## Multiple slowly oscillating periodic solutions for x'(t) = f(x(t-1)) with negative feedback

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Abstract. We consider the prototype equation

$$x'(t) = f(x(t-1))$$

for delayed negative feedback. We review known results on uniqueness and nonuniqueness of slowly oscillating periodic solutions, and present some new results and examples.

1. Introduction. Consider the scalar-valued differential equation (1.1) x'(t) = f(x(t-1))

with constant delay 1. Throughout, we shall impose the following assumptions on the function f:

(H)  $\begin{cases} f \text{ is continuous and bounded below or above;} \\ f \text{ is of negative feedback type: } xf(x) < 0 \text{ for all } x \neq 0; \\ f \text{ is } C^1 \text{ on a neighborhood of } 0, \text{ with } f'(0) < 0. \end{cases}$ 

Clearly, the continuity of f implies f(0) = 0; accordingly, the zero function  $x(t) = 0, t \in \mathbb{R}$ , always satisfies (1.1) on all of  $\mathbb{R}$ .

Various applications have been proposed for (1.1). For example, taking  $f(u) = -\alpha(1 - e^{-u}), \alpha > 0$ , yields what is now called *Wright's equation*. An appropriate transformation yields the equation (also called Wright's equation)  $y'(t) = -\alpha y(t-1)[1+y(t)]$ , which was examined in [43] and was heuristically motivated by the study of the distribution of primes; see Marshall and Smith [27] for a recent recounting of this connection.

Such applications aside, it is undoubtedly the simplicity of (1.1) that has made it among the most intensively studied and best understood types of differential delay equations.

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A slowly oscillating periodic solution (SOP solution) of (1.1) is a periodic function  $p : \mathbb{R} \to \mathbb{R}$  that satisfies the equation for all  $t \in \mathbb{R}$  and has the feature that if  $z \neq z'$  are distinct zeros of p, then |z - z'| > 1. SOP solutions, when they exist, often play an important role in the global dynamics of (1.1), and they have been studied for the last several decades. Our goal in this paper is to review some known results on SOP solutions of (1.1) (especially concerning uniqueness and nonuniqueness), to sketch the prominent techniques used to obtain those results, and to provide some novel examples of equations with multiple SOP solutions. In particular, we present a criterion for (1.1) to have multiple nontrivial SOP solutions when the equilibrium at 0 is locally attracting, and a family of equations for which we can obtain multiple SOP solutions, one of which has arbitrarily long period, while keeping f uniformly bounded.

In the usual way we write  $C = C([-1,0], \mathbb{R})$  for the Banach space of all continuous functions  $\varphi : [-1,0] \to \mathbb{R}$  equipped with the norm  $\|\varphi\| = \sup_{s \in [-1,0]} |\varphi(s)|$  of uniform convergence. If x is any continuous real-valued function whose domain contains the interval [t-1,t], the segment  $x_t$  of xat t is the member of C given by the formula  $x_t(s) := x(t+s), s \in [-1,0]$ . A solution of (1.1) is either a continuously differentiable function  $x : \mathbb{R} \to \mathbb{R}$ that satisfies (1.1) for all  $t \in \mathbb{R}$ , or a continuous function  $x : [-1,\infty) \to \mathbb{R}$ that is continuously differentiable for all t > 0 and satisfies (1.1) for all t > 0. In either case,  $x_0$  is called the *initial value* of the solution x. We regard Cas the phase space for (1.1), and view solutions x as describing orbits  $\{x_t\}$ in C.

By the so-called *method-of-steps*, it is easily seen that every function  $\varphi \in C$  determines a solution  $x^{\varphi} : [-1, \infty) \to \mathbb{R}$  of (1.1) with initial value  $x_0^{\varphi} = \varphi$ . Indeed, for  $t \in [0, 1]$ , we have

$$x^{\varphi}(t) = \varphi(0) + \int_{0}^{t} f(\varphi(s-1)) \, ds,$$

and the formulas for  $x^{\varphi}$  on  $[1, 2], [2, 3], \ldots$  are then similar. We call  $x^{\varphi}$  the continuation of  $\varphi$  as a solution of (1.1). Equation (1.1) defines a semiflow

$$F: [0,\infty) \times C \ni (t,\varphi) \mapsto x_t^{\varphi} \in C$$

on the Banach space C. The semiflow F is continuous in the following sense: given  $\varphi \in C$ ,  $t \geq 0$ , and  $\varepsilon > 0$ , there exists  $\delta > 0$  such that for all  $\psi \in C$ with  $\|\psi - \varphi\| < \delta$  and all  $0 \leq s \leq t$ ,

$$|F(s,\psi) - F(s,\varphi)| = |x^{\psi}(s) - x^{\varphi}(s)| < \varepsilon.$$

Additionally, using the Arzelà–Ascoli theorem it not difficult to show that all maps  $F(t, \cdot) : C \to C, t \ge 1$ , are compact in the sense that they map bounded sets to precompact sets.

We define a solution x to be *slowly oscillating* if for any two zeros z < z'of x we have z' - z > 1. Similarly, a solution x is called *eventually slowly oscillating* if there is some  $\tau \in \mathbb{R}$  such that z' - z > 1 for any two zeros  $\tau < z < z'$  of x. Suppose that  $\varphi \in C$  is such that (for example)  $\varphi(s) > 0$  for all  $s \in [-1, 0]$ . Then either  $x^{\varphi}(t) > 0$  for all t > 0 or not; in the latter case, there is some first positive zero z of  $x^{\varphi}$ . The negative feedback condition then implies that f is strictly decreasing on [z, z + 1]. Thus the next zero z'after z, if it exists at all, satisfies z' - z > 1; repeating this argument shows that any successive zeros of  $x^{\varphi}$  must be separated by more than one unit. Thus eventually slowly oscillating solutions clearly exist. Indeed, much more is true: if f is smooth and strictly decreasing, the solution  $x^{\varphi}$  is eventually slowly oscillating for a dense subset of initial values  $\varphi \in C$  (see Mallet-Paret and Walther [26], and the earlier but somewhat more restricted work of Walther [37]); a similar conclusion seems plausible for (1.1) in general. [We point out, however, that examples are known of equations of the form

$$x'(t) = -\mu x(t) + f(x(t-1)), \quad \mu \ge 0$$

(that is, like (1.1) but with an instantaneous damping term added), for which open sets of initial conditions have continuations that are not eventually slowly oscillating. See, for example, Ivanov and Losson [15] and Stoffer [35]. For such examples, f must be nonmonotone: again see [26].]

The zero function  $0 \in C$  is a stationary point of the semiflow F: F(t, 0) = 0for all  $t \geq 0$ . The linear problem associated with (1.1) at 0 is the equation

(1.2) 
$$y'(t) = f'(0)y(t-1).$$

Using the ansatz  $y(t)=e^{\lambda t}$  we obtain the corresponding characteristic equation

(1.3) 
$$\lambda - f'(0)e^{-\lambda} = 0.$$

The roots of (1.3) coincide with the eigenvalues of the linearization of F at  $0 \in C$ . It is also known that, when f'(0) < 0, there are only finitely many eigenvalues with nonnegative real part; there are eigenvalues with positive real part if and only if  $f'(0) < -\pi/2$ . Accordingly, the zero solution x(t) = 0,  $t \in \mathbb{R}$ , of (1.1) is locally asymptotically stable when  $0 > f'(0) > -\pi/2$  and unstable when  $f'(0) < -\pi/2$  by the so-called principle of linearized stability and instability, respectively. In the former case, the trivial solution of (1.2) is globally attracting and the trivial stationary point of (1.1) is at least locally attracting. As we are interested in slowly oscillating periodic solutions we of course will only consider examples in which the trivial solution of (1.1) is at most locally, but not globally, attracting.

REMARK 1.1. Many authors have given conditions under which the trivial solution of (1.1) is globally attracting.

A long-standing conjecture, going back to Wright [43], is that 0 is globally attracting for Wright's equation when  $\alpha \in (0, \pi/2)$ ; this was proven for  $\alpha \in (0, 3/2]$  in [43]. Further progress toward this conjecture is described by Bánhelyi et al. [2].

In an extension of the Wright result, Liz et al. [23] showed that if f is sufficiently smooth, has negative Schwarzian derivative everywhere, and satisfies  $f'(0) \in [-3/2, 0)$ , then the zero solution  $x(t) = 0, t \in \mathbb{R}$ , is globally attracting.

Other results on the global attractivity of the trivial solution of (1.1) go back to the "3/2-stability" criterion given by Yorke [47]. See, for example, Liz et al. [24].

2. SOP solutions—essential methods and results. For general facts about delay differential equations we recommend the monograph [8] by Diekmann et al. and the recent review article [42] by Walther.

The map  $P: K \to K$ . For the remainder of this paper we restrict our attention to the case that f is bounded above; the case where f is bounded below is essentially the same (and can be reduced to our case by considering the equation y'(t) = -f(-y(t-1))). In particular, let us take M > 0 such that  $f(x) \leq M$  for all  $x \in \mathbb{R}$ .

Let us write

$$K := \{ \varphi \in C \mid \varphi(-1) = 0, \varphi \text{ is nondecreasing}, \|\varphi\| \le M \}.$$

Then K is a closed convex subset of C.

The following theorem is elementary and mostly intuitively clear. For a proof we refer the reader to Diekmann et al. [8, Chapter XV], or to Walther [39] where the more general equation  $x'(t) = -\mu x(t) + f(x(t-1))$  with  $\mu \ge 0$  is discussed.

PROPOSITION 2.1. Let  $\varphi \in K$ ,  $\varphi \neq 0$ . Then  $x = x^{\varphi}$  is slowly oscillating on  $[0, \infty)$ . If z > 0 is a zero of x then |x| and |x'| are bounded on the interval [z, z + 1] by

$$\max\{|f(\xi)| \mid \xi \in x([z-1,z])\},\$$

and the function  $s \mapsto |x_{z+1}(s)|$  is increasing on [-1, 0].

If the zeroset of  $x|_{[0,\infty)}$  is unbounded then it is given by a sequence of points  $z_j = z_j(\varphi), j \in \mathbb{N}$ , with

(so) 
$$z_j + 1 < z_{j+1}$$
 and  $x'(z_j) \neq 0$  for all  $j$ ,

and x is monotone on  $[0, z_1 + 1]$  and on each interval  $[z_j + 1, z_{j+1} + 1]$ .

If the zeroset of  $x|_{[0,\infty)}$  is bounded then it is given by a finite sequence of points  $z_j = z_j(\varphi)$ ,  $1 \leq j \leq J = J(\varphi)$ , with property (so), and |x(t)|decreases monotonically to 0 on the interval  $[z_J + 1, \infty)$  as  $t \to \infty$ . Given an initial value  $\varphi \in K$ ,  $\varphi \neq 0$ , there are two possibilities: either  $x^{\varphi}$  has a second positive zero z and  $x_{z+1}^{\varphi} \in K$ , or not. We define the return map  $P: K \to K$  by setting P(0) = 0,  $P(\varphi) := x_{z+1}^{\varphi}$  if z > 0 exists, and  $P(\varphi) := 0$  otherwise. Here are the basic facts about  $x^{\varphi}$  and the map P (for a proof, see Nussbaum [28, Lemma 2.7]).

PROPOSITION 2.2. The map P is continuous, and P(K) has compact closure.

Suppose that  $\varphi$  is a nonzero periodic point of P with minimal period k. Then, writing  $0 < z_1 < \cdots < z_{2k} < \cdots$  for the positive zeros of  $x^{\varphi}$ , we find that  $x_{z_{2k+1}}^{\varphi} = \varphi$  and that x is a periodic solution of (1.1) with minimal period  $z_{2k} + 1$ . Conversely, any slowly oscillating periodic solution of (1.1) has a segment that is a periodic point of P. Accordingly, perhaps the most prominent technique for proving the existence of SOP solutions of nonlinear equations like (1.1) and its generalizations has been to show that analogs of the return map  $P: K \to K$  have nonzero fixed points. Since  $0 \in K$  is a fixed point of P, merely establishing the existence of at least one fixed point (by a naive application of Schauder's Theorem, for example) is insufficient.

We do, however, have the following theorem. This is perhaps the bestknown theorem on existence of SOP solutions of (1.1) and was given by Nussbaum in 1974 [29]. The theorem generalizes an earlier result due to Jones on Wright's equation [16], which in turn drew on the ideas of Wright [43] and Browder [3].

THEOREM 2.3. If, in addition to the hypotheses (H), we assume that  $f'(0) < -\pi/2$  (and so the equilibrium at the origin is unstable), then P has a nonzero fixed point.

The idea of Nussbaum's proof is to show that the instability of the origin makes  $0 \in K$  a so-called *ejective fixed point* of P and then to apply the so-called Browder ejective fixed point principle [4], which implies that Pmust have a fixed point that is not ejective. This same basic idea has since been used to prove the existence, when the equilibrium is unstable, of nontrivial periodic solutions for several generalizations and extensions of (1.1); these generalizations include versions of (1.1) with instantaneous damping (Hadeler and Tomiuk [13]) and various equations with state-dependent delay (e.g. Nussbaum [29], Alt [1], Kuang and Smith [20], and Walther [41]).

REMARK 2.4. The above-cited [16] and the companion paper [17] of Jones have led to the conjecture that, for  $\alpha > \pi/2$ , Wright's equation has a *unique* SOP solution. For some recent work towards this conjecture, see Xie [46] and Lessard [22].

We now turn to some other important methods for studying SOP solutions of (1.1).

When f is odd—phase plane methods. An entirely different approach to proving existence of SOP solutions in the case that f in (1.1) is odd was pioneered by Kaplan and Yorke in 1974 [18]. In this case the coupled two-dimensional system of ordinary differential equations

(2.1) 
$$\begin{cases} u' = f(v), \\ v' = -f(u) \end{cases}$$

has a first integral of the form

(2.2) 
$$H(u,v) := -\int_{0}^{u} f(s) \, ds - \int_{0}^{v} f(s) \, ds.$$

Equation (2.1) is highly symmetric: if  $(u(\cdot), v(\cdot)) : I \to \mathbb{R}^2$ ,  $I \subset \mathbb{R}$ an interval, is a solution of (2.1), then so are  $(-u(\cdot), -v(\cdot)) : I \to \mathbb{R}^2$ ,  $(v(\cdot), -u(\cdot)) : I \to \mathbb{R}^2$ , and  $(-v(\cdot), u(\cdot)) : I \to \mathbb{R}^2$ , as simple calculations show. Note that H(0,0) = 0 whereas H(u,v) > 0 for all  $(u,v) \neq (0,0)$ , and  $\nabla H(u,v) = (0,0) \in \mathbb{R}^2$  if and only if (u,v) = (0,0); it follows that, given any real  $\alpha > 0$  with  $\alpha \in H(\mathbb{R}^2) := \{H(u,v) \in [0,\infty) \mid u,v \in \mathbb{R}\}$ , the level set  $H^{-1}(\{\alpha\})$  is a simple closed curve around the origin in  $\mathbb{R}^2$ . This curve is the orbit of a periodic and nonconstant solution  $(u(\cdot), v(\cdot)) : \mathbb{R} \to \mathbb{R}^2$ of (2.1). From the symmetry of H we conclude that in this situation the orbits of  $(u(\cdot), v(\cdot)), (-u(\cdot), -v(\cdot)), (v(\cdot), -u(\cdot))$ , and  $(-v(\cdot), u(\cdot))$  coincide; that is, the orbit  $H^{-1}(\{\alpha\})$  is invariant under rotations by  $\pi/2$ , and all the solutions  $(-u(\cdot), -v(\cdot)), (v(\cdot), -u(\cdot))$ , and  $(-v(\cdot), u(\cdot))$  are translations of the periodic solution  $(u(\cdot), v(\cdot))$ .

It was observed by Kaplan and Yorke in [18] that any solution  $(u(\cdot), v(\cdot))$ :  $\mathbb{R} \to \mathbb{R}^2$  of (2.1) with period 4 is associated to a periodic solution of (1.1): namely,  $x(t) := u(t), t \in \mathbb{R}$ , satisfies (1.1), is periodic of minimal period 4, and has the special symmetry  $x(t) = -x(t-2), t \in \mathbb{R}$ . We shall refer to such a solution as a Kaplan-Yorke solution of (1.1).

The existence result obtained in [18] for (1.1) is the following. The hypotheses guarantee that (2.1) has a solution of period 4.

THEOREM 2.5. Assume that, in addition to hypothesis (H), f is odd and  $\int_0^\infty |f(x)| dx = \infty$ . Write

$$a = \lim_{x \to 0} \frac{f(x)}{x}$$
 and  $A = \lim_{x \to \infty} \frac{f(x)}{x}$ .

If either  $A < -\pi/2 < a$  or  $a < -\pi/2 < A$ , then (1.1) has a Kaplan–Yorke solution.

Nussbaum [31, Remark 2.1] uses the above theorem and the ideas of its proof to indicate a basic method to construct inductively a continuous odd function f ensuring that (2.1) has at least n distinct periodic solutions of

period 4, and so (1.1) has at least n distinct Kaplan–Yorke solutions, for any fixed  $n \in \mathbb{N}$ . On the other hand, in the same work Nussbaum combines the last result with Schauder's theorem—roughly speaking applied on a subset of a variant of K—to exhibit an example where Kaplan–Yorke solutions coexist with solutions of period greater than 4. We will present some other examples of a similar flavor below.

When f is strictly decreasing. As already suggested, the case that f is smooth and strictly decreasing is particularly well understood. A first major step in this understanding was made by Kaplan and Yorke [19], who considered the orbits traced out by slowly oscillating solutions in the (x(t), x(t-1)) plane. The chief observation was that the ways in which these orbits can cross in the plane is sharply limited. The main theorem in [19] says, roughly speaking, that if  $f'(0) < -\pi/2$  these orbits tend toward an annulus in the plane whose inner and outer boundary curves are the orbits of SOP solutions. In particular, if (1.1) has a unique SOP solution in this case, this SOP solution is necessarily asymptotically stable.

In [31], Nussbaum built on the phase plane approach in [19] to give conditions under which this SOP solution is unique (it is sufficient, for example, that f be smooth and odd with f'(x) and f(x)/x increasing on the positive half-line). Also using a phase plane approach, Cao [6] established uniqueness of SOP solutions for f smooth and strictly decreasing and with a slightly different concavity condition, but without the assumption that fis odd.

Mallet-Paret and Sell [25] have given a Poincaré–Bendixson type theorem for a class of equations that includes (1.1) when f is smooth and strictly decreasing. In this case, the theorem says that  $\omega$ -limit sets of solutions are either the equilibrium at 0 or nontrivial periodic solutions (at least if the origin is hyperbolic); and that if p and q are two distinct nontrivial periodic solutions then the planar traces given by  $t \mapsto (p(t), p(t-1))$  and  $t \mapsto (q(t), q(t-1))$  are disjoint Jordan curves (we shall revisit this point later). Among the implications of this fact is that all SOP solutions correspond to fixed points (as opposed to higher-period periodic points) of the map  $P: K \to K$ .

Walther has shown that when f is smooth and strictly decreasing, the attractor of the set of slowly oscillating solutions is either trivial or homeomorphic to a closed unit disk [39].

The fixed point index. Other existence results, as well as uniqueness and nonuniqueness results, depend on the so-called fixed point index. We shall use the index below and give some of its flavor; very roughly speaking, the index sometimes allows us to use information on subsets of K where we understand P well to draw conclusions about the existence of fixed points on subsets of K where we understand P less well. We mention here a handful of earlier results. In [28], Nussbaum used a fixed point index approach to establish the existence of SOP solutions under hypotheses similar to those in Theorem 2.5, but without the assumption that f is odd; our work in Section 3 below bears a strong conceptual affinity with Nussbaum's approach. In [44] and [46], Xie combined stability estimates and the fixed point index to obtain, in certain cases, uniqueness of fixed points of P.

Stability of SOP solutions. Assessing the stability of SOP solutions is, in general, a difficult problem. The stability implied by uniqueness in the case that f is monotonic, alluded to above, is one result; and various authors (see, for example, Cao [5]) have used Kaplan–Yorke-type phase plane techniques with f monotonic to conclude, based on how solutions "spiral" in the (x(t), x(t-1)) plane, that certain SOP solutions are stable or unstable. Other authors have observed that, loosely speaking, if f is sufficiently flat on long enough intervals, the return map P becomes contractive (at least on certain subsets of K) and so we can conclude that asymptotically stable SOP solutions exist. See, for example, Walther [40] and Vas [36].

Several authors have also established results on the so-called Floquet multipliers of SOP solutions, especially in the case that f is odd. If p is an SOP solution with  $p_0 \in K$  and period  $\omega$ , then the Floquet multipliers of p constitute the nonzero spectrum of the linear monodromy operator  $\psi \mapsto$  $D_2F(\omega, p_0)\psi$ . With the possible exception of 1 (which is always an eigenvalue of the monodromy operator with corresponding eigenvector  $p'_0$ ), the nonzero spectrum of the monodromy operator coincides with the nonzero spectrum of the derivative of P at  $p_0$  (see, for example, Xie [45] or Diekmann et al. [8, Chapter XIV] for a fuller discussion).

Results on Floquet multipliers of SOP solutions include the following. Stability of Kaplan–Yorke solutions is proven by Chow and Walther [7] for a class of odd f satisfying a monotonicity and concavity condition; instability of Kaplan–Yorke solutions under conceptually similar kinds of conditions is proven by Ivanov et al. [14]. (The two sets of hypotheses in these two papers correspond, roughly speaking, to the two possibilities  $A < -\pi/2 < a$  and  $a < -\pi/2 < A$  in the statement of Theorem 2.5.) For additional results and references we also mention the more recent work of Skubachevskii and Walther [34].

**Bifurcation results.** Several authors have given bifurcation results for parameterized versions of (1.1). In 1975 Nussbaum [30] augmented Theorem 2.3 by showing that, for the parameterized equation

$$x'(t) = \alpha f(x(t-1)), \quad f'(0) = -1,$$

there is a continuous branch  $\{(\alpha, p_0^{\alpha}) : \alpha > \pi/2\}$  of initial conditions of nontrivial SOP solutions, with  $(\pi/2, 0)$  in its closure. Further branches of periodic solutions emanating from bifurcation points on this branch (and hence nonuniqueness of SOP solutions) have been described, especially in the case that f is odd. Dormayer has described branches of periodic solutions bifurcating from the primary branch of Kaplan–Yorke solutions and consisting of solutions that are symmetric [9] and that are not symmetric [10] (here symmetric means that  $p(t) = -p(t+\tau)$ , where  $\tau$  is the half-period). Both of these results draw on earlier work of Walther [38]. See also Dormayer and Lani-Wayda [11], and the references therein.

Other results on nonuniqueness. Apart from the nonuniqueness results on periodic solutions arising from the bifurcation scenarios mentioned above or from the above-discussed work [31] of Nussbaum, we mention some more nonuniqueness results below.

For instance, Peters [32] and Siegberg [33] have exhibited examples of (1.1) where P is semiconjugate, on a suitable subset of K, to a chaotic interval map; in this case P has many periodic points, of many different periods. A more recent result in this vein is given by Lani-Wayda in [21], which contains additional references.

Cao [5] and Vas [36] have both given further examples of equations with many SOP solutions, many of which are stable. In [5], f is monotonic, and the above-mentioned results of Kaplan and Yorke on the behavior of orbits in the (x(t), x(t-1)) plane are exploited to devise a condition that guarantees the existence of several (possibly infinitely many) SOP solutions whose planar orbits are nested within one another. In [36], f is unbounded and similar to a decreasing step function (the equation in [36] also has a damping term added). The many intervals where f is nearly flat cause the analog of P to be contractive on many different subdomains of C, resulting in many distinct periodic solutions. In Section 5 we present an example in a similar spirit, but the feedback function is bounded and the nonmonotonicity of fplays a key role.

## 3. A class of equations with at least two SOP solutions

Facts about the fixed point index. We here recall only those properties of the fixed point index that we shall need and even those in a restricted setting. For a deeper discussion consult Granas and Dugundji [12, Chapter IV] and also Nussbaum [29, Section 1].

Suppose that X is a closed, bounded, convex subset of a Banach space. Then X has the subspace topology induced by the Banach space topology. (This is the topology to which we shall be referring throughout this paragraph—so the open sets we refer to just below are open relative to X, and may have empty interior with respect to the underlying Banach space.) Suppose that  $G: X \to X$  is continuous and compact in the sense that it maps bounded sets of X into precompact sets. Then, if V is any open (relative to X) subset of X such that G is fixed-point free on  $\partial V := \overline{V} \setminus V$ , the integer  $i_X(G, V)$ , the fixed point index of G on V with respect to X, is defined. The following two properties hold:

- (I) If  $V \subset X$  is convex and  $G(\overline{V}) \subset V$ , then G has a fixed point in V and  $i_X(G, V) = 1$ . In particular,  $i_X(G, X) = 1$ .
- (II) If  $i_X(G, U)$  and  $i_X(G, V)$  are defined,  $\overline{U} \cap \overline{V} = \emptyset$ , and the fixed point set of G is contained in  $U \cup V$ , then

$$1 = i_X(G, X) = i_X(G, U) + i_X(G, V).$$

Property (I) follows from the so-called normalization and homotopy properties of the index, and (II) is called the additivity property. We will also need the so-called mod p theorem for the fixed point index. The following is a very special case of this theorem; a general statement appears, for instance, in Granas and Dugundji [12, p. 460]:

LEMMA 3.1. Suppose that X is a closed, bounded, convex subset of a Banach space and that  $G : X \to X$  is a continuous compact map. Let  $p \in \mathbb{N}$  be prime and write  $\operatorname{Fix}(G)$  and  $\operatorname{Fix}(G^p)$  for the fixed point sets of G and  $G^p$  in X, respectively. Suppose that  $V \subset X$  is relatively open and that  $\operatorname{Fix}(G) \cap \overline{V} = \operatorname{Fix}(G^p) \cap \overline{V} = \{q\}, q \in V$ . Then  $i_X(G, V)$  is congruent to  $i_X(G^p, V)$  modulo p.

As a consequence of this result we obtain the next corollary, which will be essential for our approach.

COROLLARY 3.2. Suppose that X is a closed, bounded, convex subset of a Banach space and that  $G: X \to X$  is a continuous compact map. Let U be a relatively open convex subset of X. Suppose that  $q \in U$  is a fixed point of G and that there is some  $n \in \mathbb{N}$  such that, for all integers  $m \ge n$ ,

 $G^m(\overline{U}) \subset U \quad and \quad \mathrm{Fix}(G^m) \cap \overline{U} = \{q\}.$ 

Then  $i_X(G, U) = 1$ .

Proof. Since  $\operatorname{Fix}(G) \cap \overline{U} \subset \operatorname{Fix}(G^m) \cap \overline{U}$  for all  $m \geq n$  and  $q \in \operatorname{Fix}(G)$ by assumption, we actually have  $\operatorname{Fix}(G) \cap \overline{U} = \operatorname{Fix}(G^m) \cap \overline{U} = \{q\}$  for all such m. By property (I) of the fixed point index given above,  $i_X(G^m, U) = 1$ for all sufficiently large m, and in particular for all sufficiently large primes p. Thus by Lemma 3.1,  $i := i_X(G, U)$  is congruent to 1 mod p for all sufficiently large primes p. It follows that i must equal 1. Indeed, for let p be any prime greater than m. To say that  $i \equiv 1 \mod p$  is to say that i = kp + 1 for some integer k. Now choose a prime p' > |kp+1| + 1; if  $k \neq 0$ , then kp + 1 cannot be congruent to 1 mod p'.

We are now in a position to discuss our first example of an instance of (1.1) with at least two distinct slowly oscillating periodic solutions. An important aspect of this example is the fact that the trivial solution x(t) = 0,  $t \in \mathbb{R}$ , is locally asymptotically stable. The result thus bears some intuitive resemblence to the  $f'(0) > -\pi/2$  case of Theorem 2.5 (and even more so the closely related Theorem 2.1 in [28]) in that there is a periodic solution that is in some sense "between" two stable sets. Here are the main ideas, which can be applied under several sets of hypotheses other than those we formulate below. We will apply the fixed point index to the continuous compact map P on the closed, bounded, convex set K. We will exhibit two relatively open convex subsets V and U of K such that  $\overline{V} \cap \overline{U} = \emptyset$ ,  $P(\overline{V}) \subset V$ , and U satisfies the hypotheses of Corollary 3.2 (U will be an open set containing 0). We will accordingly have  $i_K(P,U) = i_K(P,V) = 1$ ; the contrapositive of (II) will then guarantee the existence of another nontrivial fixed point of P in  $K \setminus (U \cup V)$ . Since P is fixed-point free on  $\partial V$  and  $\partial U$ , this fixed point actually lies in the open set  $K \setminus (U \cup V).$ 

For the rest of this section we impose, in addition to (H), the following restriction on the feedback function  $f : \mathbb{R} \to \mathbb{R}$ :

(H') 
$$\begin{cases} |f| \text{ is bounded by some } \mu > 0;\\ \text{there are } \beta, \sigma > 0 \text{ such that } |f(x)| \ge \sigma \text{ whenever } |x| \ge \beta. \end{cases}$$

Here is the main result of the section.

PROPOSITION 3.3. Suppose that (H) and (H') hold, and moreover that

$$-\frac{\pi}{2} < f'(0) < 0 \quad and \quad \beta < \frac{\sigma}{2+\mu/\sigma}$$

Then (1.1) has at least two distinct slowly oscillating periodic solutions. Each of these solutions corresponds to a fixed point of P (and so has minimal period given by the length of a minimal interval containing three consecutive zeros).

The first step in the proof of Proposition 3.3 is to exhibit the subset  $V \subset K$  described above.

LEMMA 3.4. Suppose that (H) and (H') hold, and that

$$\beta < \frac{\sigma}{2 + \mu/\sigma}.$$

Then there is a relatively open convex subset  $V \subset K$  such that  $0 \notin \overline{V}$  and  $P(\overline{V}) \subset V$ . Accordingly,  $i_K(P, V) = 1$  and P has a fixed point in V; this fixed point is a segment of a slowly oscillating periodic solution of (1.1).

*Proof.* Note that our conditions on  $\beta$  guarantee that  $\beta/\sigma < 1/2$ . Choose  $\gamma > 0$  small enough that  $\gamma + \beta/\sigma < 1$  and

(3.1) 
$$\beta \le \frac{\sigma - (\sigma + \mu)\gamma}{2 + \mu/\sigma}$$

Now define the following set:

$$V := \{ \varphi \in K \mid \varphi(t) > \sigma t + \sigma(1 - \gamma) \text{ for } t \in [-1 + \gamma, -\beta/\sigma] \}.$$

This set is clearly open and convex, and its closure is

 $\overline{V} = \{ \varphi \in K \mid \varphi(t) \ge \sigma t + \sigma(1 - \gamma) \text{ for } t \in [-1 + \gamma, -\beta/\sigma] \}.$ 

The members of  $\overline{V}$  are all the members of K whose graphs do not go below the graph pictured in Figure 1.



Fig. 1. Definition of  $\overline{V}$ 

Suppose now that  $x : [-1, \infty) \to \mathbb{R}$  is a solution of (1.1) with  $x_0 \in \overline{V}$ . Notice that  $x(0) \ge \sigma(1-\gamma) - \beta$ , and that this latter quantity is greater than  $\beta$  by assumption (3.1). Moreover,  $x(t) \ge \beta$  for all  $t \in [-1 + (\gamma + \beta/\sigma), 0]$  since  $x_0 \in \overline{V}$ . We claim that in fact  $x(t) \ge \beta$  for all  $t \in [0, \gamma + \beta/\sigma]$  as well. For by the first point of (H'), we have  $x'(t) \ge -\mu$  for all  $t \in [0, \gamma + \beta/\sigma]$ , and so accordingly

$$x(\gamma + \beta/\sigma) \ge \sigma(1 - \gamma) - \beta - \mu(\gamma + \beta/\sigma);$$

our assumption (3.1) on  $\beta$  guarantees that the quantity on the right is greater than or equal to  $\beta$ . Therefore, at the first positive time  $\tau$  such that  $x(\tau) = \beta$ , we see that  $x(t) \ge \beta$  on the entire interval  $[\tau - 1, \tau]$ .

Hypothesis (H') and the negative feedback condition now imply that  $x'(t) \leq -\sigma$  for all  $t \in [\tau, \tau + 1]$ . Thus  $z_1$ , the first positive zero of x, occurs in the interval  $(\tau, \tau + \beta/\sigma]$ . We therefore see that  $x'(t) \leq -\sigma$  on the interval

$$[z_1, \tau + 1] \supset [z_1, z_1 + 1 - \beta/\sigma]$$

and that  $x'(t) \leq 0$  on  $[\tau + 1, z_1 + 1]$ . It follows that

$$-x_{z_1+1}(s) \ge \sigma s + \sigma > \sigma s + \sigma(1-\gamma)$$
 for all  $s \in [-1, -\beta/\sigma]$ 

and so, in particular,  $x_{z_1+1} \in -V$ .

By taking into account that the assumptions (H') are symmetric, even though f is not assumed odd, a symmetric argument shows that  $x_{z_2+1} \in V$ , which is the desired conclusion.

REMARK 3.5. Let  $p : \mathbb{R} \to \mathbb{R}$  denote the slowly oscillating periodic solution found in the last result, translated so that  $p_0 \in V \subset K$ . By imposing further conditions on f(x) for  $|x| \ge \beta$ , one can guarantee that  $p_0 \in V$  is a stable fixed point of P. One particularly simple such condition is to insist that f be constant on  $[\beta, \infty)$ . In this case, the proof of Lemma 3.4 shows that P is in fact constant on  $\overline{V}$  since, in the notation of the above proof,  $x_{\tau}$ can only be continued in one possible way. More generally, by insisting that f be Lipschitz and have a small enough Lipschitz constant on  $[\beta, \infty)$ , it is possible to show, in the spirit of Walther [40], that P is contractive on the subset  $\overline{V}$  of K.

We now turn our attention to our desired open convex subset  $U \subset K$ about  $0 \in K$ . The lemma below follows from standard results about the behavior of solutions of (1.1) near the trivial solution  $x(t) = 0, t \in \mathbb{R}$ , when the latter is locally asymptotically stable by the so-called principle of linearized stability.

LEMMA 3.6. Suppose that (H) holds and that  $f'(0) \in (-\pi/2, 0)$ . Then there is an open convex subset U of K with  $0 \in U$  such that the hypotheses of Corollary 3.2 are satisfied. More precisely, there is an open convex set  $U \subset K$  with  $0 \in U$  and some  $n \in \mathbb{N}$  such that, for all  $m \geq n$ ,

 $P^m(\overline{U}) \subset U$  and  $\operatorname{Fix}(P) \cap \overline{U} = \operatorname{Fix}(P^m) \cap \overline{U} = 0.$ 

Thus  $i_K(P,U) = 1$ . Moreover,  $U \subset K$  can be chosen so that  $\overline{U} \cap \overline{V} = \emptyset$ where V is as in Lemma 3.4.

*Proof.* Since  $f'(0) \in (-\pi/2, 0), 0 \in C$  is a locally exponentially stable stationary point of the semiflow F. Thus there are numbers  $\epsilon > 0, \kappa > 0$ , and  $\omega < 0$  such that  $||F(t, \varphi)|| \le \kappa e^{\omega t}$  whenever  $||\varphi|| \le \epsilon$  (see, for instance, [8, Chapter VII]). Set

$$U := \{ \varphi \in K \mid \|\varphi\| < \epsilon \}.$$

Notice that there is no loss of generality in taking  $\epsilon > 0$  small enough that  $\overline{U} \cap \overline{V} = \emptyset$  where V is as in Lemma 3.4. Given  $\varphi \in \overline{U}$ , we have  $x^{\varphi}(t) \to 0$  as  $t \to \infty$ ; it is therefore clear that no power of P has any fixed points on  $\overline{U}$  other than  $0 \in K$ .

Now choose T > 0 so large that  $\kappa e^{\omega T} < \epsilon$ , and let *n* be an integer larger than T/2. We claim that for any  $\varphi \in \overline{U}$  we have  $||P^n(\varphi)|| < \epsilon$ . In this situation  $x^{\varphi}$  has fewer than 2n positive zeros,  $P^n(\varphi) = 0$  by definition. Otherwise, the slowly oscillating behavior of  $x^{\varphi}$  implies that the 2nth positive

zero z of  $x^{\varphi}$  occurs after time 2n - 1 > T - 1, and thus

 $\|P^n(\varphi)\| = \|x_{z+1}^{\varphi}\| = \|F(z+1,\varphi)\| \le \kappa e^{\omega(z+1)} < \kappa e^{\omega T} < \epsilon. \quad \blacksquare$ 

Proposition 3.3 now follows from Lemmas 3.4, 3.6, and the discussion at the beginning of the section.

REMARK 3.7. Under the stronger assumption that 0 > f'(0) > -1, it is possible to use much more elementary arguments to exhibit an open neighborhood U in K about 0 such that  $P(\overline{U}) \subset U$  and  $\overline{U} \cap \overline{V} = \emptyset$ . In this case, we still have  $i_K(P, U) = 1$  by property (I) of the index above, and the rest of the proof of Proposition 3.3 is the same as before. Here is a sketch of this simpler approach.

Since f is  $C^1$  on a neighborhood of 0 and |f'(0)| < 1, there is an  $\epsilon > 0$  such that  $0 < |f'(\xi)| < 1$  for all  $\xi \in [-\epsilon, \epsilon]$ . Set

$$U := \{ \varphi \in K \mid \|\varphi\| < \epsilon \}.$$

By shrinking the constant  $\epsilon$  if necessary, we can guarantee that  $\overline{U}$  and  $\overline{V}$  are disjoint.

Now let  $\varphi \in \overline{U}$  be given with  $\varphi \neq 0$ . We claim that  $||P(\varphi)|| < ||\varphi||$ . If  $J(\varphi) < 2$  with  $J(\phi)$  introduced in Proposition 2.1,  $P(\varphi) = 0$  and the claim is immediate. Suppose that  $J(\varphi) \geq 2$  and write  $z_1 = z_1(\varphi), z_2 = z_2(\varphi)$ , and  $x = x^{\varphi}$  Since  $x(0) = ||\varphi||$  and x is decreasing on  $[0, z_1]$ , we see that  $||x_{z_1}|| \leq ||\varphi||$ . For all  $s \in [-1, 0]$ , since  $|f'(\xi)| < 1$  for  $\xi \in [-\epsilon, \epsilon]$  and f(0) = 0, we have

$$|f(x_{z_1}(s))| < |x_{z_1}(s)| \le ||\varphi||.$$

Therefore,  $0 > x'(t) > -\|\varphi\|$  for all  $t \in (z_1, z_1 + 1)$ , and  $\|x_{z_1+1}\| < \|\varphi\|$ . A symmetric argument shows that  $\|P(\varphi)\| < \|x_{z_1+1}\| < \|\varphi\|$ .

4. Solutions of long period. The interest of our next example for (1.1) with multiple slowly oscillating periodic solutions lies in the fact that we obtain solutions of "long period" even though f can be taken bounded by some fixed number: the nonmonotonicity of f is key. A similar example, that is, (1.1) with a bounded right-hand side f and periodic solutions with arbitrarily large periods, was discussed by Nussbaum [31, Theorem 3.1 & Corollary 3.1]. However, in contrast to the example in [31], here we will have both a slowly oscillating periodic solution of long period and at least one Kaplan–Yorke solution for (1.1) regardless of the stability of the zero solution.

In order to specify the feedback function  $f : \mathbb{R} \to \mathbb{R}$  in detail, we need positive reals  $a, c, \delta, \gamma$  satisfying the following conditions:

(i)  $c < \min(a, \delta);$ (ii)  $\gamma \ge 4a \text{ and } \gamma > \delta;$ (iii)  $\delta + \frac{c}{\delta}\gamma \le a.$ 

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The following result is obvious and shows that these conditions are not vacuous.

PROPOSITION 4.1. Given any a > 0 and any  $\gamma > 4a$ , there is some  $\delta_* > 0$  such that, for any  $\delta \in (0, \delta_*)$ , there is a real  $c_*(\delta) > 0$  such that all of conditions (i)–(iii) are satisfied whenever  $c \in (0, c_*(\delta))$ .

Using the reals  $a, c, \gamma, \delta > 0$ , we are now able to formulate the additional assumptions on  $f : \mathbb{R} \to \mathbb{R}$  that will be essential throughout the present section:

(H") 
$$\begin{cases} f \text{ is odd and } |f| \leq \gamma \text{ on } [0, 2a]; \\ f(x) = \begin{cases} -\gamma & \text{for } x \in [a, 2a-c], \\ -\delta & \text{for } x \in [2a, \gamma]. \end{cases}$$

Provided that, apart from (H''), the function f also satisfies the imposed standard hypothesis (H), equation (1.1) has a slowly oscillating periodic solution as we now show.

PROPOSITION 4.2. Given reals  $a, c, \gamma, \delta > 0$  satisfying conditions (i)–(iii), suppose that (H) and (H") hold. Then (1.1) has a slowly oscillating periodic solution with minimal period greater than 4 and given by the length of an interval containing three consecutive zeros.

*Proof.* The reader may find it helpful to refer to Figure 2.



Fig. 2. Proof of Proposition 4.2

1. Let  $\varphi \in C$  be given with  $\varphi(0) = -2a$  and with  $-\gamma \leq \varphi(s) \leq -2a$  for all  $s \in [-1, 0]$ , and let  $x = x^{\varphi} : [-1, \infty) \to \mathbb{R}$  denote the uniquely determined

solution of (1.1) with  $x_0 = \varphi$ . Then, in view of assumption (H"), we have  $x(t) = -2a + \delta t$  for all  $0 \le t \le 1$ . In particular, x is strictly increasing on [0, 1].

Now, note that  $c/\delta < 1$  by condition (i) and  $-2a + \delta < -a$  by condition (iii). Thus  $x(t) \ge x(c/\delta) = -2a + c$  for  $t \in [c/\delta, 1]$  and  $x(1) = -2a + \delta < -a$ . Moreover, for all  $t \in [1, 1 + c/\delta]$ , x(t) is increasing with derivative at most  $\gamma$ . We thus have the estimate

$$x(1+c/\delta) \le x(1) + \gamma \frac{c}{\delta} = -2a + \delta + \gamma \frac{c}{\delta},$$

which together with condition (iii) guarantees that  $x(t) \in [-2a + c, -a]$  for all  $t \in [c/\delta, 1 + c/\delta]$ .

2. Define  $\tau_1 := \min\{t \ge 1 + c/\delta \mid x(t) = -a\}$ . Then the last observation in combination with assumption (H") shows that  $x(\tau_1 + s) = -a + \gamma s$  for all  $s \in [0, 1]$ . Set  $\tau_2 := \tau_1 + 2a/\gamma$ . By condition (ii) we have  $\tau_2 < \tau_2 + a/\gamma < \tau_1 + 1$ , and so we obtain

$$x(\tau_2) = a, \quad x(\tau_2 + a/\gamma) = 2a, \quad x(\tau_1 + 1) = -a + \gamma.$$

Thus, as t increases from  $\tau_1$  to  $\tau_1+1$ , x(t) increases from the value  $x(\tau_1) = -a$  with constant slope  $\gamma$ , traverses the interval [-a, a], and finishes at the value  $x(\tau_1 + 1) = -a + \gamma$ . Note that  $-a + \gamma \geq 3a$  by condition (ii).

3. Now, from the oddness of the map f and the fact that  $[\tau_1, \tau_2] \ni s \mapsto x(s)$  traverses the interval [-a, a] at constant slope  $\gamma$ , it follows that  $[0, 2a/\gamma] \ni s \mapsto x(\tau_1 + 1 + s)$  traces out a symmetric arc with maximum at  $s = a/\gamma$  and a maximal value of

$$x(\tau_1 + 1 + a/\gamma) < x(\tau_1 + 1) + \gamma \frac{a}{\gamma} = (-a + \gamma) + a = \gamma.$$

Furthermore, we have

$$x(\tau_2 + 1) = x(\tau_1 + 2a/\gamma + 1) = x(\tau_1 + 1) = -a + \gamma$$

and so  $x(t) \in [3a, \gamma]$  for all  $t \in [\tau_1 + 1, \tau_2 + 1]$ .

4. On the interval [a, 2a], f is bounded below by  $-\gamma$  but, by continuity, is not identically equal to  $-\gamma$ . Accordingly,  $[0, a/\gamma] \ni s \mapsto x(\tau_2 + 1 + s)$  is decreasing from  $x(\tau_2 + 1) = \gamma - a$  to  $x(\tau_2 + 1 + a/\gamma)$  where for this last value we crudely have

$$x(\tau_2 + 1 + a/\gamma) > x(\tau_2 + 1) - \gamma \frac{a}{\gamma} \ge -a + \gamma - a = \gamma - 2a,$$

and so  $x(\tau_2 + 1 + a/\gamma) > 2a$  by condition (ii). We have established that  $x(t) \in (2a, \gamma)$  for all  $t \in (\tau_2 + a/\gamma, \tau_2 + 1 + a/\gamma]$ .

5. As t increases from  $\tau_2 + a/\gamma + 1$ , x(t) will decrease with slope  $-\delta$  until time

$$\tau_3 := \min\{t > \tau_2 + a/\gamma + 1 \mid x(t) = 2a\} > \tau_2 + 1 + a/\gamma.$$

In particular, we have  $x(t) \in [2a, \gamma]$  for all  $t \in [\tau_3 - 1, \tau_3]$ . Thus  $-x_{\tau_3}$  satisfies the same conditions as our initial condition  $\varphi$ ; by using the oddness of fand applying the arguments above again, it follows that  $x(\tau_3 + t) = -x(t)$ and secondly that

$$x(2\tau_3 + t) = -x(\tau_3 + t) = -(-x(t)) = x(t)$$

for all  $t \ge 0$ . In particular, since x has only a single zero in  $[0, \tau_3]$ , the distance between any two zeros of x in  $[0, \infty)$  is greater than or equal to  $\tau_3 > 1$ . Hence,  $\mathcal{O} := \{x(t) \mid 0 \le t \le 2\tau_3\}$  is the orbit of a slowly oscillating periodic solution of (1.1) with period  $2\tau_3$ .

6. It remains to prove that  $\tau_3 > 2$ . In order to do so, observe that

$$\tau_3 - (\tau_2 + a/\gamma + 1) \ge \frac{x(\tau_2 + a/\gamma + 1) - x(\tau_3)}{\delta} > \frac{(\gamma - 2a) - 2a}{\delta} = \frac{\gamma - 4a}{\delta}.$$

Hence,

(4.1) 
$$\tau_3 > \tau_2 + \frac{a}{\gamma} + 1 + \frac{\gamma - 4a}{\delta} \ge \tau_1 + \frac{2a}{\gamma} + \frac{a}{\gamma} + 1 + \frac{\gamma - 4a}{\delta} \ge 2 + \frac{c}{\delta} + \frac{3a}{\gamma} + \frac{\gamma - 4a}{\delta}.$$

This shows that  $2\tau_3 > 4$  and completes the proof.

REMARK 4.3. Observe that the last step of the proof indicates that, by choosing appropriate parameters  $a, c, \gamma, \delta > 0$  and adapting the map fcorrespondingly, the period  $2\tau_3$  of the resulting slowly oscillating periodic solution can be made as large as desired. For since

$$\tau_3 > 2 + \frac{\gamma - 4a}{\delta}$$

by estimate (4.1), we can make  $2\tau_3$  as large as we like by taking  $\gamma > 4a$  and then choosing  $\delta$  (and then c) small enough given parameters a and  $\gamma$ .

Of course, the slowly oscillating periodic solution found in Proposition 4.2 has a segment in the closed bounded convex set  $K \subset C$  that is a nonzero fixed point of the return map P. Moreover, this periodic solution is so strongly attractive that P is constant in each sufficiently small neighborhood of this fixed point, as we now show.

PROPOSITION 4.4. Under the hypothesis of Proposition 4.2, there is an open convex subset  $V \subset K$  such that  $0 \notin \overline{V}$  and  $P(\overline{V}) \subset V$ . Moreover, the fixed point of P in V is a segment of the slowly oscillating periodic solution of (1.1) obtained in Proposition 4.2, and P maps all of  $\overline{V}$  to this fixed point.

*Proof.* We revisit the proof of Proposition 4.2 and use the same notation in the following. To begin with, define  $s_1 := \tau_1 + a/\gamma + 1$  and  $s_2 := \tau_2 + a/\gamma + 1$ . We have  $x_{s_1} \in K$ .

Now, the proof of Proposition 4.2 shows that there are numbers  $\Delta > 0$ and  $\epsilon > 0$  such that  $x(t) \in (2a + \epsilon, \gamma - \epsilon)$  for all  $t \in [s_2 + \Delta - 1, s_2 + \Delta]$ . Therefore by continuous dependence on initial conditions there is some open set V about  $x_{s_1}$  such that  $y_0 \in \overline{V}$  implies that  $y(s) \in (2a, \gamma)$  for  $s \in [s_2 + \Delta - 1 - s_1, s_2 + \Delta - s_1]$ . The definition of f now allows y to continue in only one way; in particular, y(s) will coincide with a translate of x for  $s \geq s_2 + \Delta - s_1$ .

Note that the stability of the zero solution of (1.1) has been, so far, completely irrelevant to our discussion of the current example. On the other hand, under the assumption of local exponential stability of the zero solution we can take the same approach as in Lemma 3.6 to show the existence of multiple slowly oscillating periodic solutions of (1.1). We accordingly omit the proof of the following proposition.

PROPOSITION 4.5. Under the hypotheses of Proposition 4.2, if moreover  $f'(0) \in (-\pi/2, 0)$  then (1.1) has at least two distinct slowly oscillating periodic solutions, each with period given by three consecutive zeros.

But also in the situation where  $f'(0) < -\pi/2$ , and thus the zero solution of (1.1) is unstable, (1.1) has at least two distinct slowly oscillating periodic solutions, provided that the conditions of Proposition 4.2 are satisfied and that (for example) f' = 0 outside some open neighborhood of the origin—for then Theorem 2.5 applies. We therefore have the following result.

PROPOSITION 4.6. Suppose that the assumptions of Proposition 4.2 are satisfied, that f' = 0 outside some open neighborhood of  $0 \in \mathbb{R}$ , and that additionally  $f'(0) < -\pi/2$ . Then (1.1) has at least two distinct slowly oscillating periodic solutions.

More precisely, apart from the slowly oscillating periodic solution with minimal period greater than 4, equation (1.1) has a slowly oscillating periodic solution  $x : \mathbb{R} \to \mathbb{R}$  with minimal period 4 and the symmetry x(t) = -x(t-2) for all  $t \in \mathbb{R}$ .

REMARK 4.7. The last statement of course remains true if the assumption that f' = 0 outside some open neighborhood of the origin is replaced by any condition that makes Theorem 2.5 hold.

It is now natural to ask whether, at least under additional conditions on f, one of the slowly oscillating periodic solutions obtained in Proposition 4.5 (and so in case of a locally asymptotically stable zero solution) may also be a Kaplan–Yorke solution. In order to address this issue, we should recall some details of the approach taken in [18] for the proof of the existence of Kaplan–Yorke solutions for (1.1).

We continue to assume that f satisfies assumptions (H) and (H") and is globally  $C^1$ -smooth, and that f' = 0 outside some neighborhood of the origin. As under these conditions we have  $H(u, v) \to \infty$  as  $u^2 + v^2 \to \infty$  for the Hamiltonian H defined by (2.2), it follows that for each initial condition of the form  $(u_0, 0) \in \mathbb{R}^2$ , where  $u_0 > 0$ , equation (2.1) has a uniquely determined periodic solution  $r(\cdot; u_0) = (u(\cdot; u_0), v(\cdot; u_0)) : \mathbb{R} \to \mathbb{R}^2$  with  $r(0; u_0) = (u_0, 0)$ . Its orbit is given by  $H^{-1}(\{\alpha\})$  for  $\alpha = -\int_0^{u_0} f(s) \, ds$ . Set now, for each  $u_0 > 0$ ,

(4.2) 
$$\tau(u_0) := \min\{t > 0 \mid u(t; u_0) = 0\},$$

which is well defined since  $r(\cdot; u_0)$  traces out a periodic orbit around the origin. Moreover, in view of  $r'(0; u_0) = (0, -f(u_0))$  with  $f(u_0) < 0$  and the symmetry of H discussed above, it follows that

$$r(0; u_0) = (u_0, 0),$$
  

$$r(\tau(u_0); u_0) = (0, u_0),$$
  

$$r(2\tau(u_0); u_0) = (-u_0, 0),$$
  

$$r(3\tau(u_0); u_0) = (0, -u_0),$$
  

$$r(4\tau(u_0); u_0) = (u_0, 0) = r(0; u_0)$$

In particular, the solution  $r(\cdot; u_0)$  has period  $4\tau(u_0)$  and this period is minimal. Accordingly, in order to find a periodic solution of (1.1) we seek  $u_0 > 0$  such that  $\tau(u_0) = 1$ . The main tool for doing so will be the following lemma (see [18]).

LEMMA 4.8. Assume that  $f : \mathbb{R} \to \mathbb{R}$  is  $C^1$ -smooth, odd, and satisfies the negative feedback condition, and that f' = 0 outside some neighborhood of the origin. Then the definition (4.2) induces a continuous function  $\tau$ :  $(0, \infty) \to (0, \infty)$  with

$$\lim_{u \to 0^+} \tau(u) = -\frac{\pi}{2} \frac{1}{f'(0)} \quad and \quad \lim_{u \to \infty} \tau(u) = \infty.$$

*Proof.* Since the proof of the continuity of  $\tau$  in [18] is somewhat cursory, we give a fuller argument. Let  $u_0 > 0$  and  $\epsilon > 0$  be given. Set  $\tau_0 := \tau(u_0)$ with  $\tau(u_0)$  defined by (4.2). As  $r(\tau_0; u_0) = (0, u_0)$  and  $r(\cdot; u_0)$  intersects each half-axis transversally, there clearly is some  $0 < \tau_1 < \min\{\tau_0, \epsilon/2\}$  such that  $u(\tau_0 + t; u_0) < 0 < u(\tau_0 - t; u_0)$  for all  $0 < t \le \tau_1$ . Fix now any  $0 < \epsilon_1 < \epsilon$ satisfying both

$$\epsilon_1 < \min\{|u(\tau_0 - \tau_1; u_0)|, |u(\tau_0 + \tau_1; u_0)|\}$$

and

$$\epsilon_1 < \min\{\|r(t; u_0)\|_{\mathbb{R}^2} \mid 0 \le t \le t + \tau_1\}$$

where  $\|\cdot\|_{\mathbb{R}^2}$  denotes the Euclidean norm in  $\mathbb{R}^2$ . Then by elementary results on continuous dependence of solutions to ordinary differential equations we find some  $\delta > 0$  such that for each  $\tilde{u} > 0$  with  $|\tilde{u} - u_0| < \delta$  and all  $0 \le t \le$   $\tau_0 + \tau_1$  we have

$$\max\{|u(t;\tilde{u}) - u(t;u_0)|, |v(t;\tilde{u}) - v(t;u_0)|\} \le ||r(t;u_0) - r(t,\tilde{u})||_{\mathbb{R}^2} < \epsilon_1.$$

Consider now any  $\tilde{u} > 0$  with  $|\tilde{u} - u_0| < \delta$  and observe that on the one hand we have

$$u(\tau_0 + \tau_1; \tilde{u}) = u(\tau_0 + \tau_1; \tilde{u}) - u(\tau_0 + \tau_1; u_0) + u(\tau_0 + \tau_1; u_0)$$
  

$$\leq \|r(\tau_0 + \tau_1; \tilde{u}) - r(\tau_0 + \tau_1; u_0)\|_{\mathbb{R}^2} + u(\tau_0 + \tau_1; u_0)$$
  

$$< \epsilon_1 + u(\tau_0 + \tau_1; u_0) \le \epsilon_1 - \epsilon_1 = 0$$

and on the other hand

$$u(\tau_0 - \tau_1; \tilde{u}) = u(\tau_0 - \tau_1; \tilde{u}) - u(\tau_0 - \tau_1; u_0) + u(\tau_0 - \tau_1; u_0)$$
  
>  $\epsilon_1 - |u(\tau_0 - \tau_1; \tilde{u}) - u(\tau_0 - \tau_1; u_0)| \ge \epsilon_1 - \epsilon_1 = 0.$ 

Hence, continuity of the solution  $r(\cdot; \tilde{u})$  and so of its first component  $u(\cdot; \tilde{u})$ implies the existence of some  $\tau_0 - \tau_1 < T(\tilde{u}) < \tau_0 + \tau_1$  with  $u(T(\tilde{u}); \tilde{u}) = 0$ .

Suppose now that  $T(\tilde{u}) \neq \min\{t > 0 \mid u(t; \tilde{u}) = 0\}$ ; that is, suppose that there is some  $0 < \xi < T(\tilde{u})$  with  $u(\xi; \tilde{u}) = 0$ . Then our discussion about the solutions of (2.1) would imply the existence of some  $0 < \zeta < T(\tilde{u})$  with  $r(\zeta; \tilde{u}) = (0, -\tilde{u})$ . But in view of

$$\begin{aligned} \epsilon_1^2 &> \|r(\zeta; u_0) - r(\zeta; \tilde{u})\|_{\mathbb{R}^2}^2 \\ &= (u(\zeta; u_0) - u(\zeta; \tilde{u}))^2 + (v(\zeta; u_0) - v(\zeta; \tilde{u}))^2 \\ &= (u(\zeta; u_0))^2 + (v(\zeta; u_0) + \tilde{u})^2 \ge \|r(\zeta; u_0)\|_{\mathbb{R}^2}^2 + \tilde{u}^2 \ge \epsilon_1^2 + \tilde{u}^2 \end{aligned}$$

that is impossible. Therefore, we have  $T(\tilde{u}) = \tau(\tilde{u})$  where the right-hand side is defined by (4.2). Moreover, it follows that

$$|\tau(u_0) - \tau(\tilde{u})| = |\tau_0 - T(\tilde{u})| < \tau_1 < \frac{\epsilon}{2} < \epsilon,$$

and this proves the continuity of  $\tau: (0, \infty) \to (0, \infty)$  defined by (4.2).

The assertions about the behavior of  $\tau$  for  $t \to 0^+$  and  $t \to \infty$  are proven in [18, Theorem 1].

REMARK 4.9. Under the assumptions of the last statement it is possible to prove that  $\tau : (0, \infty) \to (0, \infty)$  is not only continuous but continuously differentiable. This follows from straightforward application of elementary results about smooth dependence of solutions of ordinary differential equations on initial conditions, along with the implicit function theorem.

Now note that if the assumptions on f from the last lemma are satisfied, and if moreover  $f'(0) < -\pi/2$  and so the zero solution of (1.1) is unstable, then  $\tau(u) < 1$  for all sufficiently small u > 0 whereas  $\tau(u) > 1$  for all sufficiently large u > 0. Thus, the intermediate value theorem implies the existence of some  $0 < u_0 < \infty$  with  $\tau(u_0) = 1$ , and so (1.1) has at least one Kaplan–Yorke solution. However, the same reasoning does not work in the situation  $-\pi/2 < f'(0) < 0$ , that is, when the zero solution of (1.1) is locally asymptotically stable. Indeed, in this case Lemma 4.8 implies that  $\tau(u) > 1$ for all sufficiently small as well as for all sufficiently large u > 0, and so a straightforward application of the intermediate value theorem fails. In our next result we will give conditions on f, similar to those Proposition 4.6 but without any restrictions on f'(0) < 0, such that the intermediate value theorem will be applicable regardless of the stability properties of the zero solution of (1.1). In doing so, we shall refer to Figure 3, which describes the value of the vector field for (2.1) with f satisfying (H") in various regions of the first quadrant of the real plane.



Fig. 3. Value of the vector field for (2.1) in various regions for  $u, v \ge 0$ 

PROPOSITION 4.10. Given reals  $a, c, \gamma, \delta > 0$  satisfying conditions (i)–(iii), assume that f is  $C^1$ -smooth and satisfies both (H) and (H"), and that f' = 0 outside some open neighborhood of  $0 \in \mathbb{R}$ . If

(4.3) 
$$\frac{1}{\gamma} \int_{0}^{a} |f(s)| \, ds < a - c$$

then  $\tau(2a-c) < 1$  with  $\tau(2a-c)$  defined by (4.2).

*Proof.* Under the given assumptions, consider the periodic solution  $r(\cdot; 2a - c) : \mathbb{R} \to \mathbb{R}^2$  of (2.1) with initial condition r(0; 2a - c) = (2a - c, 0). We have

$$r'(0;2a-c) = \begin{pmatrix} u'(0;2a-c) \\ v'(0;2a-c) \end{pmatrix} = \begin{pmatrix} f(0) \\ -f(2a-c) \end{pmatrix} = \begin{pmatrix} 0 \\ \gamma \end{pmatrix}$$

and u'(t; 2a - c) < 0 for all sufficiently small t > 0. Hence, as t moves forward, as long as  $a \le u(t) \le 2a - c$  and  $0 \le v(t) \le a$ , we will have

$$\binom{u'(t;2a-c)}{v'(t;2a-c)} = \binom{f(v(t;2a-c))}{\gamma}$$

Thus, for all such t, we actually have  $v(t; 2a - c) = \gamma t$  and

$$u(t; 2a - c) = u(0; 2a - c) + \int_{0}^{t} f(\gamma s) \, ds = 2a - c + \frac{1}{\gamma} \int_{0}^{\gamma t} f(s) \, ds.$$

Now condition (4.3) implies that  $a \le u(t; 2a-c) \le 2a-c$  for all  $0 \le t \le a/\gamma$ , and so at precisely  $t = a/\gamma$  we have

$$u(t; 2a - c) = 2a - c + \frac{1}{\gamma} \int_{0}^{a} f(s) \, ds \in [a, 2a - c] \text{ and } v(t; 2a - c) = a.$$

We see that the solution  $r(\cdot; 2a - c)$  will move diagonally across the "box"  $[a, 2a - c] \times [a, 2a - c]$ , with constant velocity

$$r'(t;2a-c) = \begin{pmatrix} u'(t;2a-c) \\ v'(t;2a-c) \end{pmatrix} = \begin{pmatrix} -\gamma \\ \gamma \end{pmatrix}$$

This traversal will clearly take fewer than  $(a - c)/\gamma$  time units. Then, by symmetry, the solution  $r(\cdot; 2a - c)$  will take another  $a/\gamma$  units to reach r(t; 2a - c) = (0, 2a - c). Therefore

$$\tau(2a-c) < \frac{a+(a-c)+a}{\gamma} < \frac{3a}{\gamma} < 1$$

by condition (ii).  $\blacksquare$ 

A direct consequence of the last result is now the following corollary.

COROLLARY 4.11. If the hypotheses on f described in Proposition 4.10 (condition (4.3) included) hold, then (1.1) has a Kaplan–Yorke solution with maximal value larger than 2a - c.

If furthermore  $f'(0) > -\pi/2$ , then (1.1) also has a Kaplan–Yorke solution with maximum value less than 2a - c.

*Proof.* Under the given conditions, Lemma 4.8 together with the last proposition show that for the continuous map  $\tau : (0, \infty) \to (0, \infty)$  defined by (4.2) we have  $\tau(2a - c) < 1$  whereas  $\tau(u) > 1$  for all sufficiently large

u > 2a - c. Hence, there clearly exists some  $2a - c < u_0 < \infty$  with  $\tau(u_0) = 1$ such that  $x : \mathbb{R} \to \mathbb{R}$  with  $x(t) := u(t; u_0)$  for  $t \in \mathbb{R}$  forms a Kaplan–Yorke solution of (1.1) with  $\max_{t \in \mathbb{R}} x(t) = u_0 > 2a - c$ , as claimed.

Under the additional assumption  $f'(0) > -\pi/2$ , Lemma 4.8 also proves that  $\tau(u) > 1$  for all sufficiently small u > 0. Hence, apart from  $2a - c < u_0 < \infty$ , there is some  $0 < \tilde{u}_0 < 2a - c$  with  $\tau(\tilde{u}_0) = 1$ . Therefore  $\tilde{x} : \mathbb{R} \to \mathbb{R}$ defined by  $\tilde{x}(t) := u(t; \tilde{u}_0)$  for all  $t \in \mathbb{R}$  is a Kaplan–Yorke solution of (1.1) as well, but with maximum value  $\max_{t \in \mathbb{R}} \tilde{x}(t) = \tilde{u}_0 < 2a - c$ .

Let us consider an example. We take f as described before Proposition 4.2, with  $\gamma = 4$ , a = 1,  $\delta = 2/3$ , and c = 1/20. Note that conditions (i)–(iii) are satisfied. In addition to hypotheses (H"), we also take f(x) = -2x on [-a+c, a-c] (so the zero solution is unstable). Now, Proposition 4.10 shows that (1.1) will have a Kaplan–Yorke solution q with maximum value greater than 2a - c. On the other hand, in this case a solution r = (u, v) of (2.1) with initial condition (3a, 0) = (3, 0) will have period strictly greater than 4. Referring once more to Figure 3, we see that as long as (u(t), v(t)) is in the "box" bounded by the lines y = 0, y = a - c, x = 2a, and  $x = \gamma$ , we will have  $v'(t) = \delta$  and  $v(t) = \delta t = 2t/3$ , and

$$u'(t) = f(v(t)) = -2v(t) = -2\delta t, \quad u(t) = 3 - \delta t^2 = 3 - \frac{2}{3}t^2.$$

Since c = 1/20, we see that r(1) is still in the above-described "box", and hence r has period strictly greater than 4. We conclude that our Kaplan– Yorke solution q has a trace in the (x(t), x(t-1)) plane going through the points  $(\alpha, 0)$  and  $(0, \alpha)$ , where  $2a - c < \alpha < 3a$ .



Fig. 4. Graphs of p and q

On the other hand, the proof of Proposition 4.2 shows that the slowly oscillating periodic solution p found there satisfies  $|p'(z)| = \gamma$  and |p(z-1)|

 $\in [a, 2a - c]$  whenever p(z) = 0, and attains a maximum value of more than 3a. Thus the trace of p in the (x(t), x(t-1)) plane is a simple closed curve going through the points  $(\alpha_2, 0)$  and  $(0, \alpha_1)$ , where  $\alpha_1 < \alpha < \alpha_2$ . Thus the traces of q and p intersect in the (x(t), x(t-1)) plane; such an intersection of the traces of two SOP solutions is not possible when f is monotonic (again, see Mallet-Paret and Sell [25]).

Figures 4 and 5 display numerical approximations of two solutions q and p as described just above. The first figure shows the graphs of the two solutions; the second shows the traces of the solutions in the plane. (In making the approximations, we have taken f linear on the intervals [a-c, a] and [2a-c, 2a], and so there are isolated points where f is not smooth.)



Fig. 5. Traces of p and q in the (x(t), x(t-1)) plane

5. Third example: many SOP solutions. An idea similar to that in this section is described at the end of Vas's paper [36] dealing with SOP solutions for a delay equation with negative feedback and instantaneous damping term. Observe that, other than the differentiability of f at 0 and the negative feedback condition itself, the only restriction on f in [0, 2a] required in Proposition 4.2 is that  $f(x) \in [-\gamma, 0)$  for all  $x \in (0, 2a]$ . This observation now allows us to specify f at different "scales", and so create an equation with as many periodic solutions as we wish.

Let  $\gamma_n, n \in \mathbb{N}$ , be a sequence of positive reals satisfying  $\gamma_{n+1} < \gamma_n/4$ . Then, for each  $\gamma_n$ , let us take (for example)

$$a_n = \gamma_n/4, \quad \delta_n = \gamma_n/8, \quad c_n = \gamma_n/64.$$

Then  $\gamma_n$ ,  $a_n$ ,  $\delta_n$ ,  $c_n$  together satisfy conditions (i)–(iii) above. Given any  $M \in \mathbb{N} \cup \{\infty\}$ , with M > 1, there exists a bounded, odd, continuous function f satisfying both the negative feedback condition and the following version of (H") for all n < M:

(5.1) 
$$\begin{cases} f \text{ is odd and } |f| \leq \gamma_n \text{ on } [0, 2a_n]; \\ f(x) = \begin{cases} -\gamma_n & \text{for } x \in [a_n, 2a_n - c_n], \\ -\delta_n & \text{for } x \in [2a_n, \gamma_n]. \end{cases}$$

If  $M < \infty$ , we can moreover take f to be differentiable at 0. According to our work in Section 4, for such f equation (1.1) has M - 1 stable SOP solutions whose segments are fixed points of P.

Suppose  $\gamma_1 = 5$  and  $\gamma_2 = 1$ , and  $a_n, c_n, \delta_n$  are as described above for  $n \in \{1, 2\}$ . We take f piecewise linear. The two stable periodic solutions are readily seen in numerical approximation; see Figure 6.



Fig. 6. Two stable periodic solutions

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## References

- W. Alt, Periodic solutions of some autonomous differential equations with variable time delay, in: H. O. Peitgen and H.-O. Walther (eds.), Functional Differential Equations and Approximation of Fixed Points (Bonn, 1978), Lecture Notes in Math. 730, Springer, Berlin, 1979, 16–31.
- [2] B. Bánhelyi, T. Csendes, T. Krisztin, and A. Neumaier, Global attractivity of the zero solution for Wright's equation, SIAM J. Appl. Dynam. Systems 13 (2014), 537–563.
- F. E. Browder, On a generalization of the Schauder fixed point theorem, Duke Math. J. 26 (1959), 291–303.
- F. E. Browder, A further generalization of the Schauder fixed point theorem, Duke Math. J. 32 (1965), 575–578.
- Y. Cao, Multiexistence of slowly oscillating periodic solutions for differential delay equations, SIAM J. Math. Anal. 26 (1995), 436–445.
- Y. Cao, Uniqueness of periodic solution for differential delay equations, J. Differential Equations 128 (1996), 46–57.
- [7] S.-N. Chow and H.-O. Walther, Characteristic mutipliers and stability of symmetric periodic solutions of  $\dot{x}(t) = g(x(t-1))$ , Trans. Amer. Math. Soc. 307 (1988), 127–142.
- [8] O. Diekmann, S. A. van Gils, S. M. Verduyn Lunel, and H.-O. Walther, *Delay Equations: Functional, Complex, and Nonlinear Analysis*, Appl. Math. Sci. 110, Springer, 1995.
- P. Dormayer, Smooth bifurcation of symmetric periodic solutions of functional differential equations, J. Differential Equations 82 (1989), 109–155.
- [10] P. Dormayer, Smooth symmetry breaking bifurcation for functional differential equations, Differential Integral Equations 5 (1992), 831–854.
- [11] P. Dormayer and B. Lani-Wayda, Floquet multipliers and secondary bifurcations in functional differential equations. Numerical and analytical results, Z. Angew. Math. Phys. 46 (1995), 823–858.
- [12] A. Granas and J. Dugundji, *Fixed Point Theory*, Springer Monogr. Math., Springer, New York, 2003.
- [13] K. P. Hadeler and J. Tomiuk, Periodic solutions of difference-differential equations, Arch. Ration. Mech. Anal. 65 (1977), 87–95.
- [14] A. F. Ivanov, B. Lani-Wayda, and H.-O. Walther, Unstable hyperbolic periodic solutions of differential delay equations, in: Recent Trends in Differential Equations, World Sci. Ser. Appl. Anal., World Sci., River Edge, NJ, 1992, 301–316.
- [15] A. F. Ivanov and J. Losson, Stable rapidly oscillating solutions in delay equations with negative feedback, Differential Integral Equations 12 (1999), 811–832.
- [16] G. S. Jones, The existence of periodic solutions of  $f'(x) = -\alpha f(x-1)\{1+f(x)\},$ J. Math. Anal. Appl. 5 (1962), 435–450.
- [17] G. S. Jones, On the nonlinear differential-difference equation  $f'(x) = -\alpha f(x-1) \cdot \{1 + f(x)\}$ , J. Math. Anal. Appl. 4 (1962), 440–469.
- [18] J. L. Kaplan and J. A. Yorke, Ordinary differential equations which yield periodic solutions of differential delay equations, J. Math. Anal. Appl. 48 (1974), 317–324.
- [19] J. L. Kaplan and J. A. Yorke, On the stability of a periodic solution of a differential delay equation, SIAM J. Math. Anal. 6 (1975), 268–282.
- [20] Y. Kuang and H. L. Smith, Slowly oscillating periodic solutions of autonomous state-dependent delay equations, Nonlinear Anal. 19 (1992), 855–872.
- B. Lani-Wayda, Erratic solutions of simple delay equations, Trans. Amer. Math. Soc. 351 (1999), 901–945.

- [22] J.-P. Lessard, Recent advances about the uniqueness of the slowly oscillating periodic solutions of Wright's equation, J. Differential Equations 248 (2010), 992–1016.
- [23] E. Liz, M. Pinto, G. Robledo, S. Trofimchuk and V. Tkachenko, Wright type delay differential equations with negative Schwarzian, Discrete Contin. Dynam. Systems 9 (2003), 309–321.
- [24] E. Liz, V. Tkachenko, and S. Trofimchuk, Yorke and Wright 3/2-stability theorems from a unified point of view, Discrete Contin. Dynam. Systems 2003, Suppl. Vol., 580–589.
- [25] J. Mallet-Paret and G. R. Sell, The Poincaré-Bendixson theorem for monotone cyclic feedback systems with delay, J. Differential Equations 125 (1996), 441–489.
- [26] J. Mallet-Paret and H.-O. Walther, Rapid oscillations are rare in scalar systems governed by monotone negative feedback with a time lag, preprint.
- [27] S. H. Marshall and D. R. Smith, Feedback, control, and the distribution of prime numbers, Math. Mag. 86 (2013), 189–203.
- [28] R. D. Nussbaum, Periodic solutions of some nonlinear autonomous functional differential equations II, J. Differential Equations 14 (1973), 360–394.
- [29] R. D. Nussbaum, Periodic solutions of some nonlinear autonomous functional differential equations, Ann. Mat. Pura Appl. 101 (1974), 263–306.
- [30] R. D. Nussbaum, A Global bifurcation theorem with applications to functional differential equations, J. Funct. Anal. 19 (1975), 319–338.
- [31] R. D. Nussbaum, Uniqueness and nonuniqueness for periodic solutions of x'(t) = -g(x(t-1)), J. Differential Equations 34 (1979), 25–54.
- [32] H. Peters, Chaotic behaviour of nonlinear differential-delay equations, Nonlinear Anal. 7 (1983), 1315–1334.
- [33] H.-W. Siegberg, Chaotic behavior of a class of differential delay equations, Ann. Mat. Pura Appl. 138 (1984), 15–33.
- [34] A. L. Skubachevskii and H.-O. Walther, On the Floquet multipliers of periodic soutions to non-linear functional differential equations, J. Dynam. Differential Equations 18 (2006), 257–355.
- [35] D. Stoffer, Delay equations with rapidly oscillating stable periodic solutions, J. Dynam. Differential Equations 20 (2008), 201–238.
- [36] G. Vas, Infinite number of stable periodic solutions for an equation with negative feedback, Electron. J. Qual. Theory Differential Equations 2011, no. 18, 20 pp.
- [37] H.-O. Walther, Density of slowly oscillating solutions of  $\dot{x}(t) = -f(x(t-1))$ , J. Math. Anal. Appl. 79 (1981), 127–140.
- [38] H.-O. Walther, Bifurcation from periodic solutions in functional differential equations, Math. Z. 182 (1983), 269–289.
- [39] H.-O. Walther, The 2-dimensional attractor of  $x'(t) = -\mu x(t) + f(x(t-1))$ , Mem. Amer. Math. Soc. 113 (1995), no. 544.
- [40] H.-O. Walther, Contracting return maps for monotone delayed feedback, Discrete Contin. Dynam. Systems 7 (2001), 515–540.
- [41] H.-O. Walther, A periodic solution of a differential equation with state-dependent delay, J. Differential Equations 244 (2008), 1910–1945.
- [42] H.-O. Walther, Topics in delay differential equations, Jahresber. Deutsch. Math.-Verein. 116 (2014), 87–114.
- [43] E. M. Wright, A nonlinear difference-differential equation, J. Reine Angew. Math. 194 (1955), 66–87.
- [44] X. Xie, Uniqueness and stability of slowly oscillating periodic solutions of delay equations with bounded nonlinearity, J. Dynam. Differential Equations 3 (1991), 515–540.

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- [45] X. Xie, The multiplier equation an its application to s-solutions of a differential delay equation, J. Differential Equations 95 (1992), 259–280.
- [46] X. Xie, Uniqueness and stability of slowly oscillating periodic solutions of delay equations with unbounded nonlinearity, J. Differential Equations 103 (1993), 350– 374.
- [47] J. A. Yorke, Asymptotic stability for one dimensional differential-delay equations, J. Differential Equations 7 (1970), 189–202.

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