

ANALYTICITY OF SOLUTIONS AND KOLMOGOROV'S DISSIPATION SCALE FOR 2D NAVIER-STOKES EQUATIONS

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1. Introduction. Let us consider the two-dimensional Navier-Stokes (NS) system on a torus:

$$(1) \quad \dot{u} - \nu \Delta u + (u, \nabla)u - \nabla p = \eta(t, x), \quad \operatorname{div} u = 0, \quad x = (x_1, x_2) \in \mathbf{T}^2.$$

Here $\mathbf{T}^2 = \mathbf{R}^2/2\pi\mathbf{Z}^2$, $\nu > 0$ is the viscosity, $u = u(t, x)$ is the velocity field, p is the pressure, and η is an external force. Equation (1) is supplemented with the initial condition

$$(2) \quad u(0, x) = u_0(x).$$

As is known [L], the problem (1), (2) is well-posed. Namely, for any right-hand side η and initial function u_0 that belong to appropriate functional classes there is a unique solution $u(t, x)$ for (1), (2). The aim of this article is to study analyticity of solutions regarded as functions of x and to find an asymptotic lower bound for the radius analyticity as $\nu \rightarrow 0$.

This problem is closely related to the Kolmogorov-Obukhov hypothesis on the behaviour of the energy spectrum of solutions in the turbulent regime. More precisely, let us expand a solution of (1) into the Fourier series,

$$u(t, x) = \sum_{j \in \mathbf{Z}^2} u_j(t) e^{ijx},$$

and define the energy corresponding to a wave number k by the formula

$$E_k = \sum_{k-1/2 \leq |j| \leq k+1/2} |u_j|^2.$$

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Roughly speaking, in the 2D case, the hypothesis is that there is a threshold $\lambda_{2D} = \lambda_{2D}(\nu)$ called Kolmogorov's dissipation scale such that

$$E_k \sim k^{-3} \quad \text{for } k \leq \lambda_{2D}^{-1}, \quad E_k \lesssim k^{-N} \quad \text{for } k \geq \lambda_{2D}^{-1},$$

where $N > 0$ is an arbitrary constant. Furthermore, the Kolmogorov dissipation scale is of order $\nu^{1/2}$.

Let us consider a solution for (1) that admits analytic continuation to the domain $|\operatorname{Im} x_i| \leq r\nu^\gamma$, $i = 1, 2$, where $r > 0$ is a constant. In this case, we have

$$|u_j| \leq \text{const.} e^{-r\nu^\gamma |j|},$$

whence it follows that $E_k \lesssim e^{-r\nu^\gamma k}$. Therefore,

$$E_k \lesssim k^{-N} \quad \text{for } \frac{k}{\ln k} \gtrsim \nu^{-\gamma}.$$

Hence, ignoring the logarithm, one can say that, if the Kolmogorov-Obukhov hypothesis is true, then $\lambda_{2D} \gtrsim \nu^\gamma$. Thus, *an asymptotic estimate for the radius of analyticity implies a lower bound for the Kolmogorov dissipation scale*. See [Fr, Ga, HKR, Ku] and references therein for a more detailed discussion.

The problem of analyticity of solutions for deterministic Navier-Stokes equations was studied in many papers (see, e.g., [FT, HKR]). The aim of this article is to present some results in the case when the right-hand side η is a random process analytic in the space variables and white in time. This case was investigated earlier in [M, BKL]. In particular, as is shown in [BKL], for any initial function u_0 the solution of the problem (1), (2) is analytic in x for $t > 0$, and its radius of analyticity ρ_ν can be estimated asymptotically from below by $\nu^{3+\delta}$ for any $\delta > 0$. Our estimates for solutions of (1), (2) imply that $\rho_\nu \gtrsim \nu^{2+\delta}$ for any $\delta > 0$. We note that this assertion is true for any stationary solution of Eq. (1).

2. Preliminaries. In this section, we introduce necessary functional spaces, recall the definition of a solution for Eq. (1) and the notion of a stationary solution, and formulate some known results.

Let $H^s = H^s(\mathbf{T}^2, \mathbf{R}^2)$ be the space of vector functions $u = (u_1, u_2)$ on \mathbf{T}^2 whose components belong to the Sobolev space of order s . For $s = 0$, we obtain the usual space $L^2 = L^2(\mathbf{T}^2, \mathbf{R}^2)$ with natural norm $|\cdot|$. Let H be the subspace of $u \in L^2$ such that $\operatorname{div} u = 0$ and $\int_{\mathbf{T}^2} u(x) dx = 0$ and let $\Pi: L^2 \rightarrow H$ be the orthogonal projection onto H .

Applying Π to Eq. (1), we write it in the form (see [VF])

$$(3) \quad \dot{u} + \nu Lu + B(u, u) = \eta(t),$$

where L is the restriction of the operator $-\Delta$ to H , $B(u, u) = \Pi(u, \nabla)u$ is the nonlinear term, and $\eta(t)$ is the projection of the external force. (To simplify the notation, we denote the external force and its projection to H by the same symbol.)

To describe the class of right-hand sides for (3), we introduce a trigonometric basis in H . Namely, let \mathbf{Z}' be a subset of $\mathbf{Z}_0^2 = \mathbf{Z}^2 \setminus \{0\}$ such that $\mathbf{Z}' \cup (-\mathbf{Z}') = \mathbf{Z}_0^2$ and let

$$e_j(x) = \frac{\sin(jx) j^\perp}{\sqrt{2\pi}|j|}, \quad e_{-j}(x) = \frac{\cos(jx) j^\perp}{\sqrt{2\pi}|j|}, \quad j = (j_1, j_2) \in \mathbf{Z}',$$

where $j^\perp = (-j_2, j_1)$ and $|j| = (j_1^2 + j_2^2)^{1/2}$. It is clear that $Le_j = |j|^2 e_j$, $j \in \mathbf{Z}_0^2$, and that $\{e_j, j \in \mathbf{Z}_0^2\}$ is a basis in H .

We assume that η has the form

$$(4) \quad \eta(t, x) = \frac{\partial}{\partial t} \zeta(t, x), \quad \zeta(t, x) = \sum_{j \in \mathbf{Z}_0^2} b_j \beta_j(t) e_j(x),$$

where β_j are independent standard Brownian motions defined on a complete probability $(\Omega, \mathcal{F}, \mathbf{P})$ with filtration \mathcal{F}_t and b_j are real constants satisfying the condition

$$(5) \quad \sum_{j \in \mathbf{Z}_0^2} |j|^2 b_j^2 < \infty.$$

This assumption implies, in particular, that almost all sample paths $\zeta(t, \cdot)$ belong to the space $C(\mathbf{R}_+, V)$ of continuous functions on the half-line $\mathbf{R}_+ = [0, +\infty)$ with range in $V := H \cap H^1(\mathbf{T}^2, \mathbf{R}^2)$. Equation (3) is regarded as an Itô's stochastic PDE.

Let $L_{\text{loc}}^2(\mathbf{R}_+, V)$ be the space of Bochner-measurable functions $u(t) : \mathbf{R}_+ \rightarrow V$ such that $\int_0^T \|f(t)\|^2 dt < \infty$ for any $T > 0$, where $\|u\| = |L^{1/2}u|$ is the norm in V .

DEFINITION 1. A random process $u(t) = u(t, x)$ in H defined on the half-line $t \geq 0$ and progressively measurable with respect to \mathcal{F}_t is called a *strong solution* of Eq. (3) if the following two conditions hold with probability 1.

- (i) The function $u(t, x)$ belongs to $L_{\text{loc}}^2(\mathbf{R}_+, V) \cap C(\mathbf{R}_+, H)$.
- (ii) For any $t > 0$,

$$u(t) + \int_0^t (\nu Lu + B(u, u)) ds = u(0) + \zeta(t),$$

where the left- and right-hand sides of this relation are regarded as elements of the space $H^{-1}(\mathbf{T}^2, \mathbf{R}^2)$.

A proof of the following result can be found in [VF, Chapter X] (also see [DaZ, Chapter 15]).

PROPOSITION 2. *Suppose that condition (5) holds. Then for any \mathcal{F}_0 -measurable random variable u_0 with range in H Eq. (3) has a unique solution on $[0, \infty)$ that satisfies the initial condition (2).*

Let $u(t)$ be a solution for Eq. (3) and let $\mu(t)$ be its distribution at time t . Thus, $\mu(t)$ is a probability Borel measure in the functional space H .

DEFINITION 3. The solution $u(t)$ is said to be *stationary* if $\mu(t)$ does not depend on t . In this case, $\mu(t) \equiv \mu$ is called a *stationary measure* for Eq. (3).

A proof of the following theorem can be found in [DaZ, Chapter 15].

PROPOSITION 4. *Suppose that condition (5) holds. Then Eq. (3) has a stationary measure.*

3. Main results. We begin with an estimate for a second exponential moment for stationary solutions.

THEOREM 5. *Suppose that condition (5) holds. There are positive constants σ and C not depending on ν such that, if μ is a stationary measure for Eq. (3) with some $\nu \in (0, 1]$, then*

$$\int_H \exp(\sigma\nu\|u\|^2) \mu(du) \leq C.$$

We now assume that the coefficients b_j entering the right-hand side of (3) (see (4)) satisfy the following inequality for some $\rho > 0$:

$$(6) \quad \sum_{j \in \mathbf{Z}_0^2} e^{2\rho|j|} b_j^2 < \infty.$$

Recall that, for a function $u(t, x)$ with range in H , we denote by $u_j(t)$ its Fourier coefficients. The following theorem establishes the analyticity of solutions for the Cauchy problem (3), (2) and gives an asymptotic lower bound for the radius of analyticity.

THEOREM 6. *Suppose that condition (6) holds. Let $u_0 = u_0^\nu(x)$ be a family of random initial functions that satisfy the inequality*

$$\mathbf{E} \exp(\sigma\nu\|u_0^\nu\|^2) \leq R \quad \text{for } 0 < \nu \leq 1,$$

where the positive constants σ and R do not depend on ν . Then for any $t_0 \geq 1$, $T > 0$ and $\delta \in (0, 1]$ there are positive random variables $r_\nu = r_\nu(t_0, T, \delta)$ and $C_\nu = C_\nu(t_0, T, \delta)$ such that, with probability 1,

$$|u_j(t)| \leq C_\nu e^{-r_\nu \nu^{2+\delta}|j|}, \quad t_0 \leq t \leq t_0 + T, \quad j \in \mathbf{Z}_0^2.$$

Moreover, for any integer $m \geq 1$ there is a constant $K_m = K_m(\sigma, T, \delta, R) > 0$ not depending on t_0 and ν such that

$$\mathbf{E} r_\nu^{-m} \leq K_m, \quad \mathbf{E} C_\nu^m \leq K_m \nu^{-m/2} \quad \text{for } 0 < \nu \leq 1.$$

In particular, any solution of Eq. (3) with deterministic initial function $u_0 \in V$ is analytic in x with probability 1, and its radius of analyticity can be estimated from below by $r_\nu \nu^{2+\delta}$. Theorem 5 implies that the above assertions are valid for any stationary solution of Eq. (3).

Proofs of Theorems 5 and 6 are given in [S].

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