

*THE NONLINEAR NEUMANN PROBLEM
AND SHARP WEIGHTED SOBOLEV INEQUALITIES*

BY

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Abstract. We prove sharp inequalities in weighted Sobolev spaces. Our approach is based on the blow-up technique applied to some nonlinear Neumann problems.

1. Introduction. The main purpose of this work is to prove some new weighted Sobolev inequalities. These inequalities are obtained by analyzing the asymptotic behaviour of solutions of nonlinear Neumann problems involving the critical Sobolev exponent.

Let S denote the best Sobolev constant for the embedding of $D^{1,2}(\mathbb{R}^N)$ into $L^{2^*}(\mathbb{R}^N)$, where $2^* = 2N/(N-2)$, $N \geq 3$, that is,

$$(1) \quad S = \inf \left\{ \frac{\int_{\mathbb{R}^N} |\nabla u|^2 dx}{\left(\int_{\mathbb{R}^N} |u|^{2^*} dx\right)^{2/2^*}} : u \in D^{1,2}(\mathbb{R}^N) \setminus \{0\} \right\}.$$

It is known that the constant S is attained by

$$U(x) = \left[\frac{N(N-2)}{N(N-2) + |x|^2} \right]^{(N-2)/2}.$$

The function U , called an *instanton*, satisfies the equation

$$-\Delta U = U^{2^*-1} \quad \text{in } \mathbb{R}^N.$$

We also have $\int_{\mathbb{R}^N} |\nabla U|^2 dx = \int_{\mathbb{R}^N} U^{2^*} dx = S^{N/2}$.

Let $\varepsilon > 0$ and $y \in \mathbb{R}^N$. We define

$$U_{\varepsilon,y}(x) = \varepsilon^{-(N-2)/2} U((x-y)/\varepsilon).$$

Then any minimizer for S is of the form $U_{\varepsilon,y}$.

Let $\Omega \subset \mathbb{R}^N$ be a bounded domain with a smooth boundary $\partial\Omega$. Let Q denote a Hölder continuous and positive function defined on $\bar{\Omega}$. Also, let $Q_M = \max_{x \in \bar{\Omega}} Q(x)$ and $Q_m = \max_{x \in \partial\Omega} Q(x)$. In what follows we write $p+1 = 2^*$. In this paper the following inequalities are proved:

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(I) Let $N \geq 5$. Suppose that $Q_M < 2^{2/(N-2)}Q_m$. Then there exists a constant $\Lambda_1(\Omega) > 0$ such that

$$\left(\int_{\Omega} Q(x)|u|^{p+1} dx \right)^{2/(p+1)} \leq \frac{2^{2/N}Q_m^{(N-2)/N}}{S} \int_{\Omega} |\nabla u|^2 dx + \Lambda_1(\Omega) \int_{\partial\Omega} u^2 dx$$

for every $u \in H^1(\Omega)$.

(II) Let $N \geq 4$ and $\tau = 2N/(N - 1)$. Suppose that $Q_M \leq 2^{2/(N-2)}Q_m$. Then there exists a constant $\Lambda_2(\Omega) > 0$ such that

$$\begin{aligned} \left(\int_{\Omega} Q(x)|u|^{p+1} dx \right)^{2/(p+1)} &\leq \frac{2^{2/N}Q_m^{(N-2)/N}}{S} \int_{\Omega} |\nabla u|^2 dx + \Lambda_2(\Omega) \left(\int_{\Omega} |u|^{\tau} dx \right)^{2/\tau} \end{aligned}$$

for every $u \in H^1(\Omega)$.

(III) Suppose $Q_M > 2^{2/(N-2)}Q_m$. Let $N \geq 5$ and $2 \leq \tau \leq 2N/(N - 1)$, or $N = 4$ and $2 < \tau \leq 2N/(N - 1)$. Then there exists a constant $\Lambda_3(\Omega) > 0$ such that

$$\left(\int_{\Omega} Q(x)|u|^{p+1} dx \right)^{2/(p+1)} \leq \frac{Q_M^{(N-2)/N}}{S} \int_{\Omega} |\nabla u|^2 dx + \Lambda_3(\Omega) \left(\int_{\Omega} |u|^{\tau} dx \right)^{2/\tau}$$

for every $u \in H^1(\Omega)$.

These inequalities should be compared with the following ones, established in the papers [6], [26]:

(A) There exists a constant $\lambda(\Omega) \geq ((N - 2)/2)H(\Omega)$ such that

$$\left(\int_{\Omega} |u|^{p+1} dx \right)^{2/(p+1)} \leq \frac{2^{2/N}}{S} \int_{\Omega} |\nabla u|^2 dx + \lambda(\Omega) \int_{\partial\Omega} u^2 dx$$

for every $u \in H^1(\Omega)$, where $H(\Omega) = \max_{x \in \partial\Omega} H(x)$ and $H(x)$ denotes the mean curvature at $x \in \partial\Omega$.

(B) Let $\tau = 2N/(N - 1)$. Then there exists a constant $\tilde{\lambda}(\Omega) > 0$ such that

$$\left(\int_{\Omega} |u|^{p+1} dx \right)^{2/(p+1)} \leq \frac{2^{2/N}}{S} \int_{\Omega} |\nabla u|^2 dx + \tilde{\lambda}(\Omega) \left(\int_{\Omega} |u|^{\tau} dx \right)^{2/\tau}$$

for every $u \in H^1(\Omega)$.

It is evident that none of (I)–(III) is a direct consequence of (A) and (B). We also point out that an inequality of type (III) in the case $Q_M \leq 2^{2/(N-2)}Q_m$ with $2 \leq \tau < 2N/(N - 1)$ is not possible. This will be clear from our analysis (see Proposition 2.7). The inequalities (I)–(III) will be

established by applying a blow-up technique to solutions of the following Neumann problems:

$$(1_\lambda) \quad \begin{cases} -\Delta u = Q(x)u^p & \text{in } \Omega, \\ \partial u / \partial \nu + \lambda u = 0 & \text{on } \partial \Omega, \end{cases}$$

and

$$(1_{\lambda,\tau}) \quad \begin{cases} -\Delta u + \lambda (\int_\Omega |u|^\tau dx)^{2/\tau-1} u^{\tau-1} = Q(x)u^p & \text{in } \Omega, \\ \partial u / \partial \nu = 0 & \text{on } \partial \Omega, \end{cases}$$

where ν is the outward normal to $\partial \Omega$, $2 \leq \tau \leq 2N/(N - 1)$.

The proofs of the three inequalities are similar and proceed by contradiction. One assumes that least energy solutions of problems (1_λ) and $(1_{\lambda,\tau})$ exist for all positive λ and shows that they are close to some instantons. This enables one to give a lower bound for the energy of the solutions and to arrive at a contradiction.

In the case corresponding to inequality (I) the instantons concentrate at the boundary of Ω and therefore we can apply the arguments used in the proof of inequality (A). In the case corresponding to inequality (III) the instantons concentrate in the interior of Ω and the estimates are slightly different from the ones in the proof of inequality (B). In the case corresponding to inequality (II) the instantons either concentrate in the interior of Ω , or concentrate on the boundary of Ω , in which case the estimates are similar to those of (B).

This paper is organized as follows. Section 2 is concerned with the existence of least energy solutions of problems (1_λ) and $(1_{\lambda,\tau})$. In Section 3 we prove inequality (I) and in Section 4 we prove inequalities (II) and (III).

It is natural to ask if there exists a constant $\Lambda_4(\Omega) > 0$ such that

$$\left(\int_\Omega Q(x)|u|^{p+1} dx \right)^{2/(p+1)} \leq \frac{Q_M^{(N-2)/N}}{S} \int_\Omega |\nabla u|^2 dx + \Lambda_4(\Omega) \int_{\partial \Omega} u^2 dx,$$

for all $u \in H^1(\Omega)$, if $Q_M \geq 2^{2/(N-2)}Q_m$. We have not been able to answer this question.

2. Solvability of problems (1_λ) and $(1_{\lambda,\tau})$. Solutions of problems (1_λ) and $(1_{\lambda,\tau})$ will be obtained as minimizers on $H^1(\Omega) \setminus \{0\}$ of the functionals

$$J_\lambda(u) = \frac{\int_\Omega |\nabla u|^2 dx + \lambda \int_{\partial \Omega} u^2 dx}{\left(\int_\Omega Q(x)|u|^{p+1} dx \right)^{2/(p+1)}}$$

and

$$J_{\lambda,\tau}(u) = \frac{\int_\Omega |\nabla u|^2 dx + \lambda (\int_\Omega |u|^\tau dx)^{2/\tau}}{\left(\int_\Omega Q(x)|u|^{p+1} dx \right)^{2/(p+1)}}.$$

We set

$$\begin{aligned}
 S_\lambda &= \inf\{J_\lambda(u) : u \in H^1(\Omega) \setminus \{0\}\} \\
 &= \inf\left\{ \int_\Omega |\nabla u|^2 dx + \lambda \int_{\partial\Omega} u^2 dx : u \in V_Q \right\}
 \end{aligned}$$

and

$$\begin{aligned}
 S_{\lambda,\tau} &= \inf\{J_{\lambda,\tau}(u) : u \in H^1(\Omega) \setminus \{0\}\} \\
 &= \inf\left\{ \int_\Omega |\nabla u|^2 dx + \lambda \left(\int_\Omega |u|^\tau dx \right)^{2/\tau} : u \in V_Q \right\},
 \end{aligned}$$

where $V_Q = \{u \in H^1(\Omega) : \int_\Omega Q(x)|u|^{p+1} dx = 1\}$.

To show that S_λ and $S_{\lambda,\tau}$ are achieved we need the following version of P. L. Lions' [14] concentration-compactness principle. Let $\{u_n\} \subset H^1(\Omega)$ be a weakly convergent sequence to u in $H^1(\Omega)$ and such that $|u_n|^{p+1} \rightharpoonup \mu$ and $|\nabla u_n|^2 \rightharpoonup \tilde{\mu}$ weakly in the sense of measures. Then there exist numbers $\mu_j > 0$, $\tilde{\mu}_j > 0$ and points $x_j \in \bar{\Omega}$, $j \in J$, where J is at most a countable set, such that

$$\mu = |u|^{p+1} + \sum_{j \in J} \mu_j \delta_{x_j}, \quad \tilde{\mu} \geq |\nabla u|^2 + \sum_{j \in J} \tilde{\mu}_j \delta_{x_j}.$$

Moreover, if $x_j \in \Omega$, then

$$(2) \quad S(\mu_j)^{(N-2)/N} \leq \tilde{\mu}_j,$$

and if $x_j \in \partial\Omega$, then

$$(3) \quad \frac{S}{2^{2/N}} (\mu_j)^{(N-2)/N} \leq \tilde{\mu}_j.$$

The following lemmas give criteria for the existence of minima for J_λ and $J_{\lambda,\tau}$.

LEMMA 2.1. *If*

$$(4) \quad S_\lambda < \frac{S}{2^{2/N} Q_m^{(N-2)/N}}$$

and

$$(5) \quad Q_M < 2^{2/(N-2)} Q_m,$$

then S_λ is achieved.

Proof. Let $\{u_n\} \subset H^1(\Omega)$ be such that $\int_\Omega Q|u_n|^{p+1} = 1$ for each n , and $\int_\Omega |\nabla u_n|^2 dx + \lambda \int_{\partial\Omega} u_n^2 dx \rightarrow S_\lambda$ as $n \rightarrow \infty$. We may assume that $u_n \rightharpoonup u$ in $H^1(\Omega)$, $u_n \rightarrow u$ in $L^2(\Omega)$ and $u_n \rightarrow u$ a.e. on Ω . Applying the concentration-compactness principle we can write

$$\int_\Omega Q|u|^{p+1} + \sum_{j \in J} Q(x_j)\mu_j = 1$$

and

$$S_\lambda = \lim_{n \rightarrow \infty} \left(\int_{\Omega} |\nabla u_n|^2 dx + \lambda \int_{\partial\Omega} u_n^2 dx \right) \geq \int_{\Omega} |\nabla u|^2 dx + \lambda \int_{\partial\Omega} u^2 dx + \sum_{j \in J} \tilde{\mu}_j.$$

Using (2), (3) and (5) we derive the following estimate from below for S_λ :

$$\begin{aligned} S_\lambda &\geq \int_{\Omega} |\nabla u|^2 dx + \lambda \int_{\partial\Omega} u^2 dx + \sum_{x_j \in \Omega} \tilde{\mu}_j + \sum_{x_j \in \partial\Omega} \tilde{\mu}_j \\ &\geq S_\lambda \left(\int_{\Omega} Q|u|^{p+1} dx \right)^{2/(p+1)} + \sum_{x_j \in \Omega} \frac{S}{Q(x_j)^{(N-2)/N}} (\mu_j Q(x_j))^{(N-2)/N} \\ &\quad + \sum_{x_j \in \partial\Omega} \frac{S}{2^{2/N} Q(x_j)^{(N-2)/N}} (\mu_j Q(x_j))^{(N-2)/N} \\ &\geq S_\lambda \left(\int_{\Omega} Q|u|^{p+1} dx \right)^{2/(p+1)} + \sum_{x_j \in \Omega} \frac{S}{Q_M^{(N-2)/N}} (Q(x_j) \mu_j)^{(N-2)/N} \\ &\quad + \sum_{x_j \in \partial\Omega} \frac{S}{2^{2/N} Q_m^{(N-2)/N}} (Q(x_j) \mu_j)^{(N-2)/N} \\ &\geq S_\lambda \left(\int_{\Omega} Q|u|^{p+1} dx \right)^{2/(p+1)} + \sum_{j \in J} \frac{S}{2^{2/N} Q_m^{(N-2)/N}} (Q(x_j) \mu_j)^{(N-2)/N}. \end{aligned}$$

Since S_λ satisfies (4) we must have $\mu_j = 0$ for all $j \in J$ and the result follows. ■

LEMMA 2.2. *If*

$$(6) \quad S_\lambda < \frac{S}{Q_M^{(N-2)/N}}$$

and

$$(7) \quad Q_M \geq 2^{2/(N-2)} Q_m,$$

then S_λ is achieved.

The same method can be used to obtain conditions guaranteeing the solvability of problem $(1_{\lambda,\tau})$.

LEMMA 2.3. *If $2 \leq \tau \leq 2N/(N-1)$,*

$$(8) \quad S_{\lambda,\tau} < \frac{S}{2^{2/N} Q_m^{(N-2)/N}}$$

and (5) holds, then $S_{\lambda,\tau}$ is achieved.

LEMMA 2.4. *If $2 \leq \tau \leq 2N/(N-1)$,*

$$(9) \quad S_{\lambda,\tau} < \frac{S}{Q_M^{(N-2)/N}}$$

and (7) holds, then $S_{\lambda,\tau}$ is achieved.

If $Q \equiv 1$, the functionals J_λ and $J_{\lambda,\tau}$ will be denoted by I_λ and $I_{\lambda,\tau}$, respectively.

Adimurthi and Mancini [1] proved that if $x_o \in \partial\Omega$, then

$$(10) \quad I_\lambda(U_{\varepsilon,x_o}) = \frac{S}{2^{2/N}} - A_N \left(\frac{N-2}{2} H(x_o) - \lambda \right) \varepsilon + O(\varepsilon^2), \quad N \geq 5,$$

where $A_N > 0$ is a constant depending on N . Using this asymptotic formula we can give a condition on Q guaranteeing the validity of the inequality (4) in Lemma 2.1.

For this we need the following assumption on Q :

(Q) There exists a point $x_o \in \partial\Omega$ such that

$$Q(x_o) = Q_m \quad \text{and} \quad |Q(x) - Q(x_o)| = o(|x - x_o|)$$

for x near x_o .

If (Q) holds, we have

$$J_\lambda(U_{\varepsilon,x_o}) = Q(x_o)^{-(N-2)/N} I_\lambda(U_{\varepsilon,x_o}) + o(\varepsilon)$$

and it follows from (10) that

$$(11) \quad J_\lambda(U_{\varepsilon,x_o}) = \frac{S}{2^{2/N} Q_m^{(N-2)/N}} - \frac{A_N}{Q_m^{(N-2)/N}} \left(\frac{N-2}{2} H(x_o) - \lambda \right) \varepsilon + o(\varepsilon).$$

We now observe that if $H(x_o) > 0$ and $\lambda < ((N-2)/2)H(x_o)$, then for sufficiently small $\varepsilon > 0$ we have

$$J_\lambda(U_{\varepsilon,x_o}) < \frac{S}{2^{2/N} Q_m^{(N-2)/N}},$$

which shows that (4) is satisfied for $\lambda < ((N-2)/2)H(x_o)$.

PROPOSITION 2.5. *Let $N \geq 5$. Suppose that $Q_M \leq 2^{2/(N-2)} Q_m$ and (Q) holds with $H(x_o) > 0$. Then problem (1_λ) has a solution for $\lambda < ((N-2)/2)H(x_o)$. Moreover,*

$$\lim_{\lambda \rightarrow \infty} S_\lambda = \frac{S}{2^{2/N} Q_m^{(N-2)/N}}.$$

Proposition 2.5 also holds for $N = 3, 4$ if one uses a suitable modification of (10) ([1], [5]). Note that from the above discussion it is obvious that under the assumption of Proposition 2.5, $S_\lambda \leq S/(2^{2/N} Q_m^{(N-2)/N})$ for all $\lambda > 0$ and $S_\lambda < S/(2^{2/N} Q_m^{(N-2)/N})$ for $\lambda < ((N-2)/2)H(\Omega)$. The proof of the result on the asymptotic behaviour of S_λ is standard.

In the case $Q_M > 2^{2/(N-2)} Q_m$, condition (6) from Lemma 2.2 is difficult to check. Obviously, it is satisfied for small $\lambda > 0$. By testing J_λ with $U_{\varepsilon,y}$, where $Q(y) = Q_M$, we easily show that $S_\lambda \leq S/Q_M^{(N-2)/N}$ for all $\lambda > 0$.

LEMMA 2.6. *If $Q_M \geq 2^{2/(N-2)}Q_m$, then the condition $S_\lambda < S/Q_M^{(N-2)/N}$ is satisfied for small $\lambda > 0$ and*

$$\lim_{\lambda \rightarrow \infty} S_\lambda = \frac{S}{Q_M^{(N-2)/N}}.$$

We now turn our attention to the functional $J_{\lambda,\tau}$ and problem $(1_{\lambda,\tau})$. By testing $I_{\lambda,\tau}$ with U_{ε,x_o} , for $N \geq 5$ we get

$$\begin{aligned} I_{\lambda,\tau}(U_{\varepsilon,x_o}) &= S/2^{2/N} - A_N H(x_o)\varepsilon + \lambda B \varepsilon^{2N/\tau - (N-2)} \\ &\quad + o(\lambda \varepsilon^{2N/\tau - (N-2)}) + o(\varepsilon^2), \end{aligned}$$

where $B > 0$ is a constant depending on N and τ . If (Q) holds and $H(x_o) > 0$, then

$$J_{\lambda,\tau}(U_{\varepsilon,x_o}) = Q_m^{-(N-2)/N} I_{\lambda,\tau}(U_{\varepsilon,x_o}) + o(\varepsilon).$$

Hence the condition (8) of Lemma 2.3 is satisfied for every $\lambda > 0$ provided $2 \leq \tau < 2N/(N-1)$.

PROPOSITION 2.7. *Let $N \geq 5$ and $Q_M \leq 2^{2/(N-2)}Q_m$. Suppose that (Q) holds, $H(x_o) > 0$ and $2 \leq \tau < 2N/(N-1)$. Then problem $(1_{\lambda,\tau})$ has a solution for each $\lambda > 0$. Moreover,*

$$\lim_{\lambda \rightarrow \infty} S_{\lambda,\tau} = \frac{S}{2^{2/N} Q_M^{N/(N-2)}} \quad \text{for } 2 \leq \tau \leq \frac{2N}{N-1}.$$

As in the case of Proposition 2.5, this continues to hold for $N = 3, 4$. Finally,

LEMMA 2.8. *If $Q_M \geq 2^{2/(N-2)}Q_m$, then the condition $S_{\lambda,\tau} < S/Q_M^{(N-2)/N}$ is satisfied for small $\lambda > 0$ and*

$$\lim_{\lambda \rightarrow \infty} S_{\lambda,\tau} = \frac{S}{Q_M^{(N-2)/N}} \quad \text{for } 2 \leq \tau \leq \frac{2N}{N-1}.$$

To end this section we observe that from Propositions 2.5 and 2.7 and Lemma 2.8 we can deduce a weak form of inequalities (I)–(III). Namely, given a $\delta > 0$, there exist $\lambda_1 = \lambda_1(\Omega)$, $\lambda_2 = \lambda_2(\Omega)$ and $\lambda_3 = \lambda_3(\Omega)$ such that, for every $u \in H^1(\Omega)$,

$$\begin{aligned} (12) \quad & \left(\int_{\Omega} Q(x)|u|^{p+1} dx \right)^{2/(p+1)} \\ & \leq \left(\frac{2^{2/N} Q_m^{(N-2)/N}}{S} + \delta \right) \int_{\Omega} |\nabla u|^2 dx + \lambda_1 \int_{\partial\Omega} u^2 dx, \end{aligned}$$

and

$$(13) \quad \left(\int_{\Omega} Q(x)|u|^{p+1} dx \right)^{2/(p+1)} \leq \left(\frac{2^{2/N} Q_m^{(N-2)/N}}{S} + \delta \right) \int_{\Omega} |\nabla u|^2 dx + \lambda_2 \left(\int_{\Omega} |u|^{\tau} dx \right)^{2/\tau}$$

if $Q_M \leq 2^{2/(N-2)} Q_m$, and

$$(14) \quad \left(\int_{\Omega} Q(x)|u|^{p+1} dx \right)^{2/(p+1)} \leq \left(\frac{Q_M^{(N-2)/N}}{S} + \delta \right) \int_{\Omega} |\nabla u|^2 dx + \lambda_3 \left(\int_{\Omega} |u|^{\tau} dx \right)^{2/\tau}$$

if $Q_M \geq 2^{2/(N-2)} Q_m$.

3. Proof of inequality (I). The proof of inequality (I) is by contradiction. Throughout this section we suppose that inequality (5) is satisfied. Assume that, for each $\lambda > 0$, $S_{\lambda} < S/(2^{2/N} Q_m^{(N-2)/N})$. Let $\lambda_k \rightarrow \infty$. For each k there is a minimizer $u_k = u_{\lambda_k}$ of J_{λ_k} with $\int_{\Omega} Q u_k^{p+1} dx = 1$. It satisfies

$$(1_{u_k}) \quad \begin{cases} -\Delta u_k = S_{\lambda_k} Q u_k^p & \text{in } \Omega, \\ \partial u_k / \partial \nu + \lambda_k u_k = 0 & \text{on } \partial \Omega. \end{cases}$$

Our aim is to show that u_k is close to some instanton U_{ε_k, P_k} with $P_k \rightarrow P_{\circ}$, $Q_m = Q(P_{\circ})$. This in turn will contradict the inequality $S_{\lambda_k} < S/(2^{2/N} Q_m^{(N-2)/N})$.

We start by setting

$$M_k := \max_{\bar{\Omega}} u_k = u_k(P_k)$$

for some $P_k \in \bar{\Omega}$.

LEMMA 3.1. $M_k \rightarrow \infty$ and $u_k \rightarrow 0$ in $H^1(\Omega)$.

Proof. Since $\{u_k\}$ is bounded in $H^1(\Omega)$ we may assume that $u_k \rightharpoonup u$ in $H^1(\Omega)$ and from (4) we deduce that $u_k \rightarrow 0$ in $L^2(\partial \Omega)$. Hence $u \in H^1_{\circ}(\Omega)$.

Assume that M_k is bounded. Then $\int_{\Omega} Q u^{p+1} dx = 1$ and by the lower semicontinuity of the norm with respect to weak convergence, we have

$$\int_{\Omega} |\nabla u|^2 dx \leq \frac{S}{2^{2/N} Q_m^{(N-2)/N}}.$$

Since $u \in H^1_{\circ}(\Omega)$ we also have

$$\frac{\int_{\Omega} |\nabla u|^2 dx}{(\int_{\Omega} Qu^{p+1} dx)^{2/(p+1)}} \geq \frac{S}{Q_M^{(N-2)/N}} > \frac{S}{2^{2/N} Q_m^{(N-2)/N}},$$

which is impossible. This shows that the sequence $\{M_k\}$ is unbounded.

Also, $u_k \rightharpoonup 0$ in $H^1(\Omega)$, otherwise its weak limit u satisfies $0 < \int_{\Omega} Qu^{p+1} dx \leq 1$. The case $\int_{\Omega} Qu^{p+1} dx = 1$ is excluded by the above argument. If $0 < \int_{\Omega} Qu^{p+1} dx < 1$, we get a contradiction by applying the concentration-compactness principle. ■

Let

$$\varepsilon_k = M_k^{-2/(N-2)}.$$

LEMMA 3.2. $\lambda_k \int_{\partial\Omega} u_k^2 \rightarrow 0$ and $\lambda_k \varepsilon_k \rightarrow 0$.

Proof. Applying inequality (12) to u_k , we get

$$1 \leq \left(\frac{2^{2/N} Q_m^{(N-2)/N}}{S} + \delta \right) \int_{\Omega} |\nabla u_k|^2 dx + \lambda_1 \int_{\partial\Omega} u_k^2 dx.$$

Since $\int_{\partial\Omega} u_k^2 dx \rightarrow 0$ we obtain

$$\begin{aligned} 1 &\leq \left(\frac{2^{2/N} Q_m^{(N-2)/N}}{S} + \delta \right) \lim_{k \rightarrow \infty} \int_{\Omega} |\nabla u_k|^2 dx \\ &\leq \left(\frac{2^{2/N} Q_m^{(N-2)/N}}{S} + \delta \right) \lim_{k \rightarrow \infty} J_{\lambda_k}(u_k) \\ &= \left(\frac{2^{2/N} Q_m^{(N-2)/N}}{S} + \delta \right) \frac{S}{2^{2/N} Q_m^{(N-2)/N}}. \end{aligned}$$

Since $\delta > 0$ is arbitrary we see that

$$\lim_{k \rightarrow \infty} \int_{\Omega} |\nabla u_k|^2 dx = \frac{S}{2^{2/N} Q_m^{(N-2)/N}}$$

and $\lim_{k \rightarrow \infty} \lambda_k \int_{\partial\Omega} u_k^2 dx = 0$.

To prove the second assertion of the lemma, note that

$$\begin{aligned} \lambda_k \int_{\partial\Omega} u_k^2 dx &= \lambda_k \varepsilon_k \int_{\partial\Omega} M_k^{2/(N-2)} u_k^2 dx \\ &= \lambda_k \varepsilon_k \int_{\partial\Omega} M_k^q \frac{u_k^2}{M_k^2} dx \geq \lambda_k \varepsilon_k \int_{\partial\Omega} u_k^q dx, \end{aligned}$$

where $q = 2(N - 1)/(N - 2)$. Therefore to complete the proof of the second assertion it is sufficient to show that

$$\liminf_{k \rightarrow \infty} \int_{\partial\Omega} u_k^q dx > 0.$$

We follow the argument used in [13] (see the proof of inequality (2.6) there). In the contrary case, assume that $\lim_{k \rightarrow \infty} \int_{\partial\Omega} u_k^q dx = 0$. Then we may assume that $u_k \rightharpoonup \bar{u}$ in $H^1_0(\Omega)$, up to a subsequence. Using the Brézis–Lieb lemma [8] we have

$$\begin{aligned}
 1 &\equiv \int_{\Omega} Q(x)u_k^{p+1} dx = \int_{\Omega} Q(x)|u_k - \bar{u}|^{p+1} dx + \int_{\Omega} Q(x)|\bar{u}|^{p+1} dx + o(1), \\
 &\int_{\Omega} Q(x)\bar{u}^{p+1} dx \leq 1, \\
 &\int_{\Omega} Q(x)|u_k - \bar{u}|^{p+1} dx \leq 1 + o(1).
 \end{aligned}$$

Using the last three relations and the inequality (see [9], inequality (1.9))

$$\left(\int_{\Omega} Q(x)|u|^{p+1} dx \right)^{2/(p+1)} \leq \left(\frac{Q_M^{(N-2)/N}}{S} + \varepsilon \right) \int_{\Omega} |\nabla u|^2 dx + \frac{C}{\varepsilon} \left(\int_{\partial\Omega} |u|^q dx \right)^{2/q}$$

for $u \in H^1(\Omega)$, where $C > 0$ is a constant, we easily deduce that for every $\delta > 0$ we have

$$\begin{aligned}
 S_{\lambda_k} &= \int_{\Omega} |\nabla u_k|^2 dx + \lambda_k \int_{\partial\Omega} u_k^2 dx \\
 &= \int_{\Omega} |\nabla(u_k - \bar{u})|^2 dx + \int_{\Omega} |\nabla \bar{u}|^2 dx + \lambda_k \int_{\partial\Omega} u_k^2 dx + o(1) \\
 &\geq \left(\frac{S}{Q_M^{(N-2)/N}} - \delta \right) \left(\int_{\Omega} Q(x)|u_k - \bar{u}|^{p+1} dx \right)^{2/(p+1)} \\
 &\quad + S_{\lambda_k} \left(\int_{\Omega} Q(x)\bar{u}^{p+1} dx \right)^{2/(p+1)} + \lambda_k \int_{\partial\Omega} u_k^2 dx + o(1) \\
 &\geq \left(\frac{S}{Q_M^{(N-2)/N}} - \delta \right) \int_{\Omega} Q(x)|u_k - \bar{u}|^{p+1} dx \\
 &\quad + S_{\lambda_k} \int_{\Omega} Q(x)\bar{u}^{p+1} dx + \lambda_k \int_{\partial\Omega} u_k^2 dx + o(1) \\
 &= \left(\frac{S}{Q_M^{(N-2)/N}} - \delta - S_{\lambda_k} \right) \int_{\Omega} Q(x)|u_k - \bar{u}|^{p+1} dx \\
 &\quad + S_{\lambda_k} + \lambda_k \int_{\partial\Omega} u_k^2 dx + o(1).
 \end{aligned}$$

Since

$$Q_M^{(N-2)/N} < 2^{2/N} Q_m^{(N-2)/N} = \frac{S}{\lim_{k \rightarrow \infty} S_{\lambda_k}},$$

choosing $\delta > 0$ sufficiently small, we deduce that

$$\lim_{k \rightarrow \infty} \int_{\Omega} Q(x) |u_k - \bar{u}|^{p+1} dx = 0$$

and consequently $\int_{\Omega} Q(x) \bar{u}^{p+1} dx = 1$. This means that

$$\begin{aligned} \frac{\int_{\Omega} |\nabla \bar{u}|^2 dx}{\left(\int_{\Omega} Q(x) \bar{u}^{p+1} dx\right)^{2/(p+1)}} &\geq \frac{\int_{\Omega} |\nabla \bar{u}|^2 dx}{Q_M^{(N-2)/N} \left(\int_{\Omega} \bar{u}^{p+1} dx\right)^{2/(p+1)}} \\ &> \frac{S}{Q_M^{(N-2)/N}} > \frac{S}{2^{2/N} Q_m^{(N-2)/N}}, \end{aligned}$$

since $\bar{u} \in H_0^1(\Omega)$. On the other hand, by the lower semicontinuity of the norm with respect to weak convergence, we have

$$\int_{\Omega} |\nabla \bar{u}|^2 dx \leq \frac{S}{2^{2/N} Q_m^{(N-2)/N}},$$

which is impossible. Therefore $\lambda_k \varepsilon_k \rightarrow 0$. ■

LEMMA 3.3. *Up to a subsequence, $P_k \rightarrow P_o$, where P_o is such that $Q(P_o) = Q_m$. Moreover, $P_k \in \partial\Omega$ for large k ,*

$$(15) \quad \lim_{k \rightarrow \infty} \int_{\Omega} \left| \nabla \left(u_k - \varepsilon_k^{-(N-2)/2} U \left(\frac{S^{1/2} Q_m^{1/N}}{2^{1/N}} \frac{\cdot - P_k}{\varepsilon_k} \right) \right) \right|^2 dx = 0,$$

$$(16) \quad \lim_{k \rightarrow \infty} \int_{\Omega} \left| u_k - \varepsilon_k^{-(N-2)/2} U \left(\frac{S^{1/2} Q_m^{1/N}}{2^{1/N}} \frac{\cdot - P_k}{\varepsilon_k} \right) \right|^{p+1} dx = 0.$$

Proof. Let $v_k(x) = \varepsilon_k^{(N-2)/2} u_k(\varepsilon_k x + P_k)$ for $x \in \Omega_k = (\Omega - P_k)/\varepsilon_k$. Then the functions v_k are solutions of the Neumann problems

$$(1v_k) \quad \begin{cases} -\Delta v_k = S_{\lambda_k} Q(\varepsilon_k x + P_k) v_k^p & \text{in } \Omega_k, \\ \partial v_k / \partial \nu + \lambda_k \varepsilon_k v_k = 0 & \text{on } \partial\Omega_k, \\ 0 \leq v_k(x) \leq 1 & \text{in } \Omega_k, \\ v_k(0) = 1. \end{cases}$$

Passing to a subsequence, we can assume that $P_k \rightarrow P_o$ for some $P_o \in \bar{\Omega}$, and $\text{dist}(P_k, \partial\Omega)/\varepsilon_k$ converges in the extended real line. Using elliptic regularity theory, we show that, up to a subsequence, $v_k \rightarrow \omega$ in $C_{\text{loc}}^2(\Omega_{\infty})$, where $\Omega_{\infty} = \lim_{k \rightarrow \infty} \Omega_k$. Thus ω satisfies

$$\begin{cases} -\Delta \omega = \widehat{S} Q(P_o) \omega^p & \text{in } \Omega_{\infty}, \\ \partial \omega / \partial \nu = 0 & \text{on } \partial\Omega_{\infty}, \\ 0 \leq \omega \leq 1 & \text{in } \Omega_{\infty}, \\ \omega(0) = 1, \end{cases}$$

where $\widehat{S} = S / (2^{2/N} Q_m^{(N-2)/N})$.

We now distinguish two cases: (i) $\text{dist}(P_k, \partial\Omega)/\varepsilon_k$ converges to ∞ , and (ii) $\text{dist}(P_k, \partial\Omega)/\varepsilon_k$ converges to a real number.

If case (i) occurs then $\Omega_\infty = \mathbb{R}^N$. By [12] (proof of Theorem 2.3 on p. 34) we see that $\omega(x) = U(\beta x)$, where $\beta^2 = Q(P_\circ)\widehat{S}$. Since

$$\int_{\Omega_k} |\nabla v_k|^2 dx \rightarrow \frac{S}{2^{2/N} Q_m^{(N-2)/N}},$$

this yields

$$\beta^{2-N} S^{N/2} = \int_{\mathbb{R}^N} |\nabla \omega|^2 dx \leq \lim_{k \rightarrow \infty} \int_{\Omega_k} |\nabla v_k|^2 dx = \frac{S}{2^{2/N} Q_m^{(N-2)/N}}.$$

Thus, $2Q_m^{(N-2)/2^* + (N-2)/2} \leq Q(P_\circ)^{(N-2)/2}$. We must have $2^{2/(N-2)} Q_m \leq Q(P_\circ) \leq Q_M$, which is impossible. Therefore (ii) prevails.

We can assume Ω_∞ is a half-space, which we take to be \mathbb{R}_+^N . Notice that $P_\circ \in \partial\Omega$. Hence,

$$(17) \quad \beta^{2-N} \frac{S^{N/2}}{2} = \int_{\mathbb{R}_+^N} |\nabla \omega|^2 dx \leq \lim_{k \rightarrow \infty} \int_{\Omega_k} |\nabla v_k|^2 dx \leq \frac{S}{2^{2/N} Q_m^{(N-2)/N}},$$

which implies $Q_m \leq Q(P_\circ)$ and necessarily $Q(P_\circ) = Q_m$. Following the argument of the proof of Lemma 2.2 in [3] (see also [13], [6]), we check that $P_k \in \partial\Omega$ for large k .

Equalities (15) and (16) now easily follow. ■

Let $E_\lambda(u) = \int_\Omega |\nabla u|^2 dx + \lambda \int_{\partial\Omega} u^2 dx$ for $u \in H^1(\Omega)$ and $w_k = Q_m^{1/(p+1)} u_k$. Notice that

$$(18) \quad E_{\lambda_k}(w_k) = Q_m^{(N-2)/N} E_{\lambda_k}(u_k) < \frac{S}{2^{2/N}}$$

and by Lemma 3.3,

$$\lim_{k \rightarrow \infty} \left\| \nabla w_k - \nabla \sigma_k^{-(N-2)/2} S^{-(N-2)/4} 2^{1/(p+1)} U\left(\frac{\cdot - P_k}{\sigma_k}\right) \right\|_2 = 0,$$

where $\sigma_k = \varepsilon_k 2^{1/N} / (S^{1/2} Q_m^{1/N})$.

The sequence w_k satisfies the assumptions of Lemma 3.4 in [6]. With the aid of Lemmas 3.5–3.8 of [6] we deduce the estimate

$$E_{\lambda_k}(w_k) \geq \frac{S}{2^{2/N}} - A_N \left(\frac{N-2}{2} H(y_k) - \lambda_k \right) \varepsilon_k + O(\varepsilon_k^2) + O(\lambda_k \varepsilon_k^2),$$

where $A_N > 0$ is a constant depending on N . Therefore there exists a $\bar{\lambda} > 0$ such that $E_{\lambda_k}(w_k) > S/2^{2/N}$ for $\lambda_k \geq \bar{\lambda}$, which contradicts (18) and our assumption that $S_\lambda < S/(2^{2/N} Q_m^{(N-2)/N})$ for each $\lambda > 0$. The proof of inequality (I) is now complete.

4. Proof of inequalities (II) and (III). As with inequality (I), the proofs of (II) and (III) are by contradiction. We start with (III). The proof is

a generalization of [26]. The main difference is that now concentration occurs in the interior of Ω and the range of values of τ is $2 \leq \tau \leq 2N/(N - 1)$.

Suppose $Q_M > 2^{2/(N-2)}Q_m$. Assume that for each $\lambda > 0$, $S_{\lambda,\tau} < S_\infty := S/Q_M^{(N-2)/N}$. Let $\lambda_k \rightarrow \infty$. By Lemma 2.4, for each k there is a minimizer $u_k = u_{\lambda_k}$ of $J_{\lambda_k,\tau}$ with $\int_\Omega Q u_k^{p+1} dx = 1$. The function u_k satisfies

$$(1_{\lambda_k,\tau}) \quad \begin{cases} -\Delta u_k + \lambda_k (\int_\Omega |u_k|^\tau dx)^{2/\tau-1} u_k^{\tau-1} = S_{\lambda_k,\tau} Q u_k^p & \text{in } \Omega, \\ \partial u_k / \partial \nu = 0 & \text{on } \partial\Omega. \end{cases}$$

First of all we observe the following result:

LEMMA 4.1. *We have*

$$(19) \quad \lim_{k \rightarrow \infty} \lambda_k \left(\int_\Omega u_k^\tau dx \right)^{2/\tau} = 0.$$

Proof. Since $J_{\lambda_k,\tau}(u_k) < S_\infty$, $\lambda_k \|u_k\|_\tau^2$ is bounded and $u_k \rightarrow 0$ in $H^1(\Omega)$. By inequality (14), for a $\delta > 0$, there exists $\lambda_3 > 0$ such that

$$\begin{aligned} 1 &= \left(\int_\Omega Q(x) |u_k|^{p+1} dx \right)^{2/(p+1)} \\ &\leq \left(\frac{Q_M^{(N-2)/N}}{S} + \delta \right) \int_\Omega |\nabla u_k|^2 dx + \lambda_3 \left(\int_\Omega |u_k|^\tau dx \right)^{2/\tau}. \end{aligned}$$

Therefore

$$1 \leq \left(\frac{Q_M^{(N-2)/N}}{S} + \delta \right) \lim_{k \rightarrow \infty} \int_\Omega |\nabla u_k|^2 dx$$

and since δ is arbitrary,

$$\frac{S}{Q_M^{(N-2)/N}} \leq \lim_{k \rightarrow \infty} \|\nabla u_k\|_2^2 \leq \lim_{k \rightarrow \infty} J_{\lambda_k,\tau}(u_k) = \frac{S}{Q_M^{(N-2)/N}}.$$

Hence $\lim_{k \rightarrow \infty} \|\nabla u_k\|_2^2 = S/Q_M^{(N-2)/N}$ and $\lim_{k \rightarrow \infty} \lambda_k (\int_\Omega u_k^\tau dx)^{2/\tau} = 0$. ■

Set $M_k = \max_{\bar{\Omega}} u_k = u_k(P_k)$ for some $P_k \in \bar{\Omega}$ and

$$v_k(x) = \varepsilon_k^{(N-2)/2} u_k(\varepsilon_k x + P_k)$$

for $x \in \Omega_k = (\Omega - P_k)/\varepsilon_k$, with

$$\varepsilon_k = M_k^{-2/(N-2)}.$$

Note that the functions v_k satisfy $0 \leq v_k(x) \leq 1$ and $v_k(0) = 1$. Define

$$(20) \quad \sigma = 2N/\tau - (N - 2).$$

Since $2 \leq \tau \leq 2N/(N - 1)$, the value σ satisfies $1 \leq \sigma \leq 2$. Lemma 4.1 implies

LEMMA 4.2. $\lim_{k \rightarrow \infty} \lambda_k \varepsilon_k^\sigma = 0$.

Proof. Changing variables, we check that

$$(21) \quad \left(\int_{\Omega} u_k^\tau dx \right)^{2/\tau} = \varepsilon_k^\sigma \left(\int_{\Omega_k} v_k^\tau dx \right)^{2/\tau}.$$

But

$$(22) \quad \int_{\Omega_k} v_k^\tau dx \geq \int_{\Omega_k} v_k^{p+1} dx = \int_{\Omega} u_k^{p+1} dx \geq \frac{1}{Q_M} \int_{\Omega} Qu_k^{p+1} dx = \frac{1}{Q_M} > 0.$$

Combining (19), (21) and (22) we conclude that $\lim_{k \rightarrow \infty} \lambda_k \varepsilon_k^\sigma = 0$. ■

In particular, $\varepsilon_k \rightarrow 0$ and $M_k \rightarrow \infty$. Next we verify that the sequence $\{u_k\}$ is close to a sequence of instantons concentrating in the interior of Ω .

LEMMA 4.3. *If $Q_M > 2^{2/(N-2)}Q_m$ then, up to a subsequence, $P_k \rightarrow P_0$, where $Q(P_0) = Q_M$. Moreover,*

$$(23) \quad \lim_{k \rightarrow \infty} \int_{\Omega} \left| \nabla \left[u_k - \varepsilon_k^{-(N-2)/2} U \left(S^{1/2} Q_M^{1/N} \frac{-P_k}{\varepsilon_k} \right) \right] \right|^2 dx = 0,$$

$$(24) \quad \lim_{k \rightarrow \infty} \int_{\Omega} \left| u_k - \varepsilon_k^{-(N-2)/2} U \left(S^{1/2} Q_M^{1/N} \frac{-P_k}{\varepsilon_k} \right) \right|^{p+1} dx = 0.$$

Proof. The functions v_k are solutions of the Neumann problems

$$\begin{cases} -\Delta v_k + \lambda_k \varepsilon_k^\sigma \left(\int_{\Omega_k} |v_k|^\tau dx \right)^{2/\tau-1} v_k^{\tau-1} = S_{\lambda_k, \tau} Q(\varepsilon_k x + P_k) v_k^p & \text{in } \Omega_k, \\ \partial v_k / \partial \nu = 0 & \text{on } \partial \Omega_k, \\ 0 \leq v_k(x) \leq 1 & \text{in } \Omega_k, \\ v_k(0) = 1. \end{cases}$$

We can assume that $P_k \rightarrow P_\circ$, for some $P_\circ \in \bar{\Omega}$, and $\text{dist}(P_k, \partial \Omega) / \varepsilon_k$ converges in the extended real line. Using elliptic regularity theory we show that, up to a subsequence, $v_k \rightarrow \omega$ in $C_{\text{loc}}^2(\Omega_\infty)$, where $\Omega_\infty = \lim_{k \rightarrow \infty} \Omega_k$. The function ω satisfies

$$\begin{cases} -\Delta \omega = \widehat{S} Q(P_\circ) \omega^p & \text{in } \Omega_\infty, \\ \partial \omega / \partial \nu = 0 & \text{on } \partial \Omega_\infty, \\ 0 \leq \omega \leq 1 & \text{in } \Omega_\infty, \\ \omega(0) = 1, \end{cases}$$

where $\widehat{S} = S / Q_M^{(N-2)/N}$.

We distinguish two cases: (i) $\text{dist}(P_k, \partial \Omega) / \varepsilon_k$ converges to a real number, and (ii) $\text{dist}(P_k, \partial \Omega) / \varepsilon_k$ converges to ∞ . In case (i) we assume that $\Omega_\infty = \mathbb{R}_+^N$. By [12] we see that $\omega(x) = U(\beta x)$, where $\beta^2 = Q(P_\circ) \widehat{S}$. This yields

$$\beta^{2-N} \frac{S^{N/2}}{2} = \int_{\mathbb{R}_+^N} |\nabla \omega|^2 dx \leq \lim_{k \rightarrow \infty} \int_{\Omega_k} |\nabla v_k|^2 dx = \frac{S}{Q_M^{(N-2)/N}}.$$

Thus $Q_M \leq 2^{2/(N-2)}Q(P_o) \leq 2^{2/(N-2)}Q_m$. So case (ii) prevails. Therefore $\Omega_\infty = \mathbb{R}^N$,

$$\beta^{2-N}S^{N/2} = \int_{\mathbb{R}^N} |\nabla\omega|^2 dx \leq \lim_{k \rightarrow \infty} \int_{\Omega_k} |\nabla v_k|^2 dx = \frac{S}{Q_M^{(N-2)/N}}$$

and $Q_M \leq Q(P_o)$. Hence $Q(P_o) = Q_M$. Equalities (23) and (24) follow. ■

We now set $W(\cdot) = U(S^{1/2}Q_M^{1/N}\cdot)$ and

$$W_{\varepsilon,y}(\cdot) = \varepsilon^{-(N-2)/2}W\left(\frac{\cdot - y}{\varepsilon}\right)$$

for $y \in \mathbb{R}^N$ and $\varepsilon > 0$. Let

$$\mathcal{M} = \{CW_{\varepsilon,y} : C \in \mathbb{R}, \varepsilon > 0, y \in \bar{\Omega}\}$$

We use the notation $d(\phi, \mathcal{M}) = \text{dist}(\phi, \mathcal{M}) = \inf\{\|\nabla(\phi - \psi)\|_2 : \psi \in \mathcal{M}\}$. The following lemma, together with the last one, guarantees the existence of an instanton closest to u_k , in the metric just defined.

LEMMA 4.4. *Let $\delta > 0$ and $\{\phi_l\} \subset H^1(\Omega)$ be such that $\phi_l \rightarrow 0$ in $H^1(\Omega)$ and*

$$d(\phi_l, \mathcal{M})^2 \leq \|\nabla\phi_l\|_2^2 - 2\delta.$$

Then there exists l_o such that for all $l \geq l_o$, $d(\phi_l, \mathcal{M})$ is achieved by some $C_lW_{\varepsilon_l,y_l}$. If $y_l \rightarrow y$, with y an interior point of Ω , and w_l is defined by

$$\phi_l = C_lW_{\varepsilon_l,y_l} + w_l,$$

then, up to a subsequence,

- (i) $\lim_{l \rightarrow \infty} \varepsilon_l = 0$;
- (ii) *if $d(\phi_l, \mathcal{M}) \rightarrow 0$ as $l \rightarrow \infty$, then $\lim_{l \rightarrow \infty} C_l = C_o \neq 0$;*
- (iii) $\int_{\Omega} W_{\varepsilon_l,y_l}^p w_l dx = O(\varepsilon_l^{(N-2)/2} \|w_l\|_{H^1(\Omega)})$;
- (iv) $\int_{\Omega} W_{\varepsilon_l,y_l}^{p-1} w_l \frac{\partial}{\partial x_i} W_{\varepsilon_l,y_l} dx = O(\varepsilon_l^{(N-2)/2} \|w_l\|_{H^1(\Omega)})$.

The proof is almost identical to that of Lemma 5.6 in [21] and is omitted (see also Lemma 3.1 in [3]).

It then follows from Lemma 4.4 that there exist sequences $\{C_kW_{\delta_k,y_k}\}$ which minimize $d(u_k, \mathcal{M})$. Equality (23) implies that

$$\lim_{k \rightarrow \infty} \|\nabla(C_kW_{\delta_k,y_k}) - \nabla W_{\varepsilon_k,P_k}\|_2 = 0.$$

We deduce

$$(25) \quad C_k \rightarrow 1, \quad y_k \rightarrow y = P_o \quad \text{and} \quad \frac{\varepsilon_k}{\delta_k} \rightarrow 1.$$

In particular, y is an interior point of Ω . We set

$$(26) \quad u_k = C_k W_{\delta_k, y_k} + w_k,$$

and define

$$W_k := W_{\delta_k, y_k}.$$

Obviously, from (23) and the definition of $C_k W_k$, we get $\|\nabla w_k\|_2 \rightarrow 0$. From (25), we get $\|W_{\varepsilon_k, P_k} - C_k W_k\|_{p+1} \rightarrow 0$, which together with (24) implies $\|w_k\|_{p+1} \rightarrow 0$.

LEMMA 4.5. *There exists $0 < \mu < 1$ such that, for sufficiently large k ,*

$$pSQ_M^{2/N} \int_{\Omega} QW_k^{p-1} w_k^2 dx \leq \mu(\|\nabla w_k\|_2^2 + \lambda_k \|w_k\|_{\tau}^2).$$

This follows from Lemma 5.9 of [21].

The next lemma is essentially due to Wang [20] and Zhu [26]. We give its proof since they do not present it for the whole range of values of τ we are interested in.

LEMMA 4.6. *Let $N \geq 5$ and $2 \leq \tau \leq 2N/(N - 1)$, or $N = 4$ and $2 < \tau \leq 2N/(N - 1)$. For any $\gamma > 0$,*

$$\int_{\Omega} W_k^{\tau-1} |w_k| dx \leq o(1)\delta_k^{\tau\sigma/2} + \gamma \|w_k\|_{\tau}^{\tau}.$$

Proof. Choose

$$r \in \left] \max \left\{ \frac{N}{N-2} \frac{1}{\tau-1}, \frac{2N}{N+2} \right\}, \frac{\tau}{\tau-1} \right[.$$

Note that this interval is nonempty since on the one hand

$$\min \frac{\tau}{\tau-1} = \frac{2N}{N+1} > \frac{2N}{N+2},$$

and on the other hand

$$\frac{N}{N-2} < 2 \leq \tau \quad \text{for } N \geq 5 \quad \text{and} \quad \tau > 2 \quad \text{for } N = 4.$$

The conjugate exponent of r satisfies $r' \in]\tau, 2N/(N - 2)[$. By the Hölder inequality,

$$\begin{aligned} \int_{\Omega} W_k^{\tau-1} |w_k| dx &\leq \left(\int_{\Omega} W_k^{(\tau-1)r} dx \right)^{1/r} \left(\int_{\Omega} |w_k|^{r'} dx \right)^{1/r'} \\ &\leq C\delta_k^{-(\tau-1)(N-2)/2+N/r} \|w_k\|_{r'}, \end{aligned}$$

and, by the interpolation inequality,

$$\|w_k\|_{r'} \leq \|w_k\|_{\tau}^a \|w_k\|_{p+1}^{1-a},$$

where

$$\frac{1}{r'} = \frac{a}{\tau} + \frac{1-a}{p+1}, \quad 0 < a < 1.$$

If $\gamma > 0$, there exists a $C(\gamma)$ such that

$$\begin{aligned} (27) \quad & \int_{\Omega} W_k^{\tau-1} |w_k| dx \\ & \leq C \delta_k^{-(\tau-1)(N-2)/2+N/r} \|w_k\|_{\tau}^a \|w_k\|_{p+1}^{1-a} \\ & \leq \gamma \|w_k\|_{\tau}^{\tau} + C(\gamma) \delta_k^{(-(\tau-1)(N-2)/2+N/r) \cdot \tau / (\tau-a)} \|w_k\|_{p+1}^{\tau(1-a)/(\tau-a)}. \end{aligned}$$

Now, from the definition of a ,

$$a \left(\frac{1}{\tau} - \frac{1}{p+1} \right) = \frac{1}{r'} - \frac{1}{p+1}$$

and from the definition of σ (eq. (20)),

$$a \left(\frac{1}{\tau} - \frac{1}{p+1} \right) = \frac{a\sigma}{2N}.$$

Hence

$$\frac{a\sigma}{2N} = \frac{1}{r'} - \frac{1}{p+1}.$$

From this we successively get

$$\begin{aligned} \frac{1}{r} - \frac{1}{N} &= \frac{1}{2} - \frac{a\sigma}{2N}, \\ \frac{1}{r} - \frac{(N-2)(\tau-1)}{2N} &= \left(\frac{\tau}{2N} - \frac{a}{2N} \right) \sigma, \\ \left(-\frac{(\tau-1)(N-2)}{2} + \frac{N}{r} \right) \cdot \frac{\tau}{\tau-a} &= \frac{\tau\sigma}{2}. \end{aligned}$$

Substituting into (27), we derive

$$\int_{\Omega} W_k^{\tau-1} |w_k| dx \leq \gamma \|w_k\|_{\tau}^{\tau} + C(\gamma) \delta_k^{\tau\sigma/2} \|w_k\|_{p+1}^{\tau(1-a)/(\tau-a)}.$$

Since (24) implies that $\|w_k\|_{p+1}^{\tau(1-a)/(\tau-a)} \rightarrow 0$ as $k \rightarrow \infty$, the proof is complete. ■

We will now prove that for large k , $J_{\lambda_k, \tau}(u_k) > S/Q_M^{(N-2)/N}$, which contradicts Lemma 2.8 and proves inequality (III). So, we estimate the terms in $J_{\lambda_k, \tau}$. By construction, $\int_{\Omega} Q|u_k|^{p+1} dx = 1$. However, it is convenient to give the following lower bound whose proof follows from (iii) of Lemma 4.4

and the arguments in [26]:

$$\begin{aligned} \left(\int_{\Omega} Q|u_k|^{p+1} dx \right)^{-2/(p+1)} &\geq Q_M^{-(N-2)/N} C_k^{-2} \|W_k\|_{p+1}^{-2} \\ &\quad \times \left(1 - \frac{(p + \gamma_1) \int_{\Omega} QW_k^{p-1} w_k^2 dx}{C_k^2 \|W_k\|_{p+1}^{p+1}} \right. \\ &\quad \left. - C\delta_k^{(N-2)/2} \|w_k\|_{H^1(\Omega)} - C(\gamma_1) \|w_k\|_{H^1(\Omega)}^{p+1} \right) \end{aligned}$$

where γ_1 is any positive number. Regarding $\|W_k\|_{p+1}^{-2}$, we have

$$\|W_k\|_{p+1}^{-2} \geq Q_M^{(N-2)/N} + O(\delta_k^N).$$

Inserting this estimate in the last inequality, we get the following lower bound for $(\int_{\Omega} Q|u_k|^{p+1} dx)^{-2/(p+1)}$:

$$\begin{aligned} C_k^{-2} \left(1 - \frac{(p + \gamma_1) \int_{\Omega} QW_k^{p-1} w_k^2 dx}{C_k^2 \|W_k\|_{p+1}^{p+1}} \right. \\ \left. - C\delta_k^{(N-2)/2} \|w_k\|_{H^1(\Omega)} - C(\gamma_1) \|w_k\|_{H^1(\Omega)}^{p+1} + O(\delta_k^N) \right). \end{aligned}$$

To estimate $(\int_{\Omega} u_k^\tau dx)^{2/\tau}$, note first that

$$\left(\int_{\Omega} W_k^\tau dx \right)^{2/\tau} = \delta_k^\sigma \left(\int_{\Omega_k} W^\tau dx \right)^{2/\tau} = C\delta_k^\sigma + o(\delta_k^\sigma)$$

since $(N - 2)\tau > N$. In fact, $(N - 2)/N < 2 \leq \tau$ except when $N = 4$, but in this case $\tau > 2$. So we can follow the argument in [26] to prove that for any $0 < \gamma_2 < 1$ there is a constant $C(\gamma_2)$ such that

$$\left(\int_{\Omega} u_k^\tau dx \right)^{2/\tau} \geq \gamma_2 \|w_k\|_\tau^2 + C(\gamma_2) \delta_k^\sigma.$$

To estimate $\int_{\Omega} |\nabla u_k|^2 dx$, note that

$$\int_{\Omega} |\nabla W_k|^2 dx = \frac{1}{S^{(N-2)/2} Q_M^{(N-2)/N}} \int_{\Omega} |\nabla U_k|^2 dx = \frac{S}{Q_M^{(N-2)/N}} + O(\delta_k^{N-2}).$$

Therefore

$$\int_{\Omega} |\nabla u_k|^2 dx = C_k^2 \frac{S}{Q_M^{(N-2)/N}} + O(\delta_k^{N-2}) + \int_{\Omega} |\nabla w_k|^2 dx.$$

We are now in a position to estimate $J_{\lambda_k, \tau}(u_k)$. Combining the previous estimates gives

$$\begin{aligned}
 J_{\lambda_k, \tau}(u_k) \geq & \left[\frac{S}{Q_M^{(N-2)/N}} + O(\delta_k^{N-2}) + \frac{\|\nabla w_k\|_2^2}{C_k^2} + \frac{\gamma_2 \lambda_k \|w_k\|_\tau^2}{C_k^2} + \frac{C(\gamma_2) \lambda_k \delta_k^\sigma}{C_k^2} \right] \\
 & \times \left[1 - \frac{(p + \gamma_1) \int_\Omega Q W_k^{p-1} w_k^2 dx}{C_k^2 \|W_k\|_{p+1}^{p+1}} \right. \\
 & \left. - C \delta_k^{(N-2)/2} \|w_k\|_{H^1(\Omega)} - C(\gamma_1) \|w_k\|_{H^1(\Omega)}^{p+1} + O(\delta_k^N) \right].
 \end{aligned}$$

For k large, this is greater than or equal to

$$\begin{aligned}
 & \frac{S}{Q_M^{(N-2)/N}} + \frac{C(\gamma_2) \lambda_k \delta_k^\sigma}{C_k^2} + O(\delta_k^{N-2}) \\
 & + \gamma_2 \frac{\|\nabla w_k\|_2^2 + \lambda_k \|w_k\|_\tau^2}{C_k^2} - \frac{(p + \gamma_1) S \int_\Omega Q W_k^{p-1} w_k^2 dx}{C_k^2 Q_M^{(N-2)/N} \|W_k\|_{p+1}^{p+1}}.
 \end{aligned}$$

By Lemma 4.5, for k sufficiently large, the difference

$$\gamma_2 \frac{\|\nabla w_k\|_2^2 + \lambda_k \|w_k\|_\tau^2}{C_k^2} - \frac{(p + \gamma_1) S \int_\Omega Q W_k^{p-1} w_k^2 dx}{C_k^2 Q_M^{(N-2)/N} \|W_k\|_{p+1}^{p+1}}$$

is greater than or equal to

$$\frac{p Q_M^{2/N}}{C_k^2} \left(\frac{\gamma_2}{\mu} - \frac{p + \gamma_1}{p} \right) \int_\Omega Q W_k^{p-1} w_k^2 dx,$$

for some $0 < \mu < 1$. Choosing γ_2 such that $\mu < \gamma_2 < 1$, and then γ_1 sufficiently small, we get

$$\frac{\gamma_2}{\mu} - \frac{p + \gamma_1}{p} > 0.$$

This implies that

$$J_{\lambda_k, \tau}(u_k) \geq \frac{S}{Q_M^{(N-2)/N}} + \frac{C(\gamma_2)}{C_k^2} \lambda_k \delta_k^\sigma + O(\delta_k^{N-2}).$$

Since $1 \leq \sigma \leq 2$, it follows that $\sigma \leq N - 2$, because if $N = 4$ we do not allow τ (and hence σ) to equal 2. So, for sufficiently large k ,

$$J_{\lambda_k, \tau}(u_k) > S/Q_M^{(N-2)/N},$$

which implies inequality (III), as explained above.

Now that we have proved (III), let us *outline the proof of inequality (II)*. Let $Q_M \leq 2^{2/(N-2)} Q_m$. One assumes that $S_{\lambda, \tau} < S/(2^{2/N} Q_m^{(N-2)/N})$ for each $\lambda > 0$, picks $\lambda_k \rightarrow \infty$ and minimizers $u_k = u_{\lambda_k}$ of $J_{\lambda_k, \tau}$ with $\int_\Omega Q u_k^{p+1} dx = 1$, and starts to prove analogues of the previous lemmas. When one gets to the analogue Lemma 4.3, one proves that, up to a subsequence, $\{u_k\}$ either concentrates in the interior of Ω or concentrates on

the boundary of Ω . In the former case concentration occurs at a point P_\circ with $Q(P_\circ) = Q_M$ and one can repeat the argument given above to derive a contradiction. In the latter case concentration occurs at a point P_\circ with $Q(P_\circ) = Q_m$.

Let $E_{\lambda,\tau}(u) = \int_\Omega |\nabla u|^2 + \lambda(\int_\Omega u^\tau)^{2/\tau}$ for $u \in H^1(\Omega)$ and $w_k = Q_m^{1/(p+1)} u_k$. Notice that

$$(28) \quad E_{\lambda_k,\tau}(w_k) = Q_m^{(N-2)/N} E_{\lambda_k}(u_k) < \frac{S}{2^{2/N}}.$$

Following the arguments in [26], one can prove that $E_{\lambda_k,\tau}(w_k) > S/2^{2/N}$ for λ_k sufficiently large, which contradicts (28) and proves inequality (II).

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