2)$. In [W] the irreducible co-representations have been described, but it was not clear if the description was complete. This problem will be studied elsewhere.

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