

Thermodynamic formalism methods in one-dimensional real and complex dynamics

Feliks Przytycki

Institute of Mathematics of the Polish Academy of Sciences

ICM 2018, Rio de Janeiro

1 – Introduction

- Pioneer of statistical physics: Marian Smoluchowski 1872 – 1917 Vienna-Lvov-Cracow.
- Application of thermodynamic methods to dynamics: Yakov Sinai, David Ruelle, Rufus Bowen 1960/70 -ties.

Lemma (finite variational principle)

For given real numbers ϕ_1, \dots, ϕ_d , the function

$$F(p_1, \dots, p_d) := \underbrace{\sum_{i=1}^d -p_i \log p_i}_{\text{entropy}} + \underbrace{\sum_{i=1}^d p_i \phi_i}_{\text{average potential}}$$

on the simplex $\{(p_1, \dots, p_d) : p_i \geq 0, \sum_{i=1}^d p_i = 1\}$ attains its maximum, called **pressure** equal to $P(\phi) = \log \sum_{i=1}^d e^{\phi_i}$, at the **equilibrium**

$$\hat{p}_j = e^{\phi_j} / \sum_{i=1}^d e^{\phi_i}.$$

Hint: $\sum_{i=1}^d -p_i \log p_i + \sum_{i=1}^d p_i \phi_i = \sum_{i=1}^d p_i \log(e^{\phi_i} / p_i)$.

1 – Introduction: dynamics setting corresponding notions

$f : X \rightarrow X$ a contin. map for a compact metric space (X, ρ) ,
 $\phi : X \rightarrow \mathbb{R}$ a continuous function (*potential*).

Definition (variational topological pressure)

$$P_{\text{var}}(f, \phi) := \sup_{\mu \in \mathcal{M}(f)} \left(h_{\mu}(f) + \int_X \phi d\mu \right),$$

where $\mathcal{M}(f)$ is the set of all f -invariant Borel probability measures on X and $h_{\mu}(f)$ is measure-theoretical entropy.

Any measure where sup is attained is called *equilibrium state*.

Definition (topological pressure via separated sets)

$P_{\text{sep}}(f, \phi) := \lim_{\varepsilon \rightarrow 0} \overline{\lim}_{n \rightarrow \infty} \frac{1}{n} \log \left(\sup_Y \sum_{y \in Y} \exp S_n \phi(y) \right)$,
supremum over all $Y \subset X$ such that for distinct $x, y \in Y$,
 $\rho_n(x, y) := \max\{\rho(f^i(x), f^i(y)), 0 \leq i \leq n\} \geq \varepsilon$.

- $h_\mu(f) := \sup_{\mathcal{A}} \lim_{n \rightarrow \infty} \frac{1}{n+1} \sum_{A \in \mathcal{A}^n} -\mu(A) \log \mu(A)$,
 supremum over finite partitions \mathcal{A} of X ,
 $\mathcal{A}^n := \bigvee_{j=0, \dots, n} f^{-j} \mathcal{A}$.

Theorem (variational principle: Ruelle, Walters, Misiurewicz, Denker, ...)

$$P_{\text{var}}(f, \phi) = P_{\text{sep}}(f, \phi).$$

FP & M. Urbański "Conformal Fractals: Ergodic Theory Methods" Cambridge 2010.

Theorem (Gibbs measure – uniform case)

Let $f : X \rightarrow X$ be a **distance expanding, topologically transitive continuous open map** of a compact metric space X and $\phi : X \rightarrow \mathbb{R}$ be a **Hölder continuous potential**. Then, there exists exactly one $\mu_\phi \in \mathcal{M}(f, X)$, called **Gibbs measure**, s.t.

$$C < \frac{\mu_\phi(f_x^{-n}(B(f^n(x), r_0)))}{\exp(S_n\phi(x) - nP)} < C^{-1},$$

called **Gibbs property**, where f_x^{-n} is the local branch of f^{-n} mapping $f^n(x)$ to x and $S_n\phi(x) := \sum_{j=0}^{n-1} \phi(f^j(x))$.

- μ_ϕ is the **unique equilibrium state** for ϕ . It is equivalent to the unique **$\exp -(\phi - P)$ -conformal measure** m_ϕ , that is an f -quasi-invariant measure with Jacobian $\exp -(\phi - P)$ for a constant P .
- $P = P(f, \phi) := \lim_{n \rightarrow \infty} \frac{1}{n} \log \sum_{x \in f^{-n}(x_0)} \exp S_n\phi(x)$. This normalizing limit exists and is equal $P_{\text{sep}}(f, \phi)$ for every $x \in X$.

2 – Introduction to dimension 1

Thermodynamic formalism is useful for studying properties of the underlying space X . In dimension 1, for f real of class $C^{1+\varepsilon}$ or f holomorphic (conformal) for an expanding repeller X , considering $\phi = \phi_t := -t \log |f'|$ for $t \in \mathbb{R}$, Gibbs property gives, as $\exp S_n(\phi_t) = |(f^n)'|^{-t}$,

$$\mu_{\phi_t}(f_x^{-n}(B(f^n(x), r_0))) \approx \exp(S_n\phi(x) - nP(\phi_t)) \approx \text{diam } f_x^{-n}(B(f^n(x), r_0))^t \exp -nP(\phi_t).$$

The latter follows from a comparison of the diameter with the inverse of the absolute value of the derivative of f^n at x , due to *bounded distortion*.

When $t = t_0$ is a zero of the function $t \mapsto P(\phi_t)$, this gives

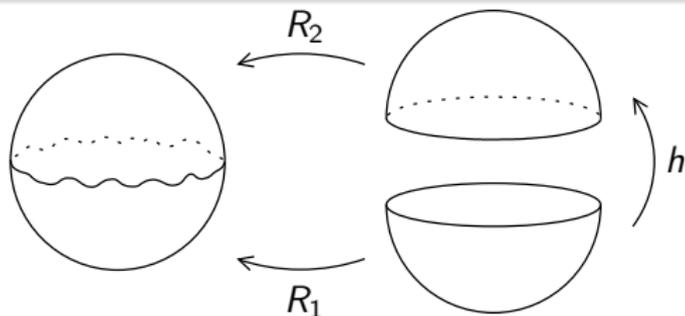
$$\mu_{\phi_{t_0}}(B) \approx (\text{diam } B)^{t_0}$$

for all small balls B , hence $\text{HD}(X) = t_0$. Moreover, the Hausdorff measure of X in this dimension is finite and nonzero.

A model application

Theorem (Bowen, Series, Sullivan)

For $f_c(z) := z^2 + c$ for an arbitrary complex number $c \neq 0$ sufficiently close to 0, the invariant Jordan curve J (Julia set for f_c) is fractal, i.e. has Hausdorff dimension bigger than 1.



If $\text{HD}(J) = 1$, then $0 < H_1(J) < \infty$ and $h = R_2^{-1} \circ R_1$ on S^1 is absolutely continuous. $g_i := R_i^{-1} \circ f_c \circ R_i$ for $i = 1, 2$ preserve length ℓ on S^1 and are ergodic. Hence h preserves ℓ so it is a rotation, identity for appropriate R_1, R_2 . Hence R_1 and R_2 glue together to a homography. Compare Mostov rigidity theorem.

complex case

In the complex case we consider f a rational mapping of degree at least 2 of the Riemann sphere $\overline{\mathbb{C}}$. We consider f acting on its **Julia set** $K = J(f)$ (generalizing the $z^2 + c$ model).

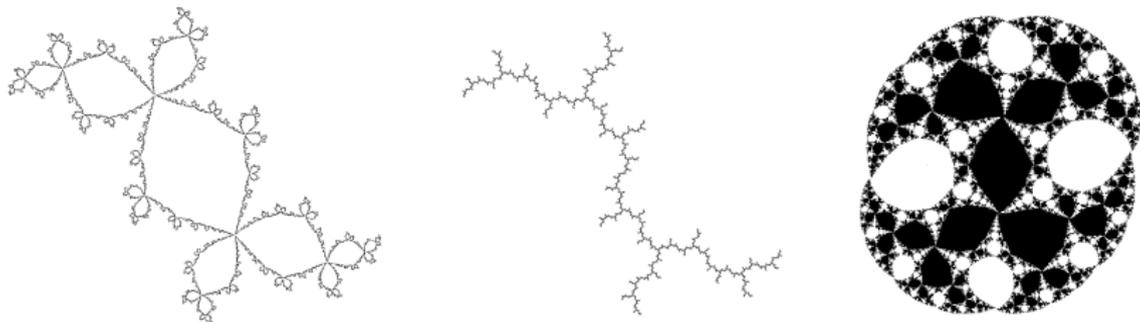


Figure: Douady's zoo: **rabbit** $f(z) = z^2 - 0.123 + 0.745i$, **dendrite** $f(z) = z^2 + i$, **basilica mated with rabbit** $f(z) = \frac{z^2+c}{z^2-1}$ for $c = \frac{1+\sqrt{-3}}{2}$ with $J(f)$ being the boundary of black and white

Definition (Real case, FP & Rivera-Letelier)

$f \in C^2$ is called a *generalized multimodal map* if defined on a neighbourhood of a compact invariant set K , critical points are not infinitely flat, *bounded distortion* property for iterates holds, abbr. BD, f is *topologically transitive* and has *positive topological entropy* on K .

Also K is a *maximal* forward invariant subset of a finite union of pairwise disjoint closed intervals whose endpoints are in K .

This maximality corresponds to Darboux property. We write $(f, K) \in \mathcal{A}_+^{\text{BD}}$, where $+$ marks positive entropy. In place of BD one can assume C^3 (and write $(f, K) \in \mathcal{A}_+^3$) and assume that all periodic orbits in K are *hyperbolic repelling*. Then changing f outside K allows to get $(f, K) \in \mathcal{A}_+^{\text{BD}}$.

Examples: Basic sets in spectral decomposition via renormalizations (de Melo, van Strien).

3 – Hyperbolic potentials

Call $\phi : K \rightarrow \mathbb{R}$ satisfying $P(f, \phi) > \sup_{\nu \in \mathcal{M}(f)} \int \phi d\nu$
hyperbolic potential (Inoquio-Renteria, Rivera-Letelier: BBMS 2012). Equiv. $P(f, \phi) > \sup_K \frac{1}{n} S_n \phi$ for some n .

Theorem (complex and real: Denker, Urbański, FP, Haydn, Rivera-Letelier, Zdunik, Szostakiewicz, H. Li, Bruin, Todd)

. If ϕ is a Hölder continuous *hyperbolic* potential, then there exists a unique equilibrium state μ_ϕ . For every Hölder $u : K \rightarrow \mathbb{R}$, the Central Limit Theorem (CLT) and Law of Iterated Logarithm (LIL) for the sequence of random variables $u \circ f^n$ and μ_ϕ hold.

CLT follows from sufficiently fast convergence of iteration of transfer operator (spectral gap). LIL is proved via LIL for a return map (inducing) to a nice domain related to μ_ϕ (Mañé, Denker, Urbański) providing a Markov structure (Infinite Iterated Function System) avoiding critical points, satisf. BD.

4 – Non-uniform hyperbolicity

a) CE. *Collet-Eckmann condition*. There exists $\lambda > 1$, $C > 0$

$$|(f^n)'(f(c))| \geq C\lambda^n.$$

for all critical points $c \in K$ whose forward orbit is disjoint from $\text{Crit}(f)$. Moreover there are no indifferent periodic orbits in K .

(b) CE2(z_0). *Backward* or *second Collet-Eckmann condition at $z_0 \in K$* . There exist $\lambda > 1$ and $C > 0$ such that for every $n \geq 1$ and every $w \in f^{-n}(z_0)$ (in a neighbourhood of K in the real case)

$$|(f^n)'(w)| \geq C\lambda^n.$$

(c) TCE. *Topological Collet-Eckmann condition* (FP & S. Rohde, Fund. Math. 1998).

There exist $M \geq 0, P \geq 1, r > 0$ such that for every $x \in K$ there exist increasing $n_j, j = 1, 2, \dots$, such that $n_j \leq P \cdot j$ and for each j and discs $B(\cdot)$ below understood in $\overline{\mathbb{C}}$ or \mathbb{R} .

$$\#\{0 \leq i < n_j : \text{Comp}_{f^i(x)} f^{-(n_j-i)} B(f^{n_j}(x), r) \cap \text{Crit}(f) \neq \emptyset\} \leq M.$$

Each component of $f^{-n}(B)$ is called a **pullback** of B .

(d) **ExpShrink.** *Exponential shrinking of components.* There exist $\lambda > 1$ and $r > 0$ such that for every $x \in K$, every $n > 0$ and every connected component W_n of $f^{-n}(B(x, r))$ for the disc (interval) $B(x, r)$ in $\overline{\mathbb{C}}$ (or \mathbb{R}), intersecting K

$$\text{diam}(W_n) \leq \lambda^{-n}.$$

(e) **LyapHyp.** *Lyapunov hyperbolicity.* There is $\lambda > 1$ such that the Lyapunov exponent $\chi(\mu) := \int_K \log |f'| d\mu$ of any ergodic measure $\mu \in \mathcal{M}(f, K)$ satisfies $\chi(\mu) \geq \log \lambda$.

(f) **UHP.** *uniform hyperbolicity on periodic orbits.* There exists $\lambda > 1$ such that every periodic point $p \in K$ of period $k \geq 1$ satisfies

$$|(f^k)'(p)| \geq \lambda^k.$$

Theorem (... , Keller, Nowicki, Sands, FP, Rohde, Rivera-Letelier, Graczyk, Smirnov)

Assume there are no indifferent periodic orbits in K . Then

- 1. The conditions (c)–(f) and else (b) for some z_0 are equivalent (in the real case under the assumption of weak isolation: any periodic orbit close to K must be in K).*
- 2. CE implies (b)–(f).*
- 3. If there is only one critical point in the Julia set in the complex case or if f is S -unimodal on $K = I$ in the real case, then all conditions above are equivalent to each other.*
- 4. TCE is topologically invariant; therefore all other conditions equivalent to it are topologically invariant.*

For polynomials (b)–(f) are equivalent to $K = J(f) = \text{Fr } \Omega_\infty(f)$, the basin of ∞ , being Hölder (Graczyk, Smirnov).

An order of proving the equivalences in Theorem above is, for z_0 safe,

$$\text{CE2}(z_0) \Rightarrow \text{ExpShrink} \Rightarrow \text{LyapHyp} \Rightarrow \text{UHP} \Rightarrow \text{CE2}(z_0)$$

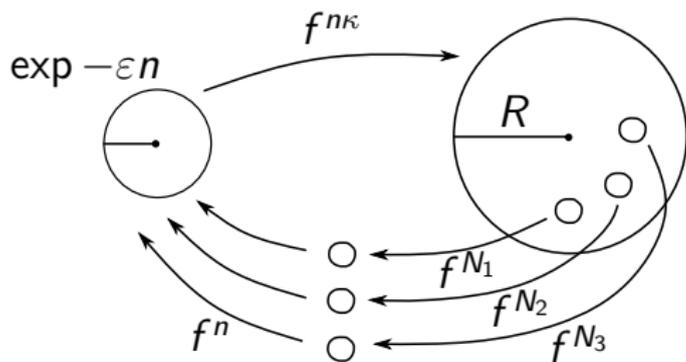
Separately one proves $\text{ExpShrink} \Leftrightarrow \text{TCE}$ using for \Rightarrow the following

Lemma (Denker, FP, Urbański, ETDS 1996)

$$\sum_{j=0}^n{}' -\log |f^j(x) - c| \leq Qn$$

for a constant $Q > 0$ every $c \in \text{Crit}(f)$, every $x \in K$ and every integer $n > 0$. Σ' means that we omit in the sum an index j of smallest distance $|f^j(x) - c|$.

Assumed UHP one proves $CE2(z_0)$ for **safe** and **hyperbolic** z_0 by “shadowing”.



Definition (safe)

We call $z \in K$ *safe* if $z \notin \bigcup_{j=1}^{\infty} (f^j(\text{Crit}(f)))$ and for every $\epsilon > 0$ and all n large enough

$$B(z, \exp(-\epsilon n)) \cap \bigcup_{j=1}^n (f^j(\text{Crit}(f))) = \emptyset.$$

Notice that this definition implies that all points except at most a set of Hausdorff dimension 0, are safe.

5 – Geometric variational pressure and equilibrium states

For $\phi = \phi_t := -t \log |f'|$, the variational definition of pressure, here

$$P(t) := P_{\text{var}}(f, \phi_t) = \sup_{\mu \in \mathcal{M}(f)} \left(h_{\mu}(f) - t \int_K \log |f'| d\mu \right)$$

still makes sense by the integrability of $\log |f'|$. Moreover

$$\int_K \log |f'| d\mu = \chi(\mu) \geq 0,$$

for all ergodic μ even in presence of critical points where $\phi = \pm\infty$, [FP: PAMS 1993, Rivera-Letelier: arXiv 2012]. By this definition $t \mapsto P(t)$ is convex, monotone decreasing.

We usually assume $t > 0$ later on.

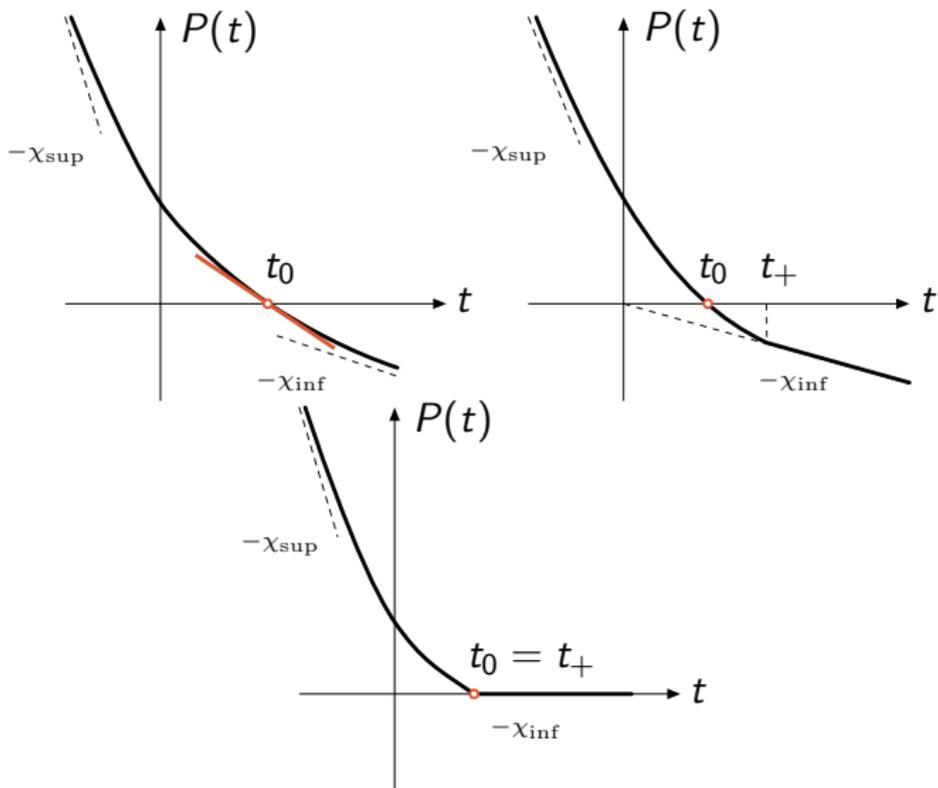


Figure: The geometric pressure: LyapHyp with $t_+ = \infty$, LyapHyp with $t_+ < \infty$, and non-LyapHyp.

$P(t)$ is equal to several other quantities (Complex: FP TAMS 1999, FP & Rivera-Letelier & Smirnov ETDS 2004). E.g.

Definition (hyperbolic pressure)

$$P_{\text{hyp}}(t) := \sup_{X \in \mathcal{H}(f, K)} P(f|_X, -t \log |f'|),$$

where $\mathcal{H}(f, K)$ is defined as the space of all compact forward inv., i.e. $f(X) \subset X$, expanding subsets of K , repellers in \mathbb{R} .

Definition (hyperbolic dimension)

$$\text{HD}_{\text{hyp}}(K) := \sup_{X \in \mathcal{H}(f, K)} \text{HD}(X).$$

For expanding $f : X \rightarrow X$, $t_0(X) = \text{HD}(X)$. Passing to sup:

Proposition (Generalized Bowen's formula)

The first zero t_0 of $t \mapsto P_{\text{hyp}}(K, t)$ is equal to $\text{HD}_{\text{hyp}}(K)$.

It may happen $\text{HD}_{\text{hyp}}(J(f)) < \text{HD}(J(f)) = 2$ for f quadratic polynomials, Avila & Lyubich.

Theorem (FP & Rivera-Letelier)

1. Real case (arXiv 2014, to appear in Memoir of the AMS).

Let $(f, K) \in \mathcal{A}_+^3$, f -periodic orbits in K be hyperbolic repelling. Then

- $t \mapsto P(t)$ is **real analytic** on an open interval (t_-, t_+) defined by $P(t) > \sup_{\nu \in \mathcal{M}(f)} -t \int \log |f'| d\nu$
- For each t in this interval **there is a unique invariant equilibrium state** μ_{ϕ_t} . It is ergodic and absolutely continuous with respect to an adequate **conformal measure** m_{ϕ_t} with $d\mu_{\phi_t}/dm_{\phi_t} \geq \text{Const} > 0$ a.e.
- If furthermore f is topologically exact on K (that is for every V an open subset of K there exists $n \geq 0$ such that $f^n(V) = K$), then this measure is **mixing, has expon. decay of corr. and satisfies CLT for Lipschitz gauge functions.**

This generalizes results by Bruin, Iommi, Pesin, Senti, Todd.

Theorem (FP & Rivera-Letelier)

2. *Complex case (Comm. Mat. Phys. 2011). The assertion is the same. One assumes a very weak expansion: **the existence of arbitrarily small nice, or pleasant, couples and hyperbolicity away from critical points.***

Remark. For real f satisfying LyapHyp and $K = \hat{I}$, we have the unique zero of pressure $t_0 = 1$ and for $-\log |f'|$ we conclude that a unique equilibrium state exists which is a.c.i.m. .

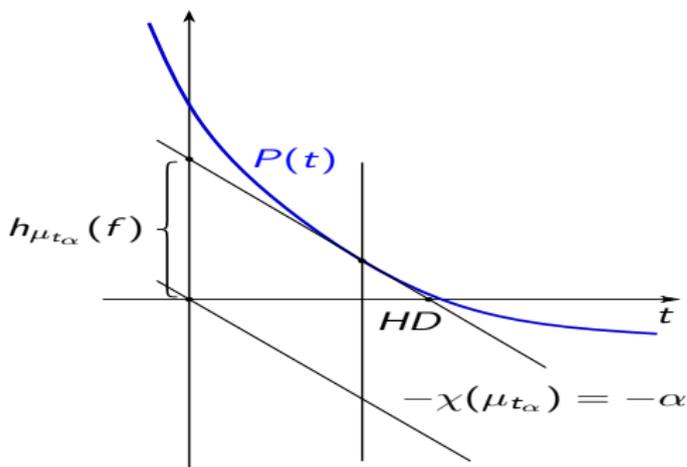
In general it holds assumed e.g. $|(f^n)'(f(c))| \rightarrow \infty$ for all $c \in \text{Crit}(f)$ (Bruin & Rivera-Letelier & Shen & van Strien: Inv. math 2008). For $t > t_+$, LyapHyp, equilibria do not exist (Rivera-Letelier & Inoquio 2012).

Proofs use inducing (Lai-Sang Young towers). For a different proof, the real case, see a recent preprint by Dobbs and Todd.

$P_{\text{var}}(t)$ allows to study **dimension spectrum for Lyapunov exponent** via Legendre transformation, proving in particular

$$\text{HD}(\{x \in K : \chi(x) = \alpha\}) = \frac{1}{|\alpha|} \inf_{t \in \mathbb{R}} (P(t) + \alpha t).$$

Proof of \geq . Given α consider t where inf is attained. The tangent to $P(t)$ at t is parallel to $-\alpha t$ and for μ_t the equilibrium, it is $h_{\mu_t}(f) - t\chi(\mu_t)$. So the infimum is $h_{\mu_t}(f)$, see Fig. (By variational definition $P(t)$ and h_{μ} are mutual Legendre type transforms.) Dividing by α gives \geq using Mañé's equality **$\text{HD}(\mu) = h_{\mu}(f)/\chi(\mu)$** .



Proof of \leq uses conformal measures.

Use of the Legendre transform of $P(t)$ allows also to give formulas for HD of irregular sets

$$\text{HD}(\{\underline{\chi}(x) = \alpha, \bar{\chi}(x) = \beta\})$$

for $\beta > 0$ [Gelfert & FP & Rams: Math. Ann. 2010, ETDS 2016].

In analogy to $\chi(\mu) \geq 0$ one has:

Theorem (Levin & FP & Shen: Inv.math. 2016))

If for a rational function $f : \mathbb{C} \rightarrow \mathbb{C}$ there is only one critical point c in $J(f)$ and no parabolic periodic orbits, then $\underline{\chi}(f(c)) \geq 0$.

For S -unimodal maps of interval this was proved much earlier by Nowicki and Sands.

6 – Other definitions of geometric pressure

Definition (tree pressure)

For every $z \in K$ and $t \in \mathbb{R}$ define

$$P_{\text{tree}}(z, t) = \limsup_{n \rightarrow \infty} \frac{1}{n} \log \sum_{f^n(x)=z, x \in K} |(f^n)'(x)|^{-t}.$$

Theorem

$P_{\text{tree}}(z, t)$ does not depend on z for z safe.

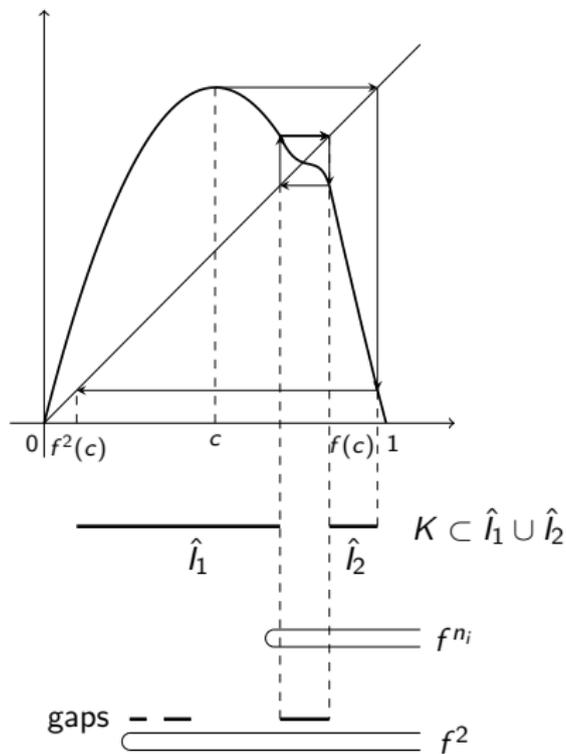
- In the complex case to prove $P_{\text{tree}}(z_1, t) = P_{\text{tree}}(z_2, t)$ one joins z_1 to z_2 with a curve not fast accumulated by critical trajectories, FP: TAMS 1999, FP & Rivera-Letelier & Smirnov: ETDS 2004.
- In the real case there is no room for such curves. Instead, one relies on topological transitivity, FP & Rivera-Letelier: arXiv 2014 & Memoir AMS 2019, FP: Monatsh. Math. 2018.

- For $\phi = -t \log |f'|$ pressure via separated sets does not make sense. Indeed, in presence of critical points for f , it is equal to $+\infty$. So it is replaced by P_{tree} .
- One can consider however *spanning geometric pressure* $P_{\text{span}}(t)$ using (n, ε) -spanning sets and infimum. Assumed *weak backward Lyapunov stability* it is indeed equal to $P(t)$ in the complex case (FP: Monatsh. Math. 2018).
- This is not so in the real case (where wbls always holds if all periodic orbits hyperbolic repelling). It happens $P_{\text{span}}(t) = \infty$ if some x with big $|(f^n)'(x)|^{-1}$ is *well ρ_n -isolated*.

Definition (weak backward Lyapunov stability, wbls)

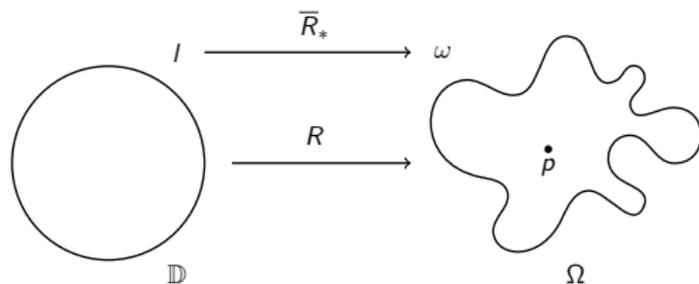
f is *weakly backward Lyapunov stable* if for every $\delta > 0$ and $\varepsilon > 0$ for all n large enough and every disc $B = B(x, \exp -\delta n)$ centered at $x \in K$, for every $0 \leq j \leq n$ and every component V of $f^{-j}(B)$ intersecting K , it holds that $\text{diam } V \leq \varepsilon$.

Question. Does wbls hold for all rational maps?



7 – Boundary dichotomy

- Let $f : \overline{\mathbb{C}} \rightarrow \overline{\mathbb{C}}$ be a rational map with $\deg(f) \geq 2$. and let $\Omega = \Omega_p(f)$ be a simply connected immediate basin of attraction to a fixed point p . Let $R : \mathbb{D} \rightarrow \Omega$ be a Riemann map $R(0) = p$ and $g : \mathbb{D} \rightarrow \mathbb{D}$ defined by $g := R^{-1} \circ f \circ R$, extended conformally beyond $\text{Fr } \Omega$ (Schwarz symmetry), thus expanding on $\partial\mathbb{D}$.
- Consider **harmonic measure** $\omega = \overline{R}_*(I)$, where I is normalized length measure on $\partial\mathbb{D}$ and \overline{R} is radial limit, defined I -a.e. I is g -invariant, hence ω is f -invariant. Denote by H_1 Hausdorff measure in dimension 1.



$$g = R^{-1} \circ f \circ R$$

f

Theorem (FP, Urbański, Zdunik: 1985 – 2006)

For f, Ω as above, $HD(\omega) = 1$. One of two cases holds:

- 1) $\omega \perp H_1$, which implies $HD_{\text{hyp}}(\text{Fr } \Omega) > 1$;
- 2) $\omega \ll H_1$ and f is a finite Blaschke product or a two-to-one holomorphic factor of a Blaschke product in some holomorphic coordinates on $\overline{\mathbb{C}}$.

Consider $\psi := \log |g'| - \log |f'| \circ \bar{R}$. Notice that

$$\int_{\partial\mathbb{D}} \psi \, dl = 0, \text{ hence } \text{HD}(\omega) = 1.$$

The latter was proved in 1985 by Makarov without assuming existence of f .

Consider the asymptotic variance

$$\sigma^2 = \sigma_\nu^2(\psi) := \lim_{n \rightarrow \infty} \frac{1}{n} \int_{\partial\mathbb{D}} (S_n \psi)^2 \, dl.$$

Then $\omega \perp H_1$ is equivalent to $\sigma^2 > 0$ and equivalent to ψ not being cohomologous to 0 (not of the form $u \circ f - u$).

Theorem (LIL-refined-HD for harmonic measure, FP, Urbański, Zdunik: Ann. Math. 1989, Studia Math. 1991)

For f, Ω with $\sigma^2 > 0$, there exists $c(\Omega) > 0$, such that for $\alpha_c(r) := r \exp(c \sqrt{\log 1/r \log \log \log 1/r})$

- i) $\omega \perp H_{\alpha_c}$ for the gauge function α_c , for all $0 < c < c(\Omega)$;*
- ii) $\mu \ll H_{\alpha_c}$ for all $c > c_1(\Omega)$.*

This theorem applies also e.g. to snowflake-type Ω 's,

Proofs.

We can find X with $\text{HD}(X) \geq \text{HD}(\omega) - \epsilon$ by Katok method and using $\text{HD} = h/\chi$. But we can do better:

$\sigma^2 > 0$ yields by CLT large fluctuations of the sums $\sum_{j=0}^{n-1} \psi \circ \varsigma^j$ from 0, allowing to find expanding X with $\text{HD}(X) > \text{HD}(\omega)$. One builds an iterated function system, for which X is the limit set. A special care is needed to get $X \subset \text{Fr } \Omega$.

Substituting in LIL $n \sim (\log 1/r_n)/\chi(\omega)$ for $r_n = |(f^n)'(x)|^{-n}$, comparing $\log |(g^n)'| - \log |(f^n)'| \circ \bar{R}$ with $\sqrt{2\sigma^2 n \log \log n}$ for a sequence of n 's, we get

Lemma (Refined Volume Lemma)

For ω -a.e. x

$$\limsup_{n \rightarrow \infty} \frac{\omega(B(x, r_n))}{\alpha_c(r_n)} = \begin{cases} \infty, & \text{for } 0 < c < c(\omega), \\ 0, & \text{for } c > c(\omega). \end{cases}$$

Using $R = f^{-n} \circ R \circ g^n$ one obtains

Theorem (radial growth)

For Lebesgue a.e. $\zeta \in \partial\mathbb{D}$

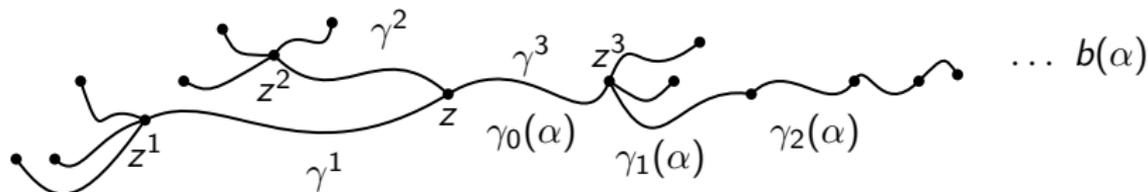
$$G^+(\zeta) := \limsup_{r \nearrow 1} \frac{\log |R'(r\zeta)|}{\sqrt{\log(1/1-r) \log \log \log(1/1-r)}} = c(\Omega).$$

Similarly $G^-(\zeta) := \liminf \dots = -c(\Omega)$.

Above theorems hold for every connected, simply connected open $\Omega \subset \mathbb{C}$, different from \mathbb{C} , without existence of f . Of course one should add ess sup over $\zeta \in \partial\mathbb{D}$ and over $z \in \text{Fr } \Omega$ in Refined Volume Lemma and reformulate the case i). There is a universal Makarov's upper bound $C_M < \infty$ for all $c(\Omega)$, $C_M \leq 1.2326$ (Hedenmalm, Kayumov: PAMS 2007). In 1989 I gave a weaker estimate.

Geometric coding trees, g.c.t.

- Above theorems hold in an abstract setting of a **geometric coding tree** in U for $f : U \rightarrow \overline{\mathbb{C}}$, $f(U) \supset U$ proper, giving a coding $\pi : \Sigma^d \rightarrow \Lambda$ to the limit set Λ (in place of $\overline{R} : \partial \mathbb{D} \rightarrow \text{Fr } \Omega$), provided f extends holomorphically beyond $\text{cl } \Lambda$ called then a **quasi-repeller**.



Curves $\gamma^j : [0, 1] \rightarrow f(U)$, $j = 1, \dots, d$, join z to z^j

$$\begin{aligned}\gamma_0(\alpha) &:= \gamma^{\alpha_0}, \\ f \circ \gamma_n(\alpha) &= \gamma_{n-1}(\varsigma(\alpha)), \\ \gamma_n(\alpha)(0) &= \gamma_{n-1}(\alpha)(1).\end{aligned}$$

- For a Hölder potential $\phi : \Sigma^d \rightarrow \mathbb{R}$ (in place of $-\log |g'|$) and Gibbs measure μ_ϕ one gets a dichotomy for $\mu := \pi_*(\mu_\phi)$ on Λ .
- For a constant potential $\mu = \mu_{\max}$ a measure of maximal entropy on Julia set $J(f)$ for $f : \overline{\mathbb{C}} \rightarrow \overline{\mathbb{C}}$ rational. Then
 - 1) If $\sigma^2 > 0$ then $\text{HD}_{\text{hyp}}(J(f)) > \text{HD}(\mu_{\max})$.
 - 2) If $\sigma^2 = 0$ then for each $x, y \in J(f)$ not postcritical, if $z = f^n(x) = f^m(y)$ for some positive integers n, m , the orders of criticality of f^n at x and f^m at y coincide. In particular all critical points in $J(f)$ are pre-periodic, f is postcritically finite with parabolic orbifold, in particular z^d , Chebyshev or some Lattès maps, (Zdunik, Inv. math. 1990).
- In the Ω version it is sufficient to assume f is defined only in a neighbourhood of $\partial\Omega$ repelling on the side of Ω , called **RB-domain**.
- This applies to f polynomial and simply connected $\Omega = \Omega_\infty$ giving again the dichotomy on $\text{Fr}\Omega$.

integral mean spectrum

- For a simply connected domain $\Omega \subset \mathbb{C}$ one considers the *integral means spectrum*:

$$\beta_{\Omega}(t) := \limsup_{r \nearrow 1} \frac{1}{|\log(1-r)|} \log \int_{\zeta \in \partial\mathbb{D}} |R'(r\zeta)|^t |d\zeta|.$$

This, in presence of f , e.g. for an RB-domain Ω and for $\phi = -\log |f'|$ for $g(z) = z^d$, e.g. Ω being a simply connected basin of ∞ for a polynomial of degree d , satisfies

$$\beta_{\Omega}(t) = t - 1 + \frac{P(t\phi)}{\log d}. \quad (\text{Makarov, FP \& Rohde})$$

One considers

$$\sigma^2(\log R') := \limsup_{r \nearrow 1} \frac{\int_{\partial\mathbb{D}} |\log R'(t\zeta)|^2 |d\zeta|}{-2\pi \log(1-r)}.$$

It holds $\sigma^2(\log R') = 2 \frac{d^2 \beta_{\Omega}(t)}{dt^2} \Big|_{t=0}$ (O. Ivrii). It is related to the Weil-Petersson metric (McMullen).

Recall $\sigma_{\mu}^2(t\phi) = \frac{d^2 P(f, t\phi)}{dt^2}$ for μ Gibbs in expanding case, Ruelle: Thermodyn. Formalism, FP & Urbański: Conformal Fractals.

8. Accessibility

Theorem (Douady-Eremenko-Levin-Petersen, accessibility of periodic sources; FP, the general case: Fund. Math. 1994)

Let Λ be a limit set for a g.c.t. \mathcal{T} for holomorphic $f : U \rightarrow \bar{\mathbb{C}}$. Assume $\text{diam}(\gamma_n(\alpha)) \rightarrow 0$, as $n \rightarrow \infty$, uniform shrinking with respect to $\alpha \in \Sigma^d$. Then every **good** $q \in \text{cl} \Lambda$ is a limit of a convergent branch $b(\alpha)$, i.e. $q \in \Lambda$. In particular, this holds for every q with $\underline{\chi}(q) > 0$ satisfying a **local backward invariance**.

Corollary (lifting of measure, FP 1994 & Proc. ICM18)

Every non-atomic hyperbolic probability measure μ , i.e. $\chi(\mu) > 0$, on $\text{cl} \Lambda$, is the π_* image of a probability ζ -invariant measure ν on Σ^d , assumed **uniform shrinking, \mathcal{T} has no self-intersections and μ -a.e. local backward invariance of U** . In part. a lift ν exists for every completely invariant RB-domain, e.g. for μ on $\text{Fr} \Omega_\infty$ for f polynomial.

THANK YOU FOR YOUR ATTENTION

