

**2.9.** LEMMA. If  $x_0^* \in \Re(X)$ , then

$$\pi(x_0^*) = \inf\{M > 0: \bigwedge_{x \in X} |x_0^*(x^2)| \leqslant M |x_0^*(x)| ||x||\} \leqslant ||2e_1 - x_0^*(e_1^2)e|| \cdot ||x_0^*||$$

We shall omit the proofs of lemmas.

Proof of Theorem 2.7. Let  $x_n^* \in \Re(X)$  and  $x_n^* \to x_0^*$ ; then

$$x_0^*(e) = \lim_{n \to \infty} x_n^*(e) = 0$$
 and  $x_0^*(e_1) = \lim_{n \to \infty} x_n^*(e_1)$ .

From Lemmas 2.8 and 2.9 we have

$$\begin{aligned} |x_n^*(x^2)| &\leqslant \pi(x_n^*) |x_n^*(x)| ||x|| \leqslant ||2e_1 - x_n^*(e_1^2) e|| \, ||x_n^*|| \, ||x|| \, ||x_n^*(x)|| \\ &\leqslant ||2e_1 - x_n^*(e_1^2) e|| \cdot M \cdot ||x|| \, ||x_n^*(x)|| \end{aligned}$$

Hence from the last inequality we obtain (as  $n \to \infty$ )

$$|x_0^*(x^2)| \leqslant ||2e_1 - x_0^*(e_1^2)e|| \cdot M |x_0^*(x)| ||x|| \leqslant M^* |x_0^*(x)| ||x||.$$

Consequently, if  $x_0^*(x) = 0$ , then  $x_0^*(x^2) = 0$  too, which completes the proof.

2.10. COROLLARY. We have

$$\overline{\left\{x^{*} \epsilon X^{*} \colon x^{*}(x) = x^{*}(e_{1}) \frac{g_{1}(x) - g_{2}(x)}{g_{1}(e_{1}) - g_{2}(e_{1})}, g_{1}, g_{2} \epsilon \mathfrak{M}\left(X\right), g_{1} \neq g_{2}\right\}} \subset \mathfrak{R}(X),$$

where the line on the left-hand side denotes weak sequential closure.

Now we can reformulate Problem 1 as follows:

PROBLEM 1'. Does the equality

$$\Re(X) = \left\{ x^* \in X^* \colon x^*(x) = \frac{g_1(x) - g_2(x)}{g_1(e_1) - g_2(e_2)}, g_1, g_2 \in \mathfrak{M}(X), \quad g_1 \neq g_2 \right\}$$

hold?

Problem 1 (or the equivalent Problem 1') can be generalized in the following manner.

Let X be an arbitrary B-algebra with unit. Let

$$U(X) = \{x^* \in R(X) : x^* = c(g_1 - g_2), g_1, g_2 \in \mathfrak{M}(X), g_1 \neq g_2, c \neq 0\}$$

$$\text{(evidently } U(X) = \Big\{ \! x_0^* \, \epsilon \Re(X) \colon \! \! \! \bigvee_{x_0 \in X} \! \! x_0^*(x_0) \, \neq 0 \, , \, 3 \left( \! \frac{x_0^*(x_0^2)}{2} \! \right)^{\! 2} \neq x_0^*(x_0) x_0^*(x_0^3) \! \Big\} ).$$

PROBLEM 2. Is it true that  $\overline{U(X)} = \Re(X)$  (where  $\overline{U(X)}$  denotes, as in the previous case, weak sequential closure U(X))?

Reçu par la Rédaction le 13. 4. 1967

# A modern version of the E. Noether's theorems in the calculus of variations, I

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## INTRODUCTION

A growing interest in problems connected with the Noether's theorems can be noticed in the last years. It is a result of a popularity of the Lagrange approach to physical theories.

In the period of time following the original paper of E. Noether [10] there were written many works developing its subject (Bessel-Hagen [2], Rosenfeld [11]) or treating about some mathematical foundations (de Donder [4]). Since 1950 we have had a lot of papers due as well to mathematicians as to physicists concerning those problems (Hill [9], Bergman and Schiller [1], Trautman [16], Fletcher [6], Schmutzer [12], Edelen [5], Funk [7], Steidel [13] and [14], Trautman [17] and [18], Demmig [3]).

In spite of the fact that the Lagrange formalism has a geometrical sense the authors use at every level of considerations a coordinate system for a description of geometrical objects. (One of such geometrical objects is a function, i.e. a scalar field, which being defined at the points of a space can not be considered as a function of their coordinates.) Such concept does not make easier to set off the gist of the matter and sometimes leads to misunderstandings.

As we will see, the notions of a differentiable manifold, a vector bundle and a jet-bundle are very useful in geometrical formulating of the variational problems (1).

A general variation of a functional defined on cross-sections of a finite-dimensional vector bundle M is considered, wherein the variation of the functional is induced by a variation of the cross-section, i.e. a one-parameter "smooth" family of cross-sections. It is easy to see that all such variations of a cross-section can be given by differentiable vector

<sup>(1)</sup> It is pointed out also in the recently appeared paper, Noether equations and conservation laws by A. Trautman in Commun. Math. Phys. 6 (1967), p. 248-261 (added in proof).

fields on the bundle M which preserve the fibre structure. Hence we are interested only in such vector fields. Our functional is so defined that it does not depend directly on cross-sections of the bundle M but on induced by them cross-sections of the jet-bundle J(M).

We get an analogue of the Euler-Lagrange equations as a necessary and sufficient condition of invariance of the functional with respect to variations of a cross-section when these vanish on the boundary of the domain of a cross-section.

A goal of this paper is to formulate The First E. Noether's Theorem, i.e. to state a necessary and sufficient condition for the functional to remain invariant with respect to a given general variation and for every restriction of a cross-section satisfying "the Euler-Lagrange condition".

At this time I would like to express my thanks to Dr. W. Tulczyjew for his help and many profitable conversations. I am also indebted to Professor K. Maurin and Professor A. Trautman for their lively interest and early encouragement.

#### PRELIMINARIES

Let M be a vector bundle with a base E (paracompact, orientable, n-dimensional differentiable manifold of class  $C^{\infty}$ ) and with a standard fibre  $F = \mathbb{R}^1$ . The following considerations concern also finite-dimensional vector bundles. (It is sufficient to treat some objects not like scalars but like vectors). The assumption  $F = \mathbb{R}^1$  is made for simplification of the notation. Let  $\pi$  be a projection of M on E. Different coordinate charts of the manifold M will be distinguished by indices  $I, J, \ldots$  If  $\stackrel{J}{U} \subset E$ , then on  $\pi^{-1}(\stackrel{J}{U})$  we have defined the coordinate map  $M = \pi^{-1}(\stackrel{J}{U}) \ni m \to I$  and I and I

The coordinate charts of the manifold E will bear also indices I, J, ... and to a domain U there corresponds the coordinate chart (x, U) where  $E \supset U \ni p \to \varkappa(p) = (\xi^1(p), \ldots, \xi^n(p)) \in \mathbb{R}^n$ . When we have some fixed coordinate chart, we will skip the indices. By  $\Gamma(M)$  we denote the space of differentiable cross-sections of the bundle M with relatively compact domains. By  $D_u$  and  $R_u$  we denote respectively the domain and the graph of a cross-section  $u \in \Gamma(M)$ . Let B be any vector bundle over E and  $\Omega \subset E$ ; then by  $\Gamma(\Omega, B)$  we denote the set of cross-sections of the bundle B, global over  $\Omega$ . Let  $(\varkappa, U)$  be a coordinate chart on E,  $u \in \Gamma(M)$  and  $p \in D_u \cap U$ ; then by  $[u]_p$  we denote the set of cross-sections  $u' \in \Gamma(M)$  such that  $p \in D_u'$ ,

u'(p) = u(p) and

$$\begin{split} \left(\frac{\partial}{\partial \, \xi^i} \, \varrho \circ u' \circ \varkappa^{-1}\right) & (\varkappa(p)) = \left(\frac{\partial}{\partial \, \xi^i} \, \varrho \circ u \circ \varkappa^{-1}\right) (\varkappa(p)), \quad i = 1, \ldots, n, \\ J^1_p : &= \bigcup_{u \in I(M)} [u]_p, \quad J^1 := \bigcup_{p \notin E} J^1_p. \end{split}$$

 $J^1=:J^1(M)$  is called the first order jet-bundle of the vector bundle M. Taking into account, in the definition of  $[u]_p$ , derivatives higher than the first order we get higher order jet-bundles. The results of this paper obtained for  $J^1(M)$  can be generalized for the case  $J^k(M)$ , where k indicates the order of the jet-bundle. In the following  $J^1(M)$  will be denoted by J. The map  $J \circ [u]_p \to \pi([u]_p) := p \circ E$  is a projection of the bundle J on E. It will be denoted by the same symbol as the projection of the bundle M on E. In the manifold J we introduce coordinate charts  $\begin{pmatrix} I \\ \chi \end{pmatrix}, \pi^{-1}(\stackrel{I}{U}) \end{pmatrix}$ 

$$egin{aligned} J &\supset \pi^{-1}(\overset{I}{U}) \circ \llbracket u 
rbracket_p & \supset \chi(\llbracket u 
rbracket_p) \ &:= \left(\overset{I}{x^1}(\llbracket u 
rbracket_p), \ldots, \overset{I}{x^n}(\llbracket u 
rbracket_p), \overset{I}{v}(\llbracket u 
rbracket_p), \overset{I}{v_1}(\llbracket u 
rbracket_p), \ldots, \overset{I}{v_n}(\llbracket u 
rbracket_p) 
ight) \epsilon R^{2n+1}, \end{aligned}$$

where

$$egin{aligned} & \stackrel{I}{x^i}(\llbracket u 
rbracket_p) \colon = \stackrel{I}{\xi^i}(p) \,, & \stackrel{I}{v}(\llbracket u 
rbracket_p) \colon = \stackrel{I}{r_p} \circ u(p) \,, \ & \stackrel{I}{v_i}(\llbracket u 
rbracket_p) \colon = \left(rac{\partial}{\partial \xi^i} \stackrel{I}{arphi} \circ u \circ \stackrel{I}{arphi^{-1}}
ight) \left(\stackrel{I}{arphi}(p)
ight) \,. \end{aligned}$$

If  $u \in \Gamma(M)$ , then  $\tilde{u} := \{[u]_p : p \in D_u\} \subset J$  is an *n*-dimensional differentiable manifold. On  $\tilde{u}$  we introduce coordinate charts bearing also the indices  $I, J, \ldots$  A coordinate map on  $\{[u]_p : p \in \tilde{U} \cap D_u\} \subset \tilde{u}$  is given by

$$\widetilde{u} \ni [u]_p \to (\widetilde{\zeta}^1([u]_p), \dots, \widetilde{\zeta}^n([u]_p)) \in \mathbf{R}^n,$$

where  $\overset{I}{\zeta}{}^{i}([u]_{p}):=\overset{I}{\xi}{}^{i}(p).$ 

Let  $i_{\widetilde{u}}$  be the imbedding of  $\widetilde{u}$  in J defined by  $i_{\widetilde{u}}([u]_p) := [u]_p$ . By  $\pi_{\widetilde{u}}$  we denote the 1-1 map of  $D_u$  on  $\widetilde{u}$  defined by  $\pi_{\widetilde{u}}(p) := [u]_p$ . If  $f \in C(J, \mathbf{R}^1)$ , then

$$\bar{f} := f \circ i_{\tilde{u}} \circ \pi_{\tilde{u}} \in C(D_u, \mathbf{R}^1).$$

When on a manifold P we have a field  $\Phi$  of curves (a curve  $\Phi(p)$  through  $p \in P$ , i.e.  $\Phi_0(p) = p$ ), then by  $[\Phi]$  ( $[\Phi(p)]$ ) we denote the vector field (the tangent vector at  $p \in P$ ) induced by it.

Let P, Q be differentiable manifolds and  $f \in C^{\infty}(P, Q)$ ; then

$$f^*: \bigwedge^k T^*(Q) \to \bigwedge^k T^*(P), \quad f_*: \bigwedge^k T(P) \to \bigwedge^k T(Q)$$

are the maps canonically induced by f.

#### THE FIRST NOETHER'S THEOREM

We introduce the notion of a differentiable curve in the space  $\Gamma(M)$ . Let Cu(M) be the set of fields  $\Psi$  that

$$\bigwedge_{\Psi \in \operatorname{Gu}(M)} \bigvee_{\varepsilon > 0} \left[ -\varepsilon, \varepsilon \right] \times M \ni (t, m) \to \Psi_t(m) \in M$$

differentiable and

1. 
$$\Psi_0(m) = m$$
,

2. 
$$(\pi(m) = \pi(m')) \Rightarrow (\pi \circ \Psi_t(m)) \underset{t \in [-s,s]}{=} \pi \circ \Psi_t(m')$$
.

The map  $[-\varepsilon,\,\varepsilon]$   ${}^{\flat}t \to \mathscr{\Psi}_t(m)\,\epsilon\,M$  is a differentiable curve in M through  $m\,\epsilon\,M$ .

Every field  $\Psi \epsilon Cu(M)$  defines the field  $\pi \Psi$  of curves in E

$$[-\varepsilon,\varepsilon] \times E \ni (t,p) \to (\pi \Psi)_t(p) := \pi \circ \Psi_t(m_p) \in E$$

where  $m_p \in M$  is such that  $\pi(m_p) = p$ .

The map  $[-\varepsilon, \varepsilon]^{\flat}t \to \mathring{\mathcal{Y}}_t(u) \, \varepsilon \, \varGamma(M), \, u \, \varepsilon \, \varGamma(M), \, \text{where} \, R_{\mathring{\mathcal{Y}}_t(u)} := \mathscr{Y}_t(R_u),$  is called a differentiable curve in  $\varGamma(M)$  through  $u \, \varepsilon \, \varGamma(M)$  and induced by the field  $\mathscr{Y}_{\varepsilon} \, Cu(M)$ .

By an integral functional on  $\Gamma(M)$  we mean a map

$$\Gamma(M) \circ u \to \mathscr{J}(u) := \int_{D_u} \mathscr{L}_u \, \epsilon \mathbf{R}^1,$$

where  $\mathcal{L}_u$  is defined by a differentiable map

$$J \ni [u]_p \to L([u]_p) \in \bigwedge^n T_p^*(E) \subset \bigwedge^n T^*(E)$$

as follows:

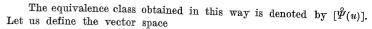
$$E \ni p \to \mathscr{L}_u(p) := L([u]_p) \in \bigwedge^n T^*(E).$$

Taking various maps L we get various integral functionals  $\mathscr{J}$ . An integral functional  $\mathscr{J}$  is called differentiable at  $u \in \Gamma(M)$  if for every  $\Psi \in Cu(M)$  there exists

$$\frac{d}{dt} \mathscr{J} \circ \mathring{\Psi}_t(u)|_{t=0}.$$

By  $\mathscr{F}$  we denote the algebra spanned by all the integral functionals differentiable on  $\Gamma(M)$ . The curves  $\mathring{\varPsi}(u)$  and  $\mathring{\varPhi}(u)$ , where  $\varPsi$ ,  $\varPhi \in Cu(M)$ , are called equivalent if

$$\bigwedge_{\mathbf{f}\in\mathcal{F}}\frac{d}{dt}\,\mathscr{F}\circ\mathring{\varPsi}_t(u)|_{t=0}=\frac{d}{dt}\,\mathscr{F}\circ\mathring{\varPhi}_t(u)|_{t=0}.$$



$$\mathfrak{W}_u(\Gamma(M)) := \{ [\mathring{\varPsi}(u)] : \varPsi_{\epsilon} Cu(M) \}$$

of linear functionals on  $\mathcal{F}$ , where, if  $\mathcal{J} \in \mathcal{F}$ , then

$$\langle \mathscr{J}, [\mathring{\varPsi}(u)] 
angle := rac{d}{dt} \, \mathscr{J} \circ \mathring{\varPsi}_t(u)|_{t=0} \, .$$

The vectors of the space  $\mathfrak{V}_u(\varGamma(M))$  will be denoted by  $[\mathring{\varPsi}(u)]$  or  $\mathfrak{X}_u$ . It is easy to see that  $\mathfrak{X}_u \epsilon \mathfrak{V}_u(\varGamma(M))$  satisfies the Leibniz formula. Hence it can be considered as a kind of a vector tangent to  $\varGamma(M)$  at  $u \epsilon \varGamma(M)$  provided that an analogue of a differentiable structure of a manifold is so chosen for  $\varGamma(M)$  that it coincides with  $\mathscr{F}$ . Any pair  $(u, \varPsi) \epsilon \varGamma(M) \times \mathcal{C}u(M)$  induces the field  $\mathring{\varPsi}$  of curves in J through points of submanifold  $\tilde{u} \subset J$ , as follows:

$$[-\varepsilon,\varepsilon]\times \tilde{u} \ni (t,[u]_p) \to \mathring{\varPsi}_t([u]_p) := [\mathring{\varPsi}_t(u)]_{(\pi^p)_t(p)} \in J.$$

The map  $\mathring{\varPsi}([u]_p)$ :  $[-\varepsilon, \varepsilon] \to J$  is a differentiable curve in J through  $[u]_p \in \mathring{u}$ . The curve  $\mathring{\varPsi}([u]_p)$  defines the vector  $X_{[u]_p} \in T_{[u]_p}(J)$ . The field on  $\widetilde{u}$  of vectors tangent to J defined by the field  $\mathring{\varPsi}$  of curves in J will be denoted by  $[\mathring{\varPsi}]$ . For interrelations between components of this field see Appendix.

Now we define a map which will be needed in the following. By j we denote an element of the bundle J. Let us recall that

$$J \ni j \to L(j) \in \bigwedge^n T^*_{\pi(j)}(E) \subset \bigwedge^n T^*(E)$$
.

We know that  $\pi^*$ :  $\bigwedge^n T^*(E) \to \bigwedge^n T^*(J)$ . Then we define

$$\pi_i^*: \bigwedge^n T_{\pi(i)}^*(E) \to \bigwedge^n T^*(J)$$

as follows: if  $\omega \in \bigwedge^n T^*_{\pi(j)}(E)$ , then  $\pi^*_j(\omega) := (\pi^*\omega)(j) \in \bigwedge^n T^*_j(J)$ . Now we define  $\mathcal{L} \in \Gamma(J, \bigwedge^n T^*(J))$  as follows:

(1) 
$$J \ni j \to \mathfrak{L}(j) := \pi_j^*(L(j)) \epsilon \wedge^n T^*(J).$$

If X is a field on  $\tilde{u}$  of vectors tangent to J, then by  $\hat{X}$  we will denote any differentiable extension of X on the manifold J such that  $\pi_* \hat{X} = \pi_* X$ .

LEMMA 1. If 
$$\mathcal{J} \in \mathcal{F}$$
,  $\Psi \in Cu(M)$  and  $u \in \Gamma(M)$ , then

$$\langle \mathscr{J}, [\mathring{\varPsi}(u)] 
angle = \int\limits_{D_{m{u}}} \pi_{\widetilde{u}}^{m{\star}} \circ i_{m{u}}^{m{\star}} \left( \stackrel{\mathcal{L}}{\mathcal{L}} \mathfrak{L} \right), \quad where \; m{X} = [\mathring{\varPsi}].$$

Proof. We introduce the following notation:

$$\begin{split} u_t &:= \mathring{\varPsi}_t(u)\,, \\ \varphi_t(p) &:= (\pi \varPsi)_t(p), \quad X := [\mathring{\varPsi}], \quad \mathsf{w}_p := \bigwedge_{i=1}^n \left(\frac{\partial}{\partial \xi^i}\right)_p, \\ \mathsf{v}_p &:= \bigwedge_{i=1}^n d\xi^i(p), \quad L'([u]_p) := \langle \mathsf{w}_p, L([u]_p) \rangle, \\ \mathsf{V}_{[u]_p} &:= \bigwedge_{i=1}^n dx^i([u]_p), \end{split}$$

where coordinate charts are fixed.

Let  $\{I_r\}_1^k$  be such that  $D_u \subset \bigcup_{r=1}^k \stackrel{I_r}{U}$ . By  $\{\eta_r\}_1^\infty$  we denote a partition of unity subordinate to the covering  $\{\overset{\iota}{U} \cap D_u\}_1^k$ . Then

$$\mathscr{J}(u_t) = \int\limits_{D_{ttt}} \mathscr{L}_{u_t} = \sum_{\mathbf{v}} \int\limits_{D_{u_t}} \eta_{\mathbf{v}} \mathscr{L}_{u_t} = \sum_{\mathbf{v}} \int\limits_{D_{u}} (\eta_{\mathbf{v}} \mathscr{L}_{u_t}) \circ \varphi_t.$$

It should be noted that in view of the relative compactness of  $D_u$ the above sum has only a finite number of non-zero terms.

$$\begin{split} \frac{d}{dt} \mathscr{J}(u_t)\Big|_{t=0} &= \sum_{\nu} \int_{\mathcal{D}_u} \frac{d}{dt} \left( \eta_{\nu} \mathscr{L}_{u_t} \right) \circ \varphi_t \Big|_{t=0} \\ &= \sum_{\nu} \int_{\mathcal{D}_u} \left[ \mathscr{L}_u \frac{d}{dt} \eta_{\nu} \circ \varphi_t \Big|_{t=0} + \eta_{\nu} \frac{d}{dt} \mathscr{L}_{u_t} \circ \varphi_t \Big|_{t=0} \right] \\ &= \int_{\mathcal{D}_u} \sum_{\nu} \frac{d}{dt} \eta_{\nu} \circ \varphi_t \Big|_{t=0} \mathscr{L}_u + \int_{\mathcal{D}_u} \frac{d}{dt} \mathscr{L}_{ut} \circ \varphi_t \Big|_{t=0} \\ &= \int_{\mathcal{D}_u} \frac{d}{dt} \mathscr{L}_{u_t} \circ \varphi_t \Big|_{t=0}. \end{split}$$

Let  $p \in \overset{I}{U}$ : then

$$\begin{split} \frac{d}{dt}\,\mathcal{L}_{u_{\boldsymbol{\ell}}} \circ \varphi_{\boldsymbol{\ell}}(p)|_{t=0} \\ &= \frac{d}{dt}\left(\langle \mathbf{w}, \mathcal{L}_{u_{\boldsymbol{\ell}}} \rangle \mathbf{v}\right) \circ \varphi_{\boldsymbol{\ell}}(p)\Big|_{t=0} = \frac{d}{dt}\left.L'([u_t]_{\varphi_{\boldsymbol{\ell}}(\mathcal{V})}) \mathbf{v}_{\varphi_{\boldsymbol{\ell}}(\mathcal{V})}\Big|_{t=0} \\ &= \frac{d}{dt}\left.L'([u_t]_{\varphi_{\boldsymbol{\ell}}(\mathcal{V})})\Big|_{t=0} \mathbf{v}_{\mathcal{V}} + L'([u]_{\mathcal{V}}) \frac{d}{dt} \mathbf{v}_{\varphi_{\boldsymbol{\ell}}(\mathcal{V})}\Big|_{t=0} \\ &= \underbrace{\overline{\mathcal{L}}}_{\mathbf{X}} L'(p) \mathbf{v}_{\mathcal{V}} + \overline{L'}(p) \left( \underbrace{\mathcal{L}}_{\mathbf{X}} \mathbf{v})(p), \end{split}$$

where  $X_0 = \pi_* X$ 

Now we will bring the right-hand side of the thesis to the form above. We must find the form of  $i_u^*\left(\pounds\,\mathfrak{L}\right)$  and therefore we are interested in  $\mathcal{L}$   $\mathfrak{L}$  only on  $\tilde{u}$ . Thus  $\hat{x}$ 

$$\begin{split} (\overset{\pounds}{\hat{x}}\,^{\mathfrak{L}})([u]_{p}) &= (\overset{\pounds}{\hat{x}}\,L'\mathsf{V})([u]_{p}) \\ &= (\overset{\pounds}{\hat{x}}\,L')([u]_{p})\,\mathsf{V}_{[u]_{p}} + L'([u]_{p})(\overset{\pounds}{\hat{x}}\,\mathsf{V})([u]_{p}). \end{split}$$

Since  $V = \pi^* \mathbf{v}$ , we have  $\underset{\hat{\mathbf{X}}}{\mathcal{E}} V = \underset{\hat{\mathbf{X}}}{\mathcal{E}} \pi^* \mathbf{v} = \pi^* \underset{\pi_* \hat{\mathbf{X}}}{\mathcal{E}} \mathbf{v} = \pi^* \underset{\mathbf{X}_0}{\mathcal{E}} \mathbf{v}$ . Above we have made use of the identity

$$f^* \underset{f_* X}{\mathcal{L}} \omega = \underset{X}{\mathcal{L}} f^* \omega,$$

where  $\omega$  is an exterior differential form.

It is easy to see that the thesis does not depend on the choice of an extension  $\hat{X}$  of the field X defined on  $\tilde{u}$  only, q. e. d.

In  $\Gamma(M) \times Cu(M)$  we can introduce three equivalence relations:

$$\begin{split} & \left( (u, \, \varPsi) \underset{\boldsymbol{r_1}}{\sim} (v, \Phi) \right) \overset{\mathrm{df}}{\Leftrightarrow} (u = v, \, [\varPsi] \underset{\overline{\mathbb{R}_u}}{\equiv} [\varPhi]), \\ & \left( (u, \, \varPsi) \underset{\boldsymbol{r_2}}{\sim} (v, \, \varPhi) \right) \overset{\mathrm{df}}{\Leftrightarrow} (u = v, \, [\mathring{\varPsi}] \underset{\overline{u}}{\equiv} [\mathring{\varPhi}]), \\ & \left( (u, \, \varPsi) \underset{\boldsymbol{r_3}}{\sim} (v, \varPhi) \right) \overset{\mathrm{df}}{\Leftrightarrow} (u = v, \, [\mathring{\varPsi}(u)] = [\mathring{\varPhi}(u)]). \end{split}$$

LEMMA 2.  $r_1 \Leftrightarrow r_2 \Rightarrow r_3$ .

Proof. Let  $\Psi \epsilon Cu(M)$ . If the field  $|\mathring{\Psi}|$  has the form

$$[\overset{*}{ extbf{Y}}] = \overset{I}{ extbf{X}} rac{\partial}{\partial x^i} + \overset{I}{ extbf{X}^0} rac{\partial}{\partial y} + \overset{I}{ extbf{Y}^i} rac{\partial}{\partial y^i}$$

at points of the set  $\pi^{-1}(\stackrel{.}{U}) \cap \tilde{u}$ , then the field  $[\Psi]$  has the form

$$[\Psi] = \overset{I}{X}{}^{i} \circ f \frac{\partial}{\partial (\overset{I}{\xi}{}^{i} \circ \pi)} + \overset{I}{X}{}^{0} \circ f \frac{\partial}{\partial \overset{I}{\rho}}$$

at points of the set  $\pi^{-1}(\stackrel{I}{U}) \cap R_u$ , where f is the canonical 1-1 map of  $R_u$  on  $\tilde{u}$  given by

$$R_u \circ u(p) \to f(u(p)) := [u]_p \in \tilde{u}.$$

Hence  $r_1 \Rightarrow r_2$ . For the proof of the implication  $r_2 \Rightarrow r_1$  see Appendix 1.

It follows from Lemma 1 that  $r_2 \Rightarrow r_3$ , q. e. d.

Remark. Let us notice that in general the implication  $r_3 \Rightarrow r_1$  does not occur. For example, a curve  $\mathring{\varPsi}(u)$  in  $\varGamma(M)$  such that

$$\partial R_{\mathcal{P}_{t}(u)} \equiv \partial R_{u}$$

can be given by a field  $\Psi \epsilon Cu(M)$  that curves  $\Psi(m)$ ,  $m \epsilon M$ , lie inside of fibres as well as by a field  $\Phi \epsilon Cu(M)$  that for some  $m \epsilon M$  curves  $\Phi(m)$  are transversal to fibres.

Now we can pass from dealing with fields of curves to dealing with vector fields. We define the following vector spaces:

$$\begin{split} T_M := & \{X \epsilon \Gamma \big(M, T(M)\big) \colon \pi(m) = \pi(m') \Rightarrow \pi_* X_m = \pi_* X_{m'}\}, \\ W_M(u) := & \{X \epsilon \Gamma \big(R_u, T(M)\big) \colon \bigwedge_{\substack{X \ W \epsilon Ou(M)}} X = [\varPsi]|_{R_u}\}, \\ W_J(u) := & \{X \epsilon \Gamma \big(\tilde{u}, T(J)\big) \colon \bigwedge_{\substack{X \ W \epsilon Cu(M)}} X = [\mathring{\varPsi}]\}. \end{split}$$

Now we can formulate Lemma 2 in a new form:

LEMMA 2'. There are canonical homomorphisms onto  $-h_u$ , H and the canonical isomorphism I:

$$T_M \xrightarrow{h_U} W_M(u) \xrightarrow{I} W_J(u) \xrightarrow{H} \mathfrak{W}_u(\Gamma(M)).$$

Proof. It is sufficient to give the explicite form of these maps:

$$\begin{split} T_M &\ni X \to h_u(X) \, := \, X \, |_{R_u} \epsilon \, W_M(u) \,, \\ W_M(u) &\ni [\Psi] \, |_{R_u} \to I([\Psi] \, |_{R_u}) \, := \, [\stackrel{*}{\varPsi}] \epsilon \, W_J(u) \,, \\ W_J(u) &\ni [\stackrel{*}{\varPsi}] \to H([\stackrel{*}{\varPsi}]) \, := \, [\stackrel{\circ}{\varPsi}(u)] \, \epsilon \, \mathfrak{V}_u \big( \varGamma(M) \big) \,. \end{split}$$

The correctness of above definitions follows from Lemma 1 and Lemma 2, q. e. d.

Now we formulate Lemma 1 in terms of vector fields.

LEMMA 1'. Let  $\mathcal{J} \in \mathcal{F}$ ,  $\mathfrak{X}_u \in \mathfrak{W}_u(\Gamma(M))$ ; then

$$\langle \mathscr{J}, \mathfrak{X}_u 
angle = \int\limits_{D_{m{u}}} \pi_{m{u}}^{m{\star}} \circ i_{m{u}}^{m{\star}} \left( \pounds \, \mathfrak{L} \, \mathfrak{L} \right), \quad ext{where} \quad m{X} \, \epsilon \, H^{-1}(\mathfrak{X}_u).$$

There is a unique decomposition  $X = X_{\perp} + (i_{\widetilde{u}} \circ \pi_{\widetilde{u}})_* X_0$  for  $X \in W_J(u)$ , where  $X_0 \in \Gamma(D_u, T(E))$  and  $(\pi)_* X_{\perp} = 0$ . Using (2) we have

$$\begin{split} \pi_{u}^{*} \circ i_{u}^{*} & (\pounds \, \mathfrak{L}) = \pi_{u}^{*} \circ i_{u}^{*} \, (\pounds \, \mathfrak{L}) + \pounds \, \pi_{u}^{*} \circ i_{u}^{*} \, \mathfrak{L} \\ \mathbf{x}_{\perp} & \mathbf{x}_{0} & \mathbf{x}_{u}^{*} \circ i_{u}^{*} \, \mathfrak{L} \end{split}$$
$$= \pi_{u}^{*} \circ i_{u}^{*} \, (\pounds \, \mathfrak{L}) + d\mathbf{X}_{0} \, \mathbf{x}_{u}^{*} \circ i_{u}^{*} \, \mathfrak{L}$$

The last expression follows from the identity  $\pounds \omega = X$  d $\omega + dX$  and from the fact that  $d\pi_u^* \circ i_u^* \mathfrak{L} = 0$ .



Let  $\{\eta_r\}_1^{\infty}$  be as before a partition of unity subordinate to a covering  $\{U \cap D_u\}_1^k$  of  $D_u$ . We define the following (n-1)-form on E:

$$N_{\boldsymbol{X}_{\perp}}(\mathscr{L}_{u}) := \sum_{\boldsymbol{r},k} (-1)^{k} \eta_{\boldsymbol{r}} \boldsymbol{X}_{\perp}^{1} \frac{\overline{I_{\boldsymbol{r}}}}{\partial L} d_{\xi^{1}}^{1_{\boldsymbol{r}}} \wedge \overset{k}{\wedge} \wedge \frac{I_{\boldsymbol{r}}}{\partial \xi^{k}},$$

where  $L, \nu = 1, 2, ...$ , are the coefficients in the expression of the form L in the coordinate system connected with the domain U.

For the proof of correctness of this definition see Appendix.

By an Euler-Lagrange's derivative of the *n*-form  $\mathscr{L}_u$  with respect to  $X \in W_M(u)$  we mean

$$[\mathscr{L}_u]_X := \pi_u^{\!\!\!\star} \circ i_u^{\!\!\!\star} (\mathscr{L}\, \mathfrak{L}) - dN_{X_\perp}(\mathscr{L}_u), \quad \text{ where } \quad X = I(X).$$

We define Noether's expression for the  $n\text{-form }\mathcal{L}_u$  and  $X \in W_M(u)$  as follows:

$$\mathscr{N}_X(\mathscr{L}_u) := N_{\pmb{X}_\perp}(\mathscr{L}_u) + \pmb{X}_0 \, \underline{\,\,\,\,} \pi_{\tilde{u}}^{\pmb{\star}} \circ i_{\tilde{u}}^{\pmb{\star}} \, \mathfrak{L},$$

where X = I(X).

Hence if  $X \in W_M(u)$  and  $\mathfrak{X}_u = H \circ I(X)$ , then

(4) 
$$\langle \mathcal{J}, \mathfrak{X}_u \rangle = \int\limits_{D_u} [\mathcal{L}_u]_{\mathbb{X}} + d\mathcal{N}_X(\mathcal{L}_u).$$

We define vector subspaces:

$$\begin{split} D_M(u) &:= \{X \in W_M(u) \colon X \underset{\partial \mathcal{R}_u}{==} 0\}, \\ D_J(u) &:= I \big( D_M(u) \big), \quad \mathfrak{D}_u \big( \varGamma(M) \big) := H \big( D_J(u) \big) \end{split}$$

A cross-section  $u \in \Gamma(M)$  is called *extremal* for  $\mathscr{J} \in \mathscr{F}$  if and only if, for every  $\mathfrak{A}_u \in \mathfrak{D}_u(\Gamma(M))$ ,  $\langle \mathscr{I}, \mathfrak{A}_u \rangle = 0$ . By  $\mathscr{S}(\mathscr{I})$  we denote the set of all  $u \in \Gamma(M)$  extremal for a given  $\mathscr{I} \in \mathscr{F}$ .

LEMMA 3. We have

$$\big(u\,\epsilon\,\mathscr{S}(\mathscr{J})\big)\Leftrightarrow \big(\underset{X\,\epsilon\,\overline{\mathscr{W}}_M(u)}{\wedge}\,[\,\mathscr{L}_u]_X \underset{\overline{D}_u}{\equiv}\,0\big).$$

Proof. If  $X \in D_M(u)$ , then  $\mathscr{N}_X(\mathscr{L}_u) \equiv 0$ . Hence for  $X \in D_M(u)$  we have  $\langle \mathscr{J}, H \circ I(X) \rangle = \int\limits_{D_u} [\mathscr{L}_u]_X$  and so the implication  $\Leftarrow$  is proved.

Let  $X \in W_M(u)$  and X = I(X). It follows directly from the definitions that for  $p \in \stackrel{\mathcal{I}}{U} \cap D_u$ 

$$[\mathcal{L}_{u}]_{X}|_{y} = \overline{X}_{\perp}^{0} \left( \frac{\overline{\partial L}}{\partial v} - \frac{\partial}{\partial \xi^{i}} \frac{\overline{\partial L}}{\partial v_{i}} \right)_{\mathbf{v}}^{I}|_{y};$$

$$[\mathcal{L}_{u}]_{X} = \sum_{\mathbf{v}} \eta_{\mathbf{v}} \overline{X}_{\perp}^{0} \left( \frac{\overline{\partial L}}{\partial v} - \frac{\partial}{\partial \xi^{i}} \frac{\overline{\partial L}}{\partial v_{i}} \right)_{\mathbf{v}}^{I}|_{\mathbf{v}}^{I}$$

and, which is of the highest importance, that  $[\mathscr{L}_u]_X(p)$  depends on  $X_p$  only. We will prove that if  $u \in \mathscr{S}(\mathscr{J})$ , then for every U

(6) 
$$A := \frac{\overline{\frac{I}{\partial L}}}{\frac{I}{\partial v}} - \frac{\partial}{\partial \xi^i} \frac{\overline{\frac{I}{\partial L}}}{\frac{I}{\partial v_i}} \equiv 0 \quad \text{ on } \stackrel{I}{U} \cap D_u.$$

Let us suppose that there exists a point  $p \in U \cap D_u$  such that  $A(p) \neq 0$ ; then there exists a neighborhood  $\Omega$  of p that  $\Omega \subset U \cap D_u$  and  $A(q) \neq 0$  for  $q \in \Omega$ . Now let  $X \in D_M(u)$  be such that  $X_{\perp} = 0$  on  $C_{\pi_u}(\Omega)$  and  $X_{\perp}^0 \geq 0$ . Then  $X_{\perp}^0 A$  has a constant sign on  $\Omega$ . Hence  $\int_{D_u} [\mathcal{L}_u]_X \neq 0$  (a contradiction). Now from (4) and (5) the implication  $\Rightarrow$  follows, q. e. d.

For every  $\mathcal{J} \in \mathcal{F}$  and for every open  $\Omega \subset E$  we define the functional  $\mathcal{J}_{\Omega} \in \mathcal{F}$  in the following way:

$$\Gamma(M) \circ u \to \mathscr{J}_{\Omega}(u) := \mathscr{J}(u|_{\Omega_{\Omega}D_{u}}) \epsilon \mathbf{R}^{1}$$

 $\mathscr{J} \in \mathscr{F}$  is called invariant at a point  $u \in \varGamma(M)$  with respect to  $X \in T_M$  if and only if

$$\langle \mathcal{J}, H \circ I \circ h_{u}(X) \rangle = 0$$
.

THE FIRST NOETHER'S THEOREM. Let  $\mathcal{J} \in \mathcal{F}$ ,  $u \in \mathcal{S}(\mathcal{J})$ ,  $X \in T_M$ ; then

The proof is an immediate consequence of (4) and Lemma 3. It is easy to see that for  $X, Y \in W_M(u), a, b \in \mathbb{R}^1$ 

$$\mathcal{N}_{aX+bY}(\mathcal{L}_u) = a\mathcal{N}_X(\mathcal{L}_u) + b\mathcal{N}_Y(\mathcal{L}_u).$$

Now we can state the second variant of

THE FIRST NOETHER'S THEOREM. Let  $\mathcal{J} \in \mathcal{F}$ ,  $u \in \mathcal{S}(\mathcal{J})$ ,  $\{X_i\}_1^r$  be a base of a subspace  $T \subset T_M$ ; then

$$\left(\begin{array}{c} \textit{For every } \Omega \subset D_u \textit{ and every} \\ X \in T, \mathscr{J}_{\varOmega} \textit{ is invariant at} \\ \textit{the point u with respect to } X \end{array}\right) \Leftrightarrow \left(d\mathscr{N}_{h_{\mathcal{U}}\!\left(X_i\right)}\!\left(\mathscr{L}_{u}\right) \underset{D_{u}}{\equiv} 0\,,\,i=1,\ldots,r\right).$$

Remark. The assumption  $D_u$  as relatively compact is not weighty. If we introduce a new definition that a cross-section of M is called *extremal* for  $\mathscr{J} \in \mathscr{F}$  if and only if

$$\bigwedge_{K \subset E} (K\text{-compact}) \Rightarrow (u|_K \in \mathcal{S}(\mathcal{J})),$$

hen it is clear that for  $\boldsymbol{u}$  extremal in this sense The First Noether's Theotem is valid.

### APPENDIX

1. Let  $X \in T(J)$ ; then in the base induced by the coordinate chart  $(\chi, \pi^{-1}(\overset{I}{U}))$  the vector  $X_{\perp}$  has the form

$$m{X}_{ot} = m{X}_{ot}^0 rac{\partial}{\partial v} + m{Y}_{ot}^I rac{\partial}{\partial v_i}.$$

But if  $X \in W_J(u)$  and  $X = [\mathring{\Psi}]$ , then

$$\begin{split} \overset{I}{Y}_{\perp}^{i}([u]_{p}) &= \frac{d}{dt} \overset{I}{v_{t}} \circ \overset{*}{\mathcal{Y}_{t}}([u]_{p})|_{t=0} = \frac{d}{dt} \left( \frac{\partial \overset{I}{\varrho} \circ \overset{*}{\mathcal{Y}_{t}}(u)}{\partial \overset{I}{\xi^{i}}}(p) \right) \Big|_{t=0} \\ &= \frac{\partial}{\partial \overset{I}{\xi^{i}}} \frac{d\overset{I}{\varrho} \circ \overset{\circ}{\mathcal{Y}_{t}}(u)(p)}{dt} \Big|_{t=0} = \frac{\partial \overset{I}{X}_{\perp}^{0}}{\partial \overset{I}{\xi^{i}}}([u]_{p}). \end{split}$$

2. The following identities are valid on  $\pi^{-1}(\overset{I}{U} \cap \overset{J}{U}) \subset J$ :

$$egin{aligned} \overset{J}{v} &= A_{JI} \circ \pi \overset{I}{v}, \ \overset{J}{v_k} &= \dfrac{\partial A_{JI}}{\partial \overset{J}{\xi^k}} \circ \pi \overset{I}{v} + A_{JI} \circ \pi & \dfrac{\partial \overset{I}{\xi^i}}{\partial \overset{J}{\xi^k}} \circ \pi \overset{I}{v_i}, \ \dfrac{\partial}{\partial v} &= A_{JI} \circ \pi & \dfrac{\partial}{\partial v} + \dfrac{\partial A_{JI}}{\partial \overset{J}{\xi^k}} \circ \pi & \dfrac{\partial}{\partial v_i}, \end{aligned}$$

$$egin{aligned} rac{\partial}{\partial v_i} &= A_{JI} \circ \pi \; rac{\partial ar{\xi}^i}{\partial ar{\xi}^k} \circ \pi \; rac{\partial}{\partial v_k}, \ m{X}^i &= m{X}^k \; rac{\partial ar{\xi}^i}{\partial ar{\xi}^k} \circ \pi, \ m{X}^0 &= m{X}^0 A_{IJ} \circ \pi + m{X}^i \; rac{\partial A_{IJ}}{\partial ar{\xi}^i} \circ \pi^J, \ m{L} &= \left| rac{D(arkappa)}{D(arkappa)} 
ight| \circ \pi^J. \end{aligned}$$

3. The above identities are used in the following proof of correctness of the definition (3). We define

$$I\{\mathscr{L}_u\}_{m{X}_\perp} := \sum_k \left(-1
ight)^k \overline{m{X}_\perp^0} \, \overline{rac{\partial L}{i}} \, d_{\xi^1}^{m{I}} \wedge \stackrel{k}{ \bigcap\limits_{j=1}^k} \wedge d_{\xi^n}^{m{I}}.$$

Taking into account that  $\overset{I}{X_{\perp}^{i}}=0,\,i=1,\ldots,n,$  we will prove that

$$(7) I\{\mathscr{L}_{u}\}_{\boldsymbol{X}_{\perp}} = J\{\mathscr{L}_{u}\}_{\boldsymbol{X}_{\perp}} on \stackrel{\boldsymbol{I}}{\boldsymbol{U}} \cap \stackrel{\boldsymbol{J}}{\boldsymbol{U}} \cap D_{u}.$$

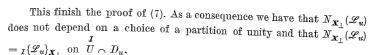
$$I\{\mathscr{L}_{u}\}_{\boldsymbol{X}_{\perp}} = \sum_{k} (-1)^{k} \stackrel{\boldsymbol{X}_{\perp}}{\boldsymbol{X}_{\perp}} A_{IJ} A_{JI} \frac{\partial \xi^{k}}{\partial \xi^{i}} \left( \frac{\partial}{\partial v^{i}} \left| \frac{D(\varkappa)}{D(\varkappa)} \right| \circ \pi \stackrel{\boldsymbol{J}}{\boldsymbol{U}} \right) d\xi^{I} \wedge \stackrel{\boldsymbol{\lambda}}{\boldsymbol{\dots}} \wedge d\xi^{n}$$

$$= \sum_{k,l} (-1)^{k} \stackrel{\boldsymbol{J}}{\boldsymbol{X}_{\perp}} \frac{\partial \xi^{k}}{\partial \xi^{i}} \left| \frac{D(\varkappa)}{D(\varkappa)} \right| \frac{\partial \boldsymbol{L}}{\partial v_{i}} \left| \frac{D(\xi^{1}, \dots, \xi^{n})}{D(\xi^{1}, \dots, \xi^{n})} \right| d\xi^{I} \wedge \stackrel{\boldsymbol{\lambda}}{\boldsymbol{\dots}} \wedge d\xi^{n}.$$

$$A_{lk} := (-1)^{k+l} \left| \frac{D(\xi^{1}, \dots, \xi^{n})}{D(\xi^{1}, \dots, \xi^{n})} \right| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \right| \right| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \right| \right| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \right| \right| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \right| \right| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \right| \right| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \right| \right| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \right| \right| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \right| \right| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \right| \right| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \right| \right| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \right| \right| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \right| \right| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \right| \right| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \right| \left| \frac{\partial}{\partial v_{i}} \right| \left| \frac{\partial}{\partial v_{i}} \right| \right| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \right| \left| \frac{\partial}{\partial v_{i}} \right| \right| \frac{\partial}{\partial v_{i}} \left| \frac{\partial}{\partial v_{i}} \right| \right| \left| \frac{\partial}{\partial v_{i}} \right| \left| \frac{\partial}{\partial v_{$$

is the algebraic complement of the element  $a_{lk}:=\partial\xi^k/\partial\xi^l$  in the matrix  $A:=D(\varkappa)/D(\varkappa)$ . The following identity is valid:

$$\sum_{k} a_{ik} A_{ik} = \delta_{ii} \det A.$$



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Reçu par la Rédaction le 7. 6. 1967